



Max-Planck-Gesellschaft

vibrational and rotational spectroscopy on surfaces

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Faradayweg 4-6
14195 Berlin

Outline



Max-Planck-Gesellschaft

- Introduction
- electron energy loss spectroscopy (EELS)
- infrared spectroscopy
- helium atom scattering (HAS)
- tip enhanced Raman spectroscopy/SERS
- SFG (sum frequency generation)
- electron spin resonance

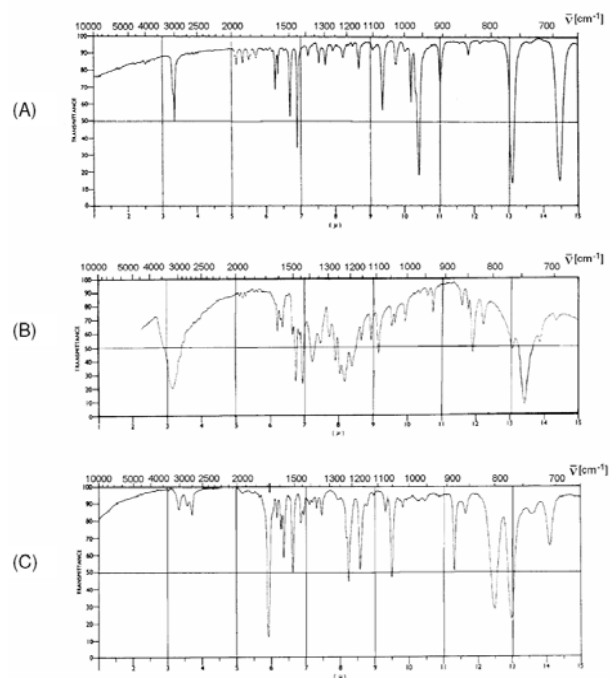
Introduction



Max-Planck-Gesellschaft

chemical perspective

Aufgabe 2: Die dargestellten IR-Spektren können von 1-Naphthaldehyd, E-Zimtaldehyd, trans-Stilben, 2,4,6-Tri-*t*-butylphenol, 2,2'-Dihydroxybiphenyl und 4,4'-Dihydroxybiphenyl stammen. Welches Spektrum gehört zu welcher Substanz? Ordnen Sie die charakteristischen Banden zu.

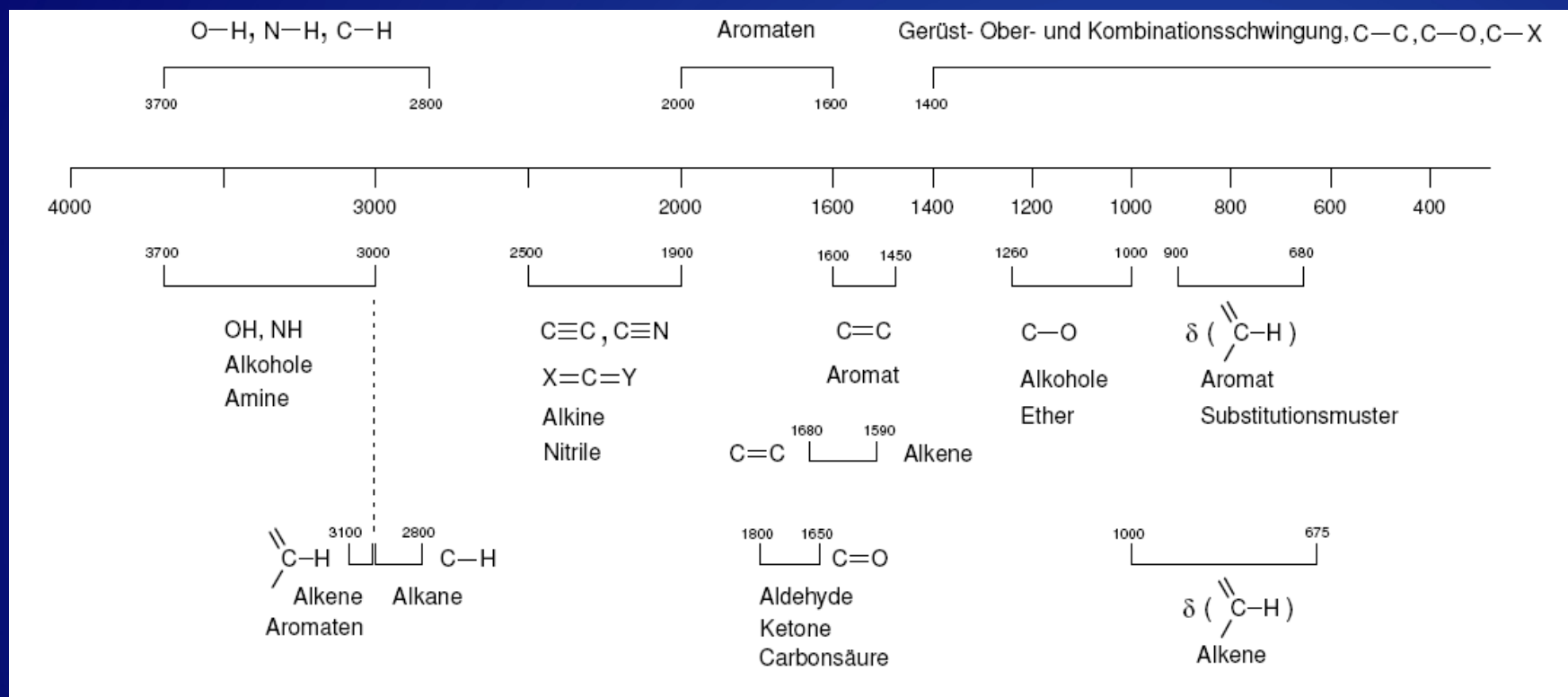


Introduction



Max-Planck-Gesellschaft

chemical perspective



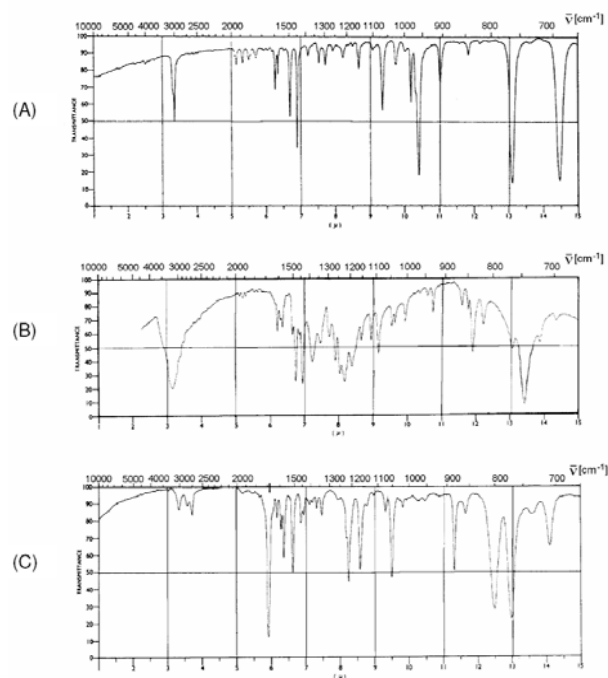
Introduction



Max-Planck-Gesellschaft

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physical perspective

PHYSICAL REVIEW B 66, 073414 (2002)

Anomalous dispersion of adsorbate phonons of Mo(110)-H

Jörg Kröger*

Institut für Experimentelle und Angewandte Physik, Christian-Albrechts-Universität zu Kiel, D-24098 Kiel, Germany

Sieghart Lehwald, Martin Balden, and Harald Ibach

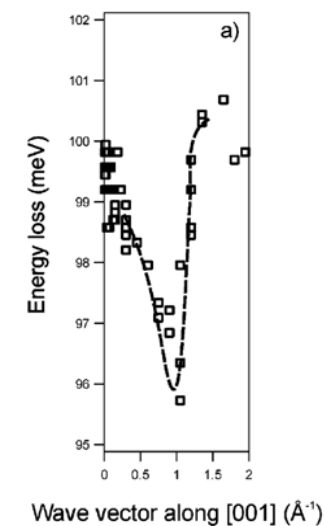
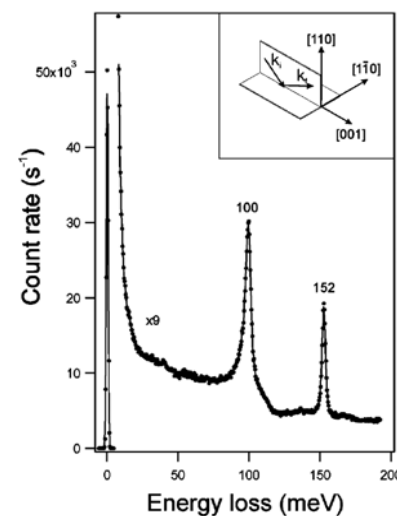
Institut für Schichten und Grenzen, Forschungszentrum Jülich, D-52425 Jülich, Germany

(Received 8 November 2001; published 21 August 2002)

The dispersion curve of the longitudinal-optical adsorbate phonon on hydrogen-saturated Mo(110) along [001] is found to exhibit an anomalous indentation. The maximum indentation is observed at a wave vector, which coincides within the experimental angular resolution with the wave vector, at which the known giant Kohn anomaly for the transverse- and longitudinal-acoustic substrate surface phonons along [001] occurs.

DOI: 10.1103/PhysRevB.66.073414

PACS number(s): 68.35.Ja, 63.20.Kr, 68.43.Pq, 73.20.At

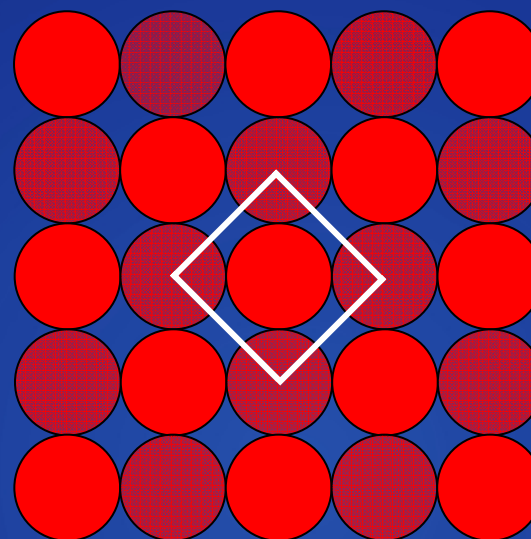


EELS

H on W(100)



Max-Planck-Gesellschaft



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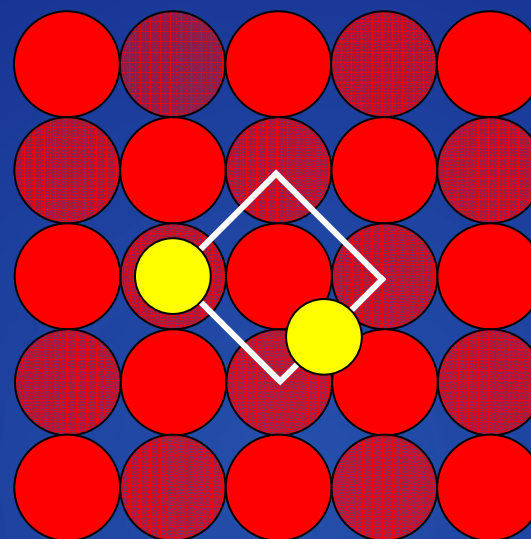
 2. layer

EELS

H on W(100)



Max-Planck-Gesellschaft



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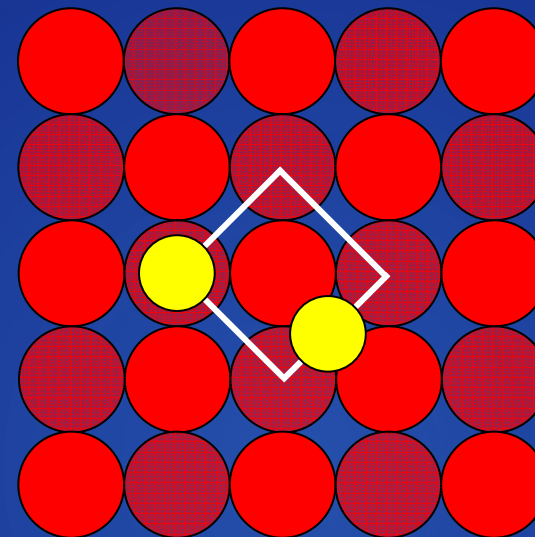
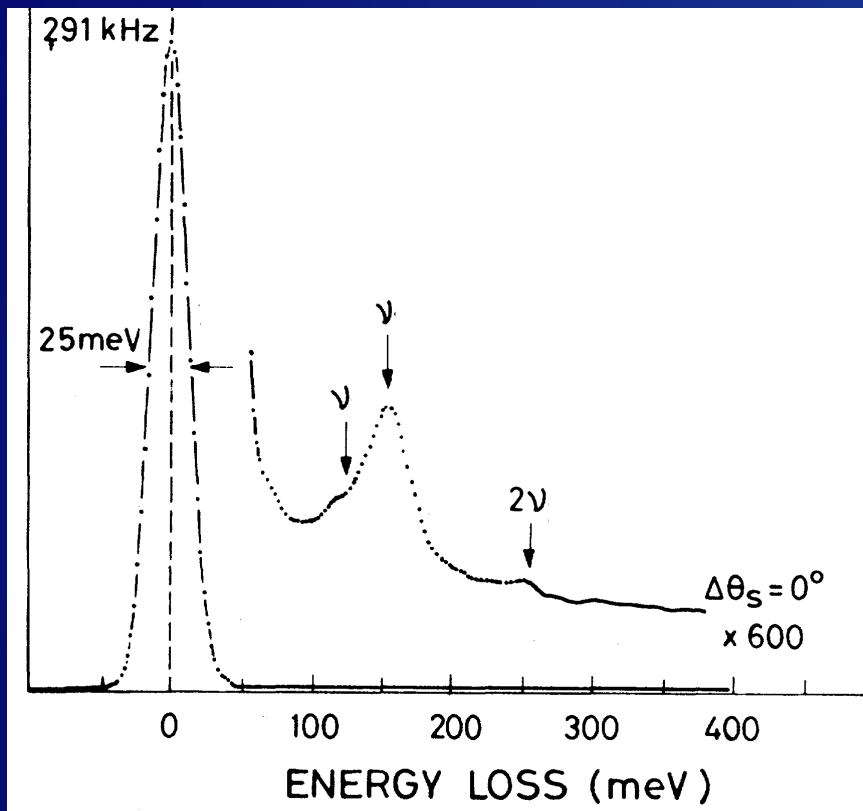
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EELS

H on W(100)



Max-Planck-Gesellschaft



1. layer

2. layer

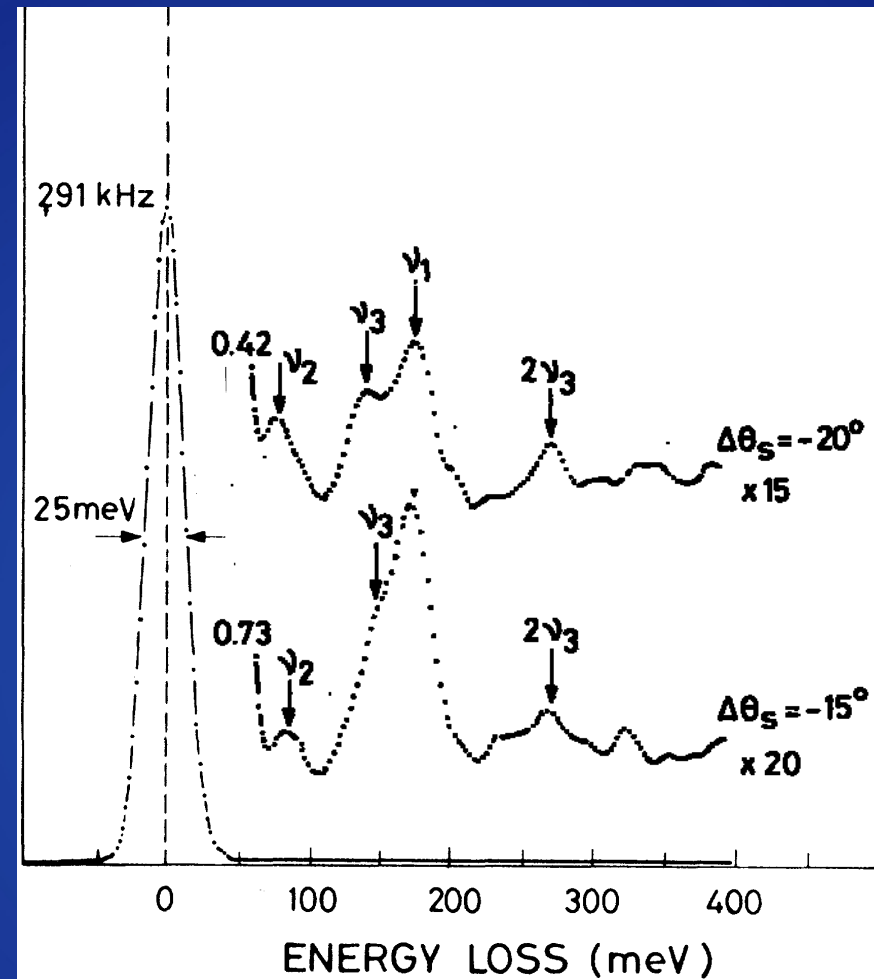
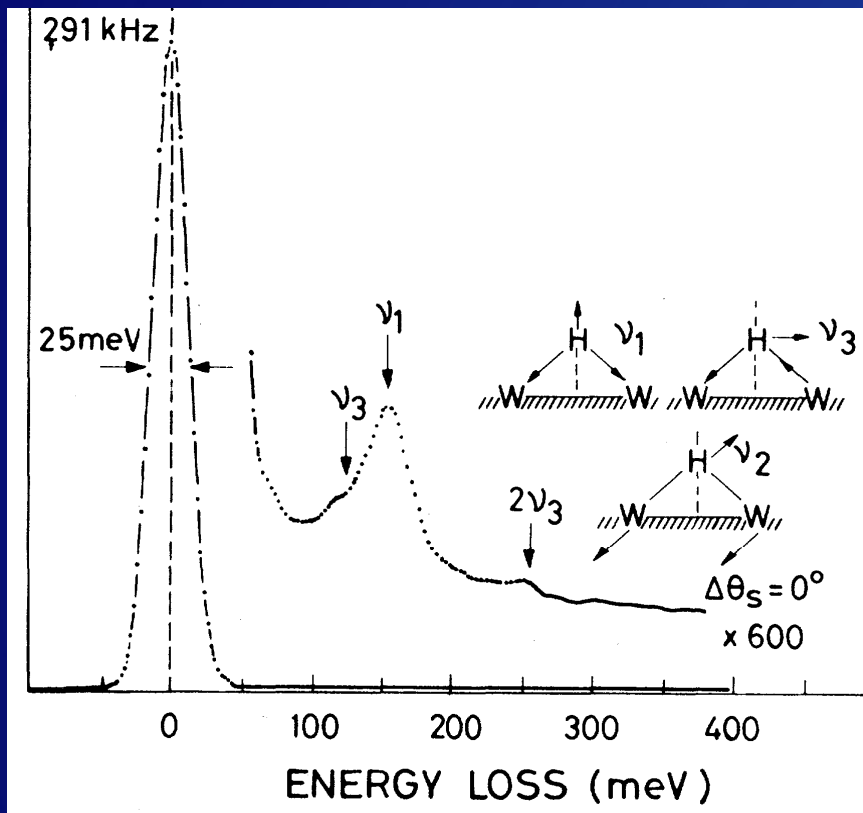
W. Ho, R. F. Willis, and E. W. Plummer, Phys. Rev. Lett. 40, 1463 (1978).

EELS

H on W(100)



Max-Planck-Gesellschaft



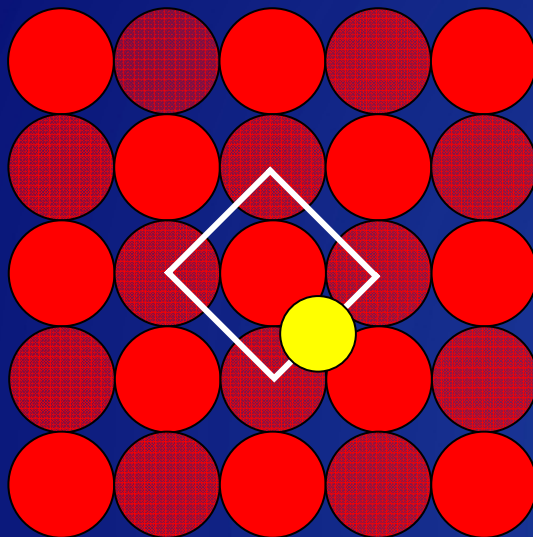
W. Ho, R. F. Willis, and E. W. Plummer, Phys. Rev. Lett. 40, 1463 (1978).

EELS

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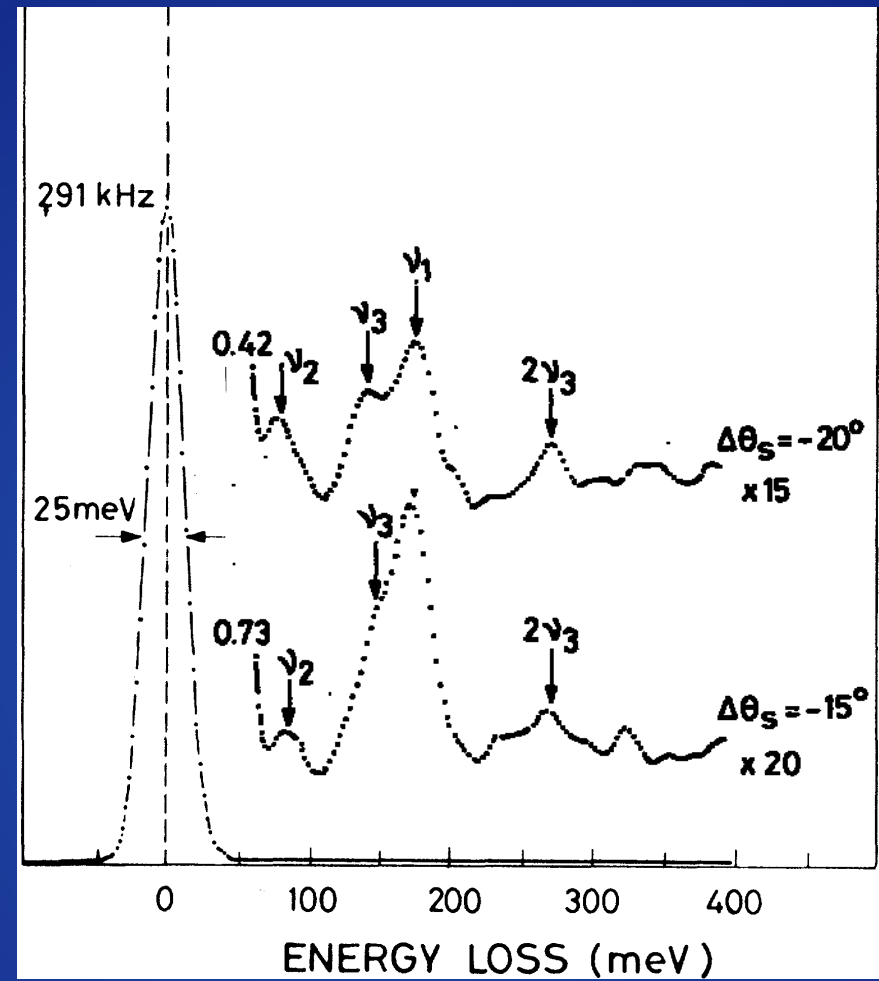


Max-Planck-Gesellschaft



 1. layer

 2. layer

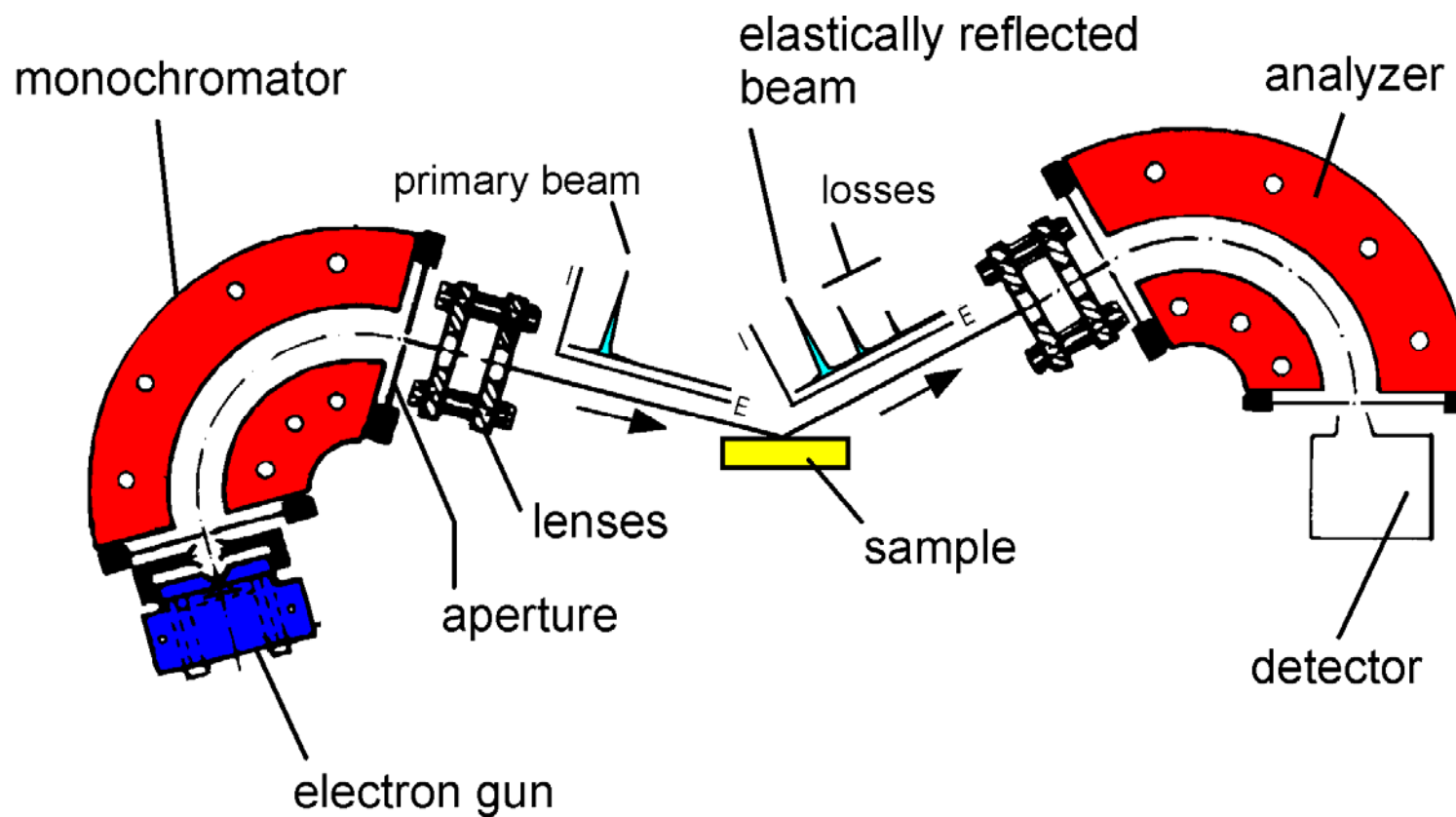


EELS

principle setup



Max-Planck-Gesellschaft



EELS

principles



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Why EELS? (pros and cons)

- lower energy vibrations detectable (down to approx. 100 cm⁻¹)
- dipole selection rule applies only in specular geometry
- off specular geometry allows detection of non IR active modes
- energy resolution to the best 4 cm⁻¹ (0.5 meV)
- electron current on the sample decreases with resolution in a power law behavior :
($I \propto (\Delta E - \Delta E_{\min})^{5/2}$)

EELS

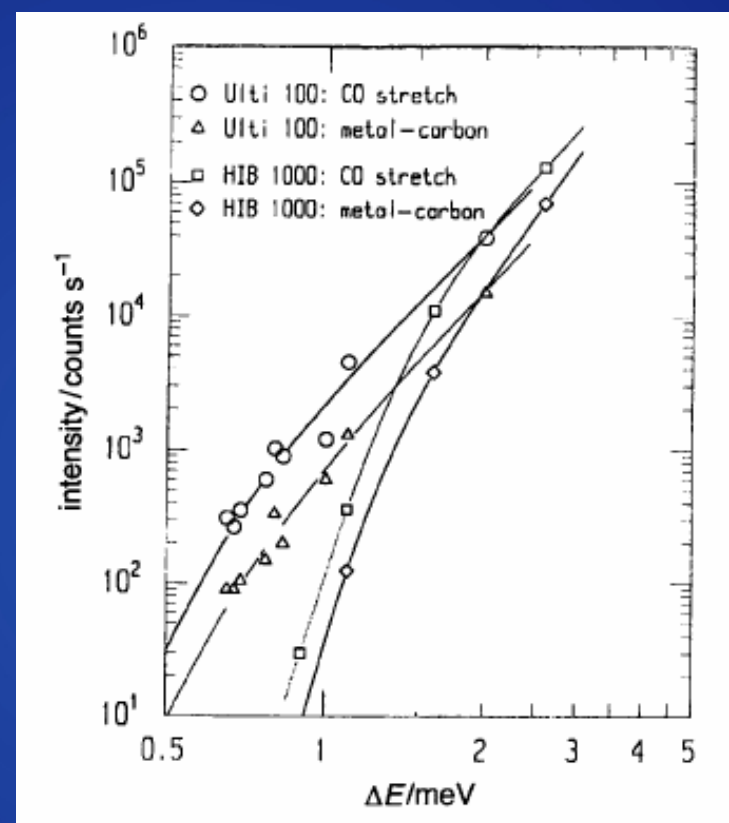
principles



Max-Planck-Gesellschaft

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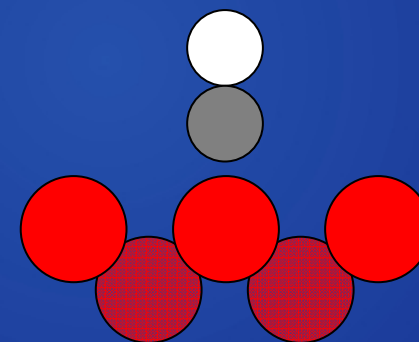
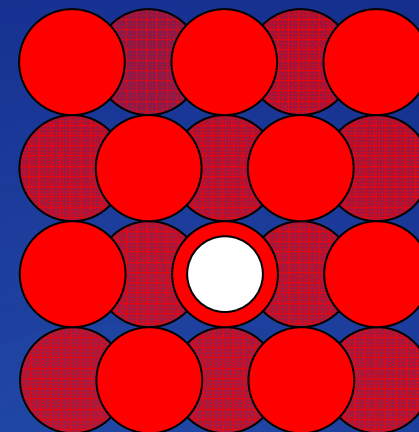
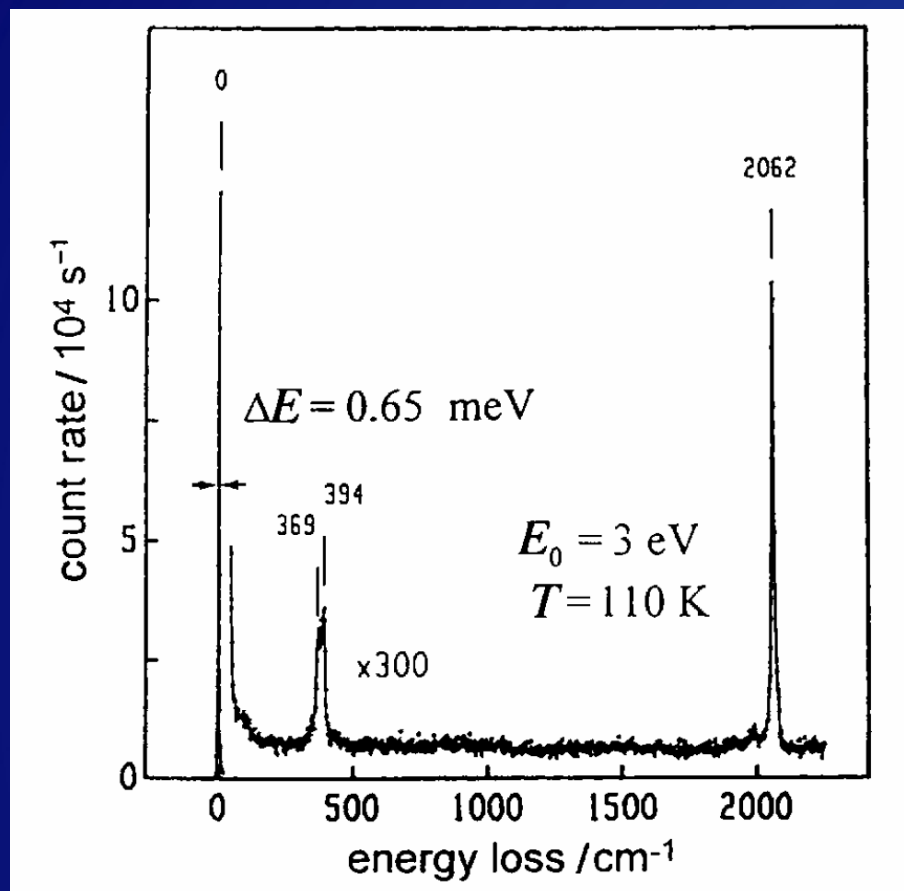
H. Ibach, M. Balden, S. Lehwald, J. Chem. Soc.-Faraday Trans. **92**, 4771 (1996).

EELS

CO/W(110)



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H. Ibach, M. Balden, S. Lehwald, J. Chem. Soc.-Faraday Trans. **92**, 4771 (1996).

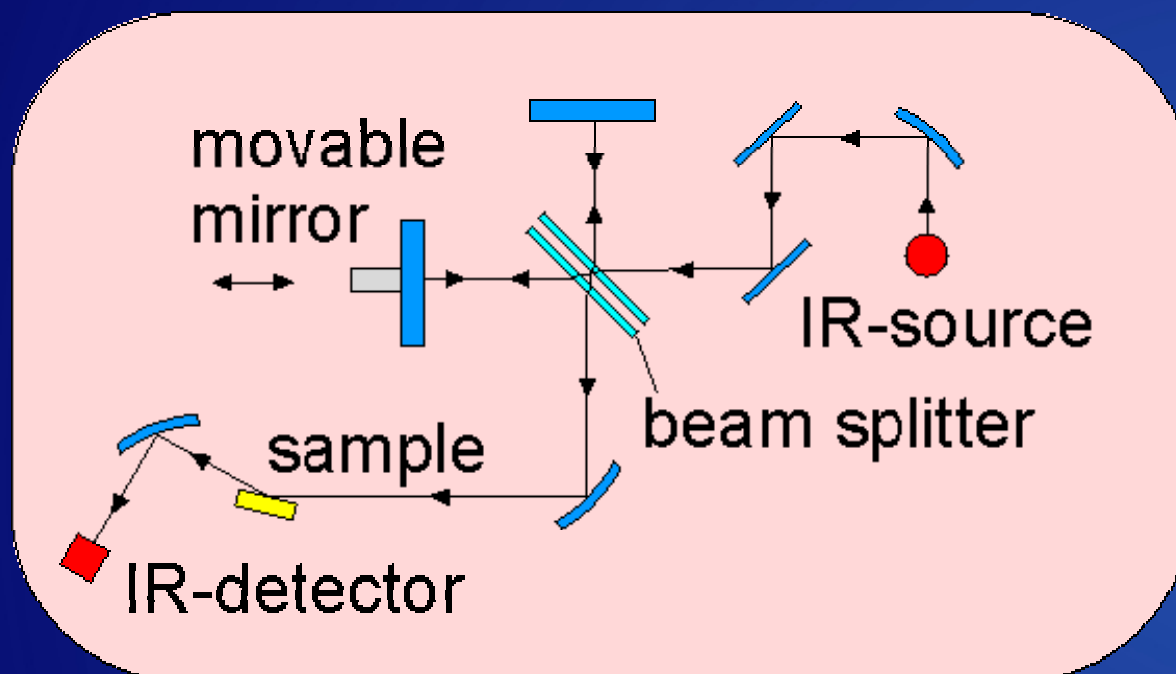
surface vibrations

IRAS



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principle experimental setup



source (white source)

- globar
- Nernst rod

detector

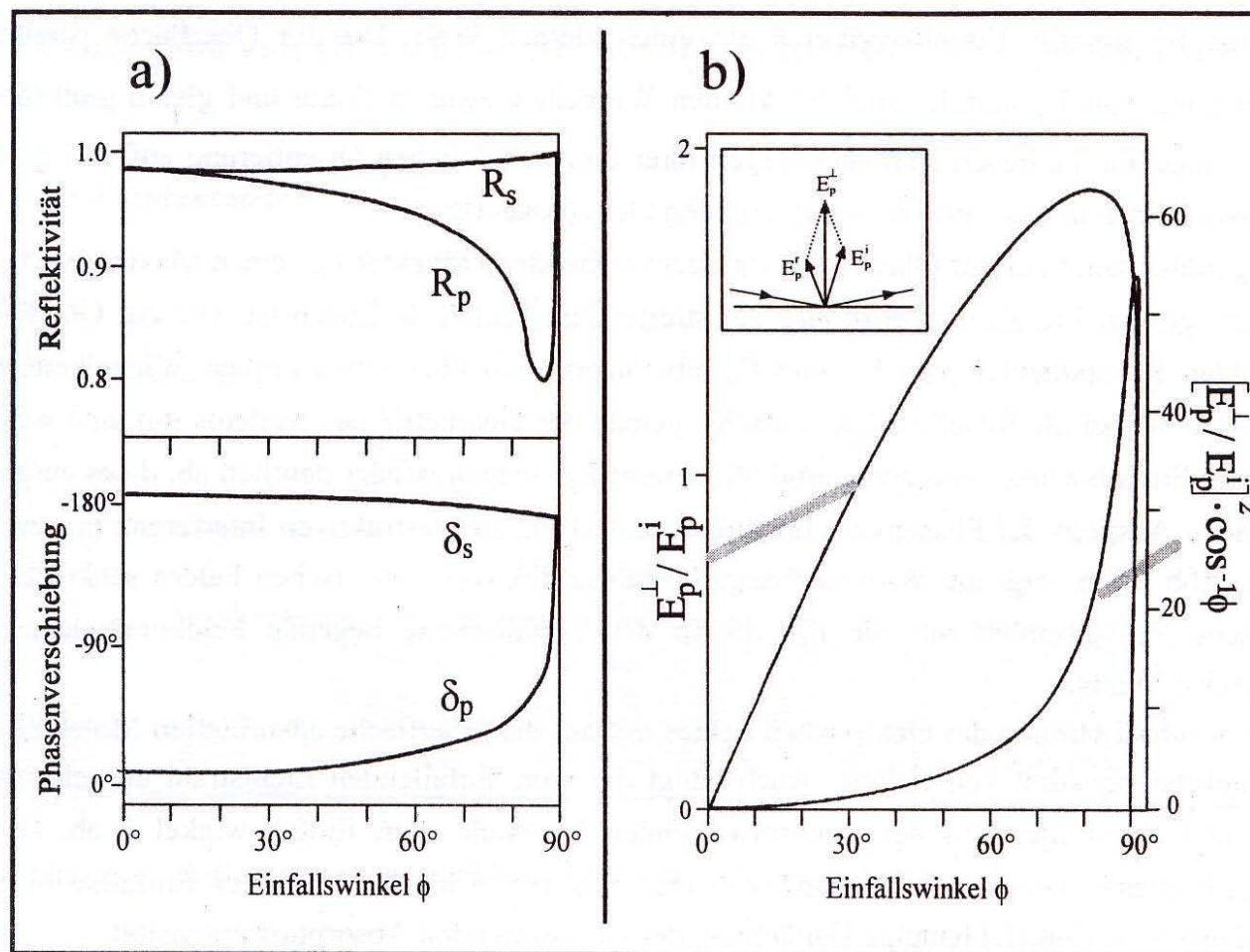
- semiconductor, broad band e.g. MCT, InSb, GeCu
- bolometer (higher sensitivity especially for low frequencies; narrow band)

surface vibrations

IRAS



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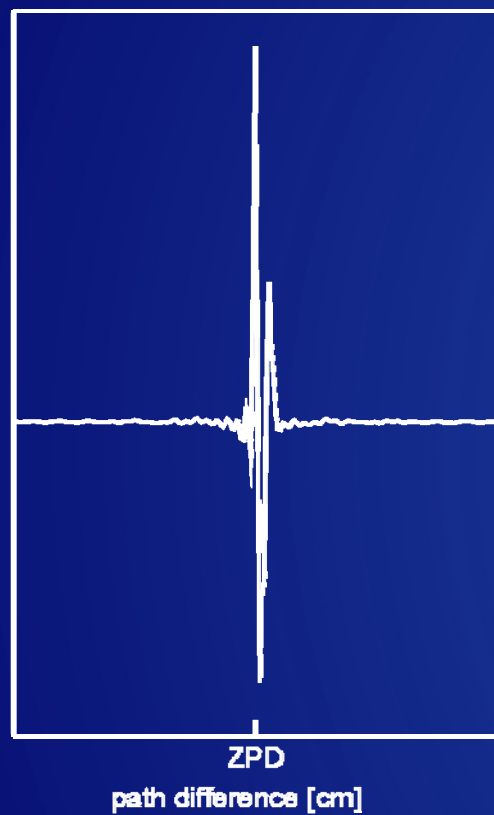
surface vibrations

IRAS



Max-Planck-Gesellschaft

interferogram



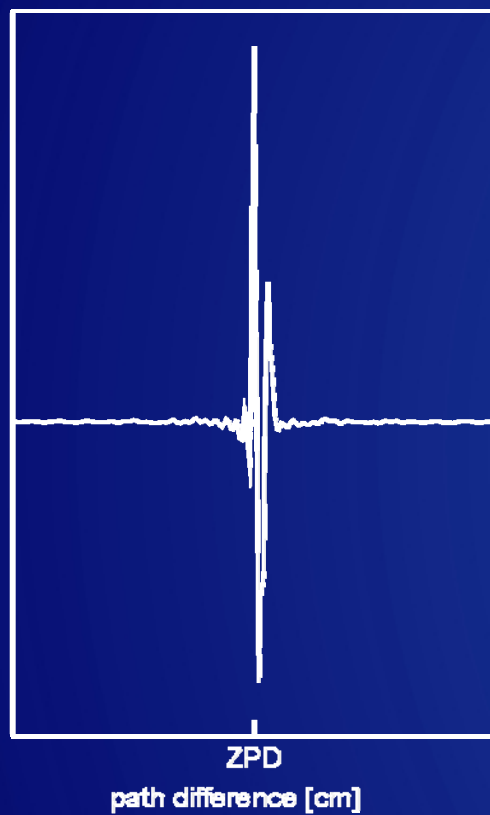
surface vibrations

IRAS



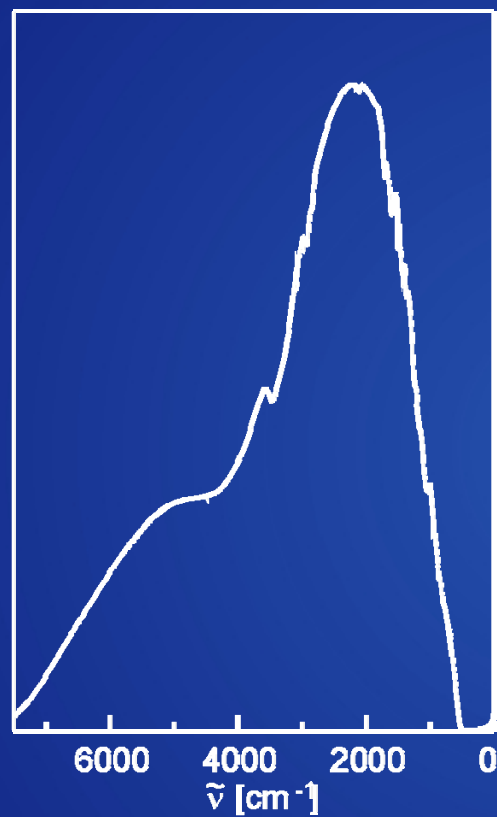
Max-Planck-Gesellschaft

interferogram



FFT

single channel
spectrum



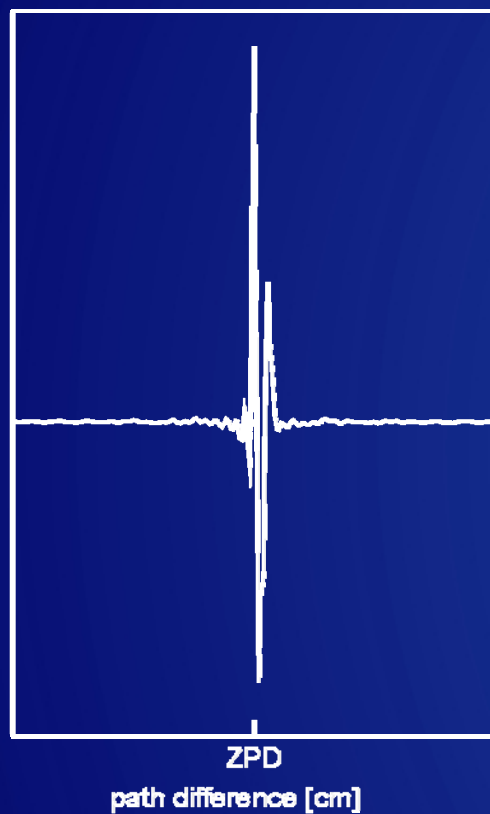
surface vibrations

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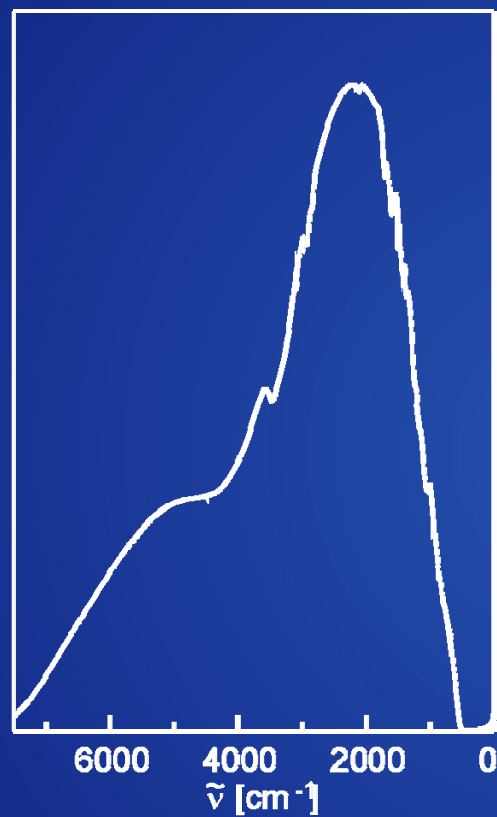
Max-Planck-Gesellschaft

interferogram

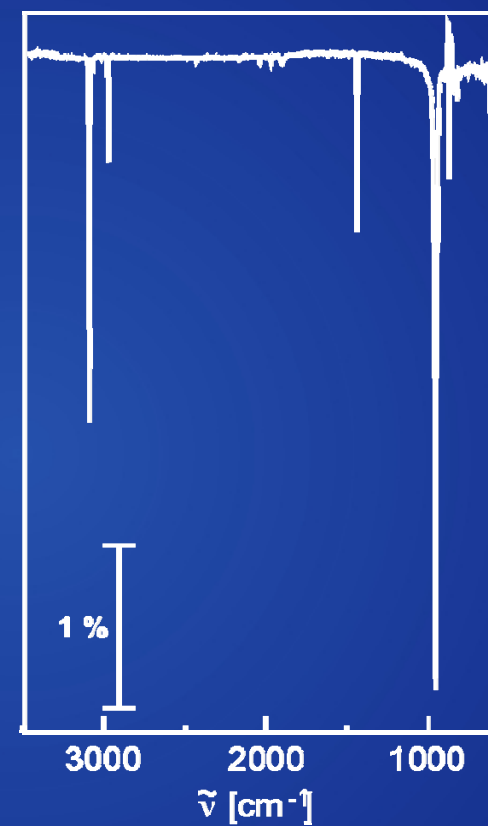


FFT

single channel
spectrum



transmission
spectrum

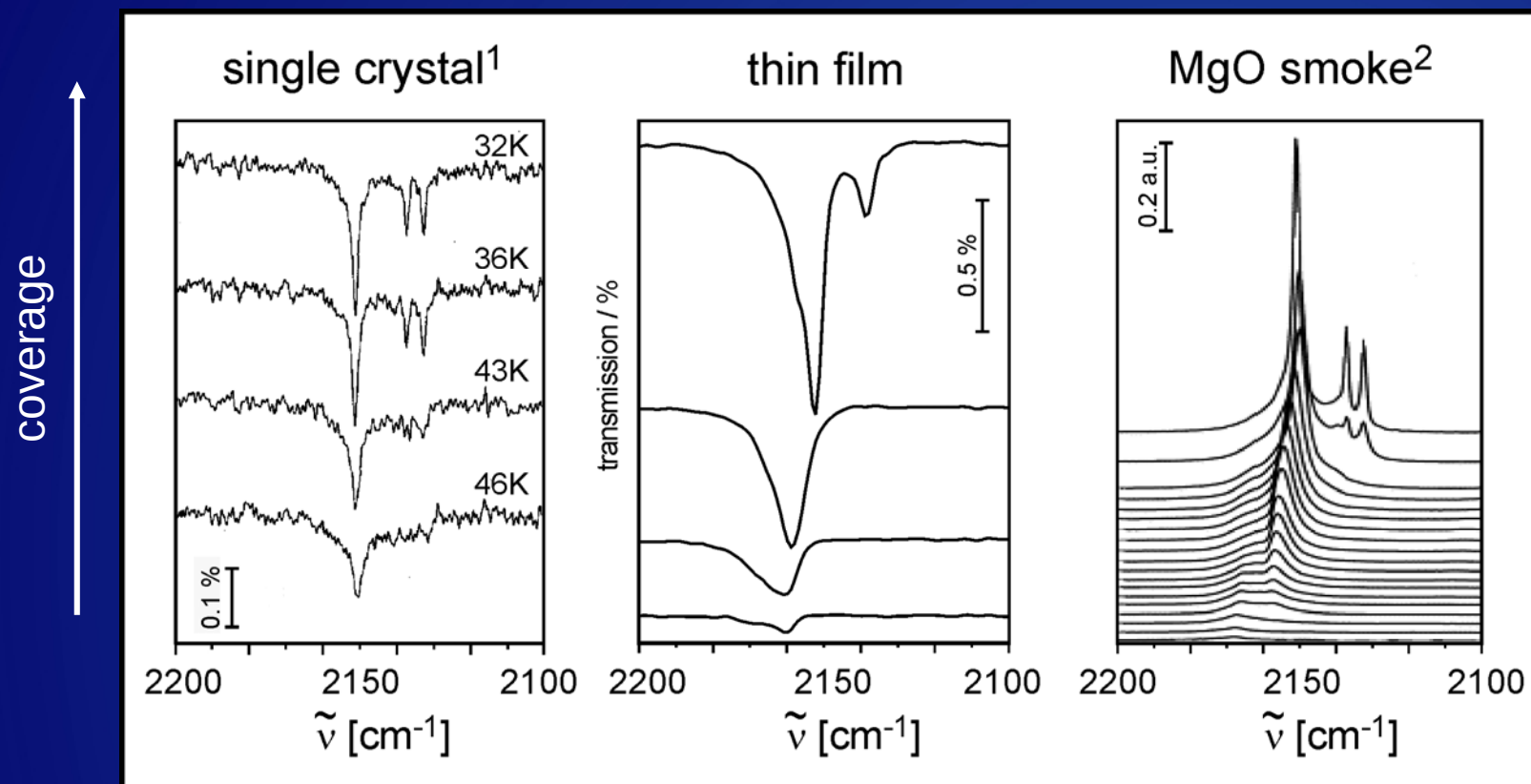


IRAS

CO on MgO(001)



Max-Planck-Gesellschaft



¹ J. Heidberg et al. *Surf. Sci.* **331-333**, 1467 (1995).

² G. Spoto et al. *Prog. Surf. Sci.* **76**, 71, (2004).

M. Sterrer et al. *Surf. Sci.* **596**, 222 (2005).

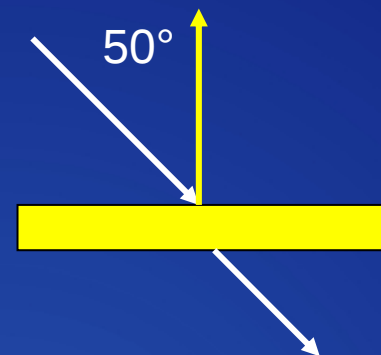
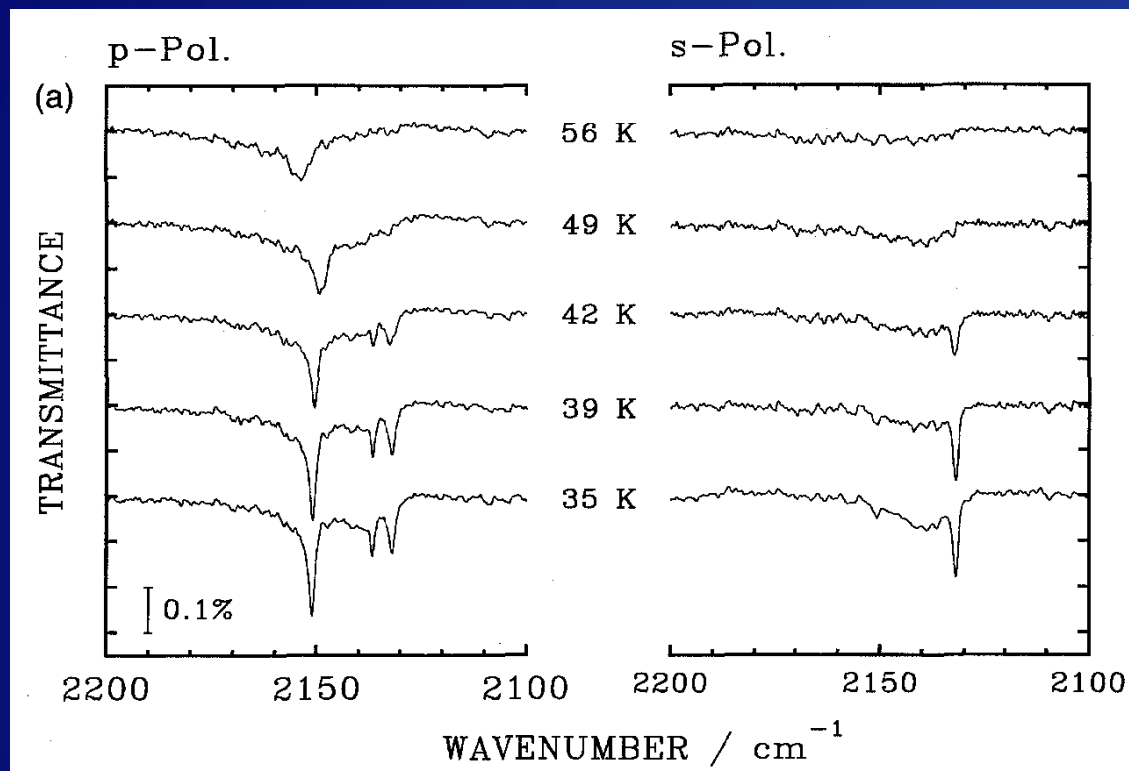
IRAS

CO on MgO(001)



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transmission experiment



- c(4x2) superstructure (LEED)
- 3 molecules in the unit cell (LEED, HAS)
- 1 molecule perpendicular
- 2 molecules tilted

¹ J. Heidberg et al. *Surf. Sci.* **331-333**, 1467 (1995).

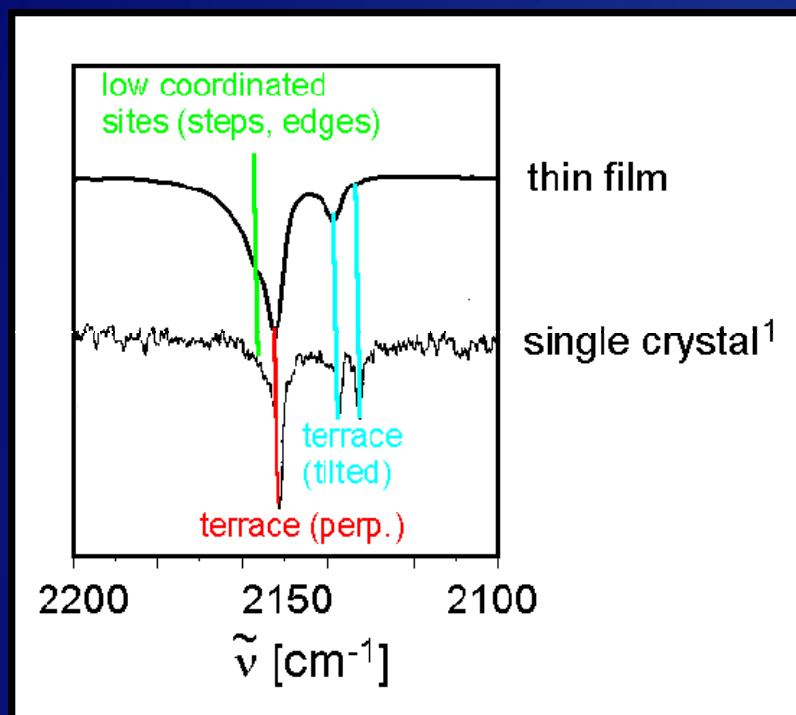
IRAS

CO on MgO(100)



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IR



¹ J. Heidberg et al. *Surf. Sci.* **331-333**, 1467 (1995).

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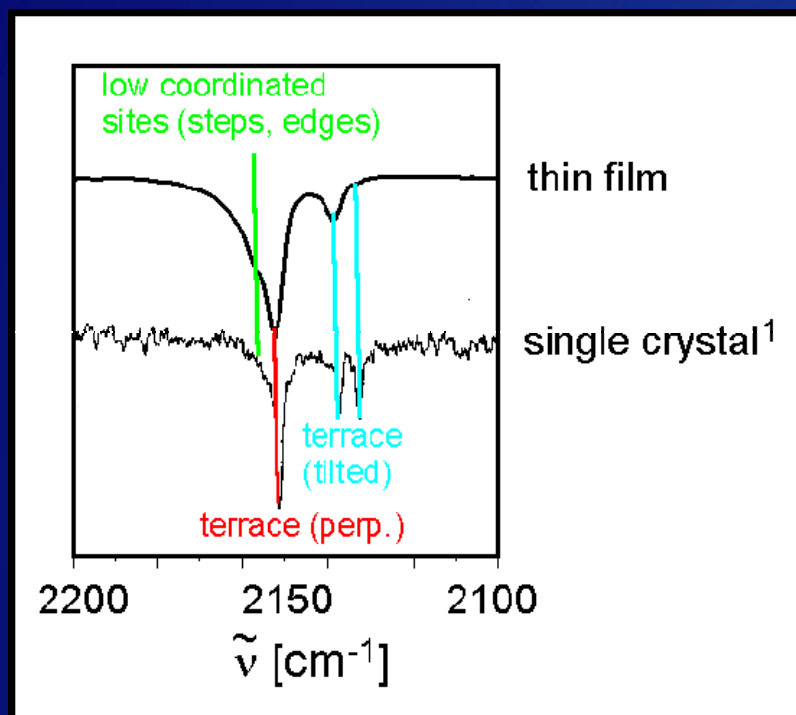
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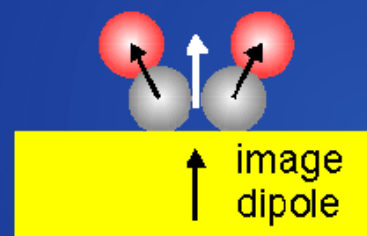
IR



2 tilted molecules

symmetric stretch
high frequency

dynamic dipole



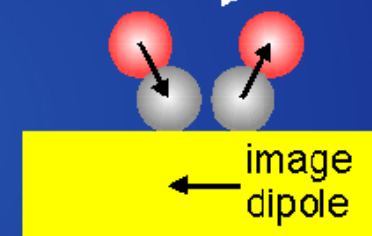
metal

$$\mu_{\text{dyndip}} > 0$$

IR active

asymmetric stretch
low frequency

dynamic dipole



metal

$$\mu_{\text{dyndip}} = 0$$

IR inactive

¹ J. Heidberg et al. *Surf. Sci.* **331-333**, 1467 (1995).

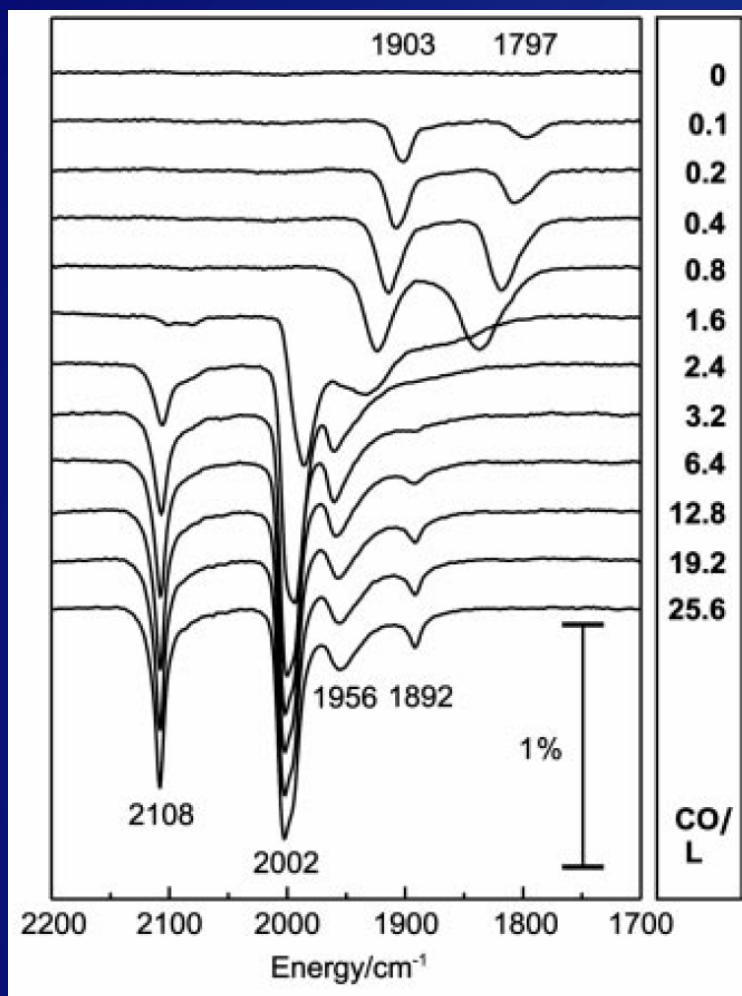
M. Sterrer et al. *Surf. Sci.* **596**, 222 (2005).

Pd particles on thin alumina

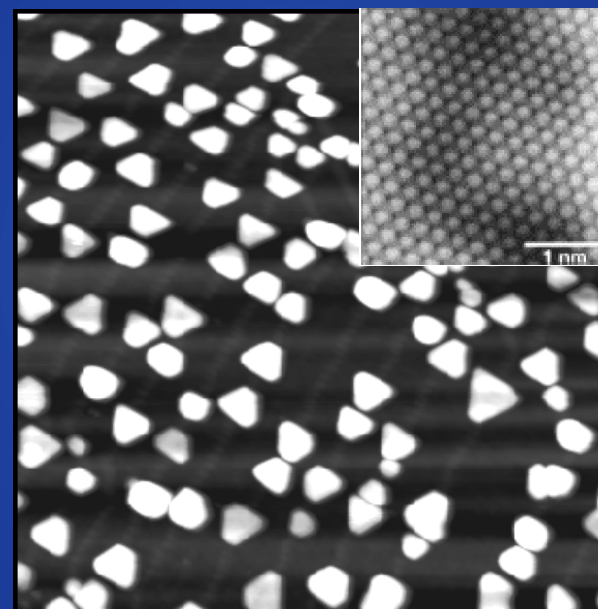
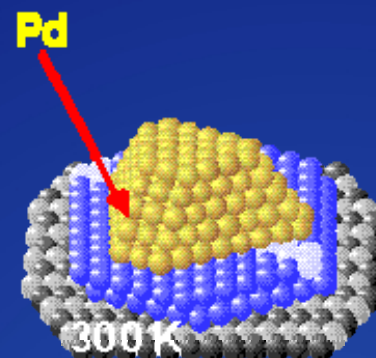
CO IRAS



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M. Frank and M. Bäumer, PCCP **2**, 3723 (2000).



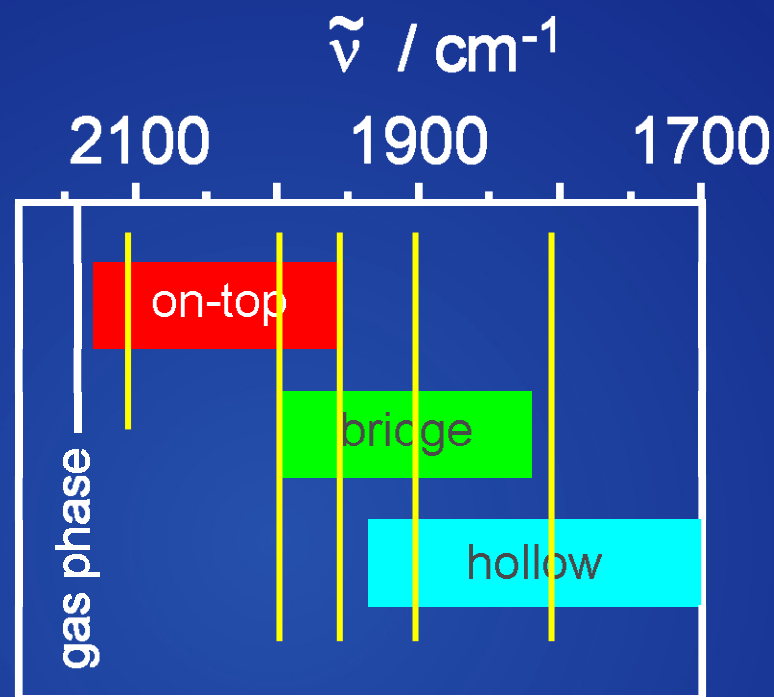
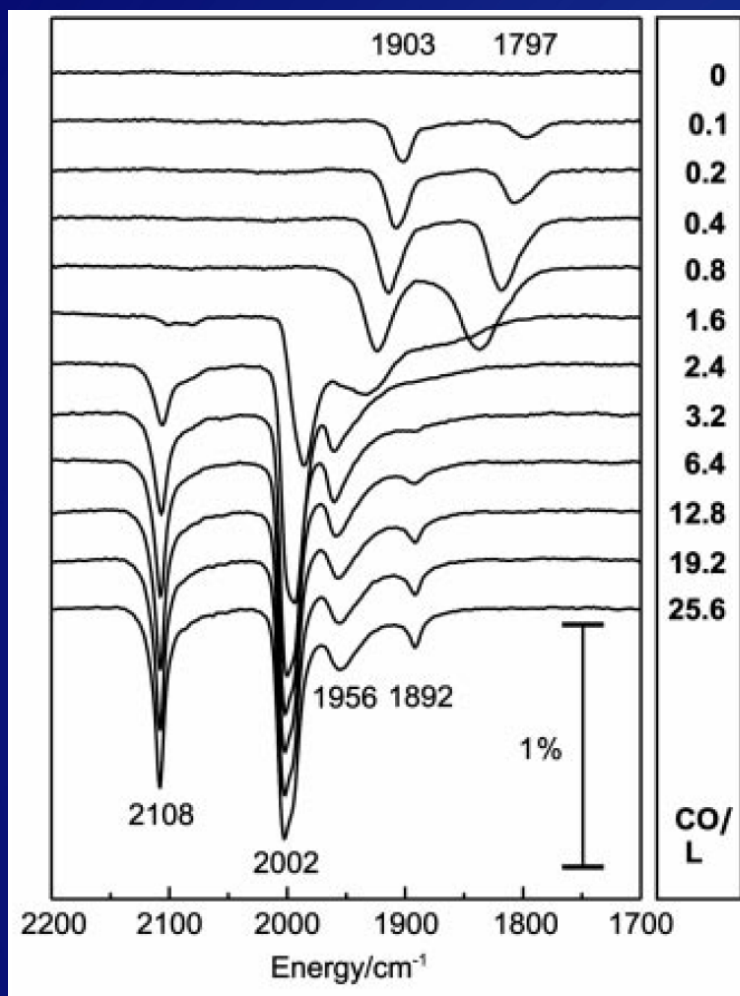
K. H. Hansen *et al.* Phys. Rev. Lett. **83**, 4120 (1999).

Pd particles on thin alumina

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Max-Planck-Gesellschaft



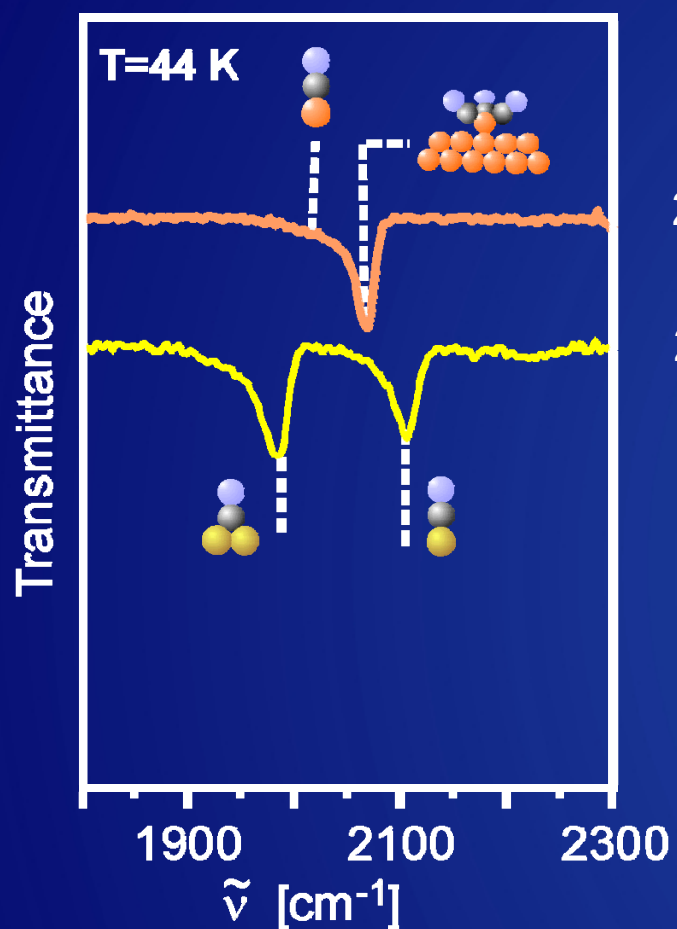
M. Frank and M. Bäumer, PCCP **2**, 3723 (2000).

Metal particles

CO IRAS

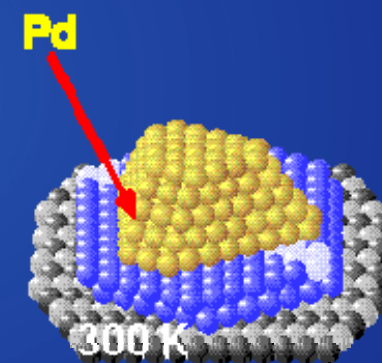
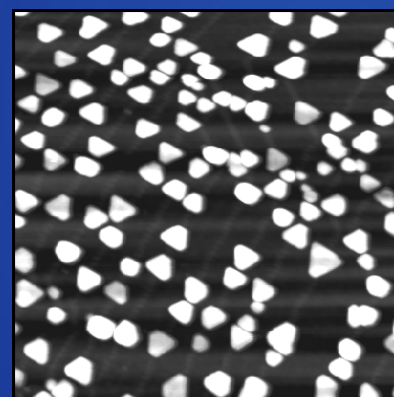
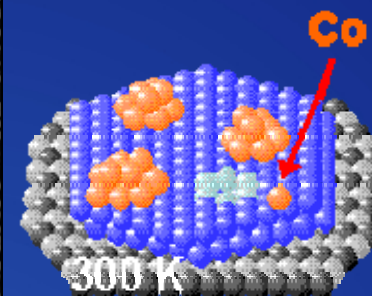
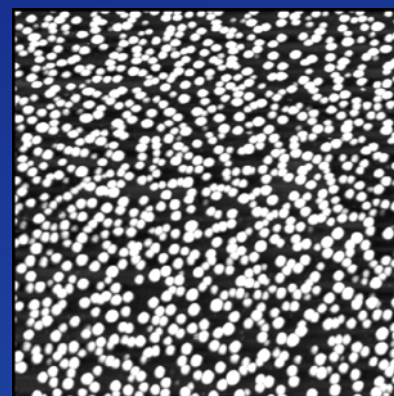


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2 Å Co

2 Å Pd

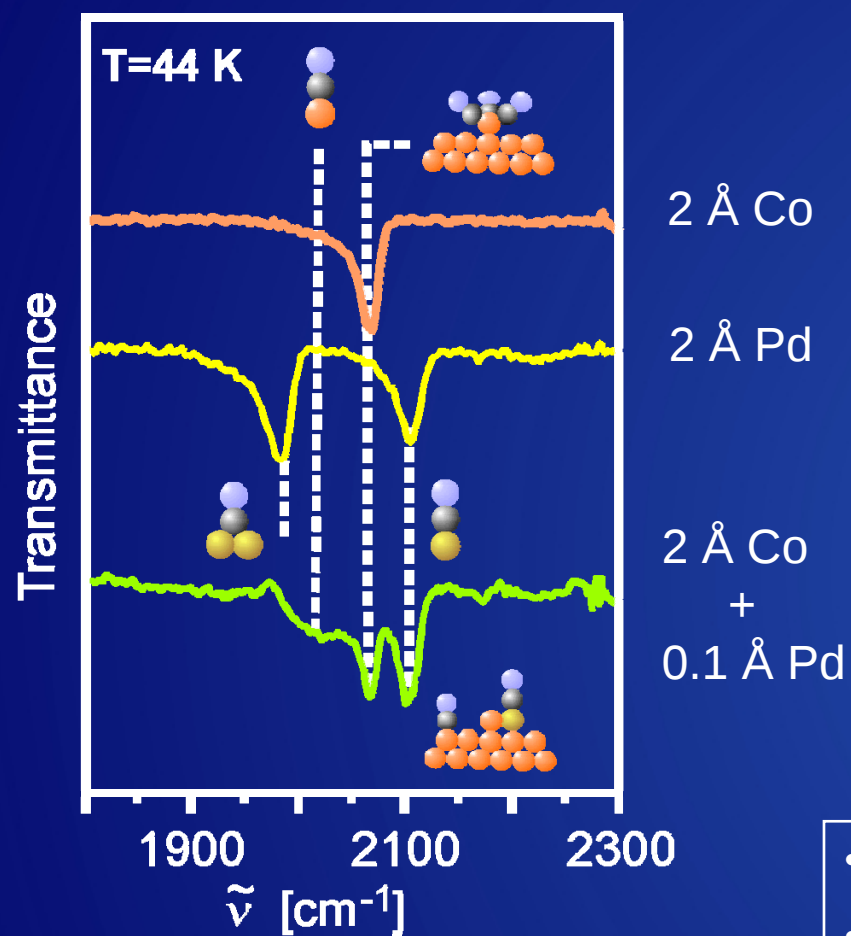


Metal particles

CO IRAS



Max-Planck-Gesellschaft



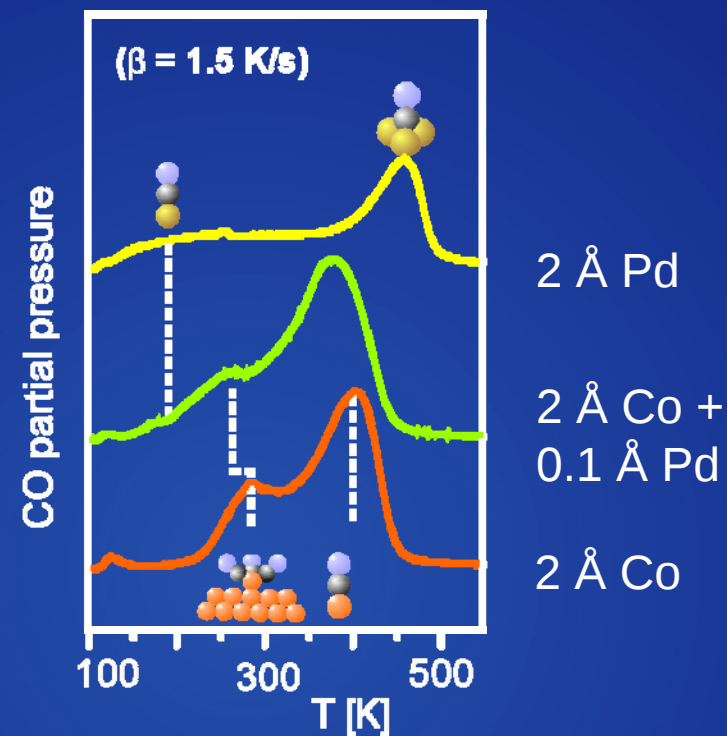
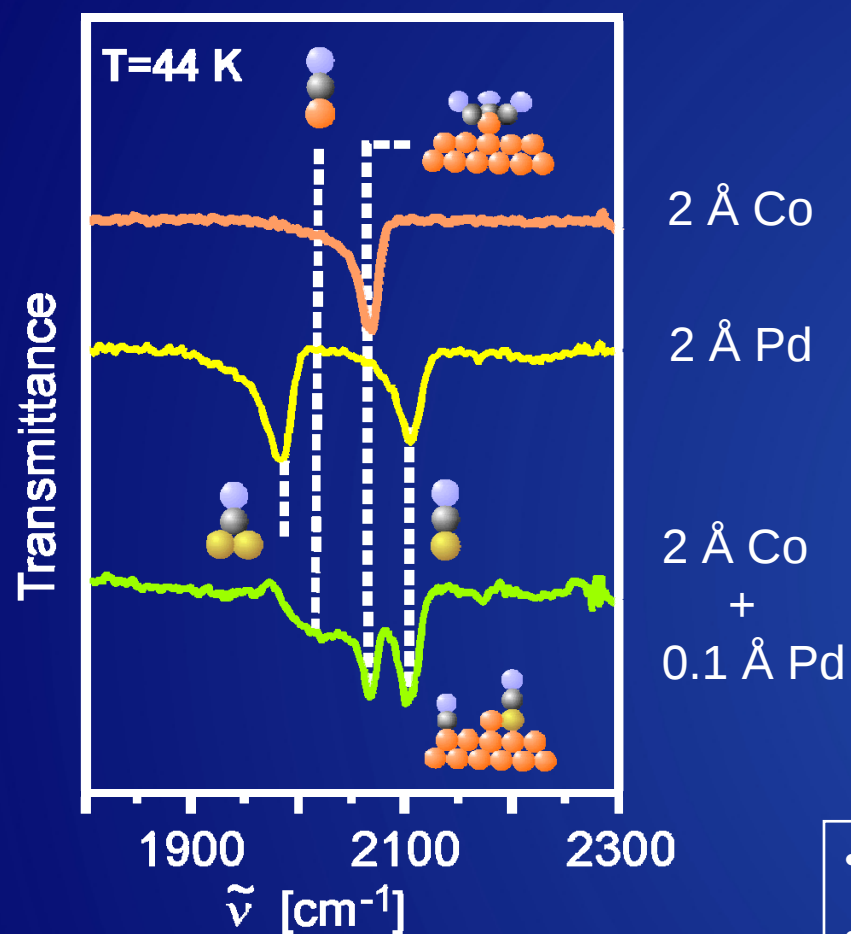
- Pd nucleates at low coordinated Co sites
- virtually no chemical contrast!

Metal particles

CO IRAS



Max-Planck-Gesellschaft



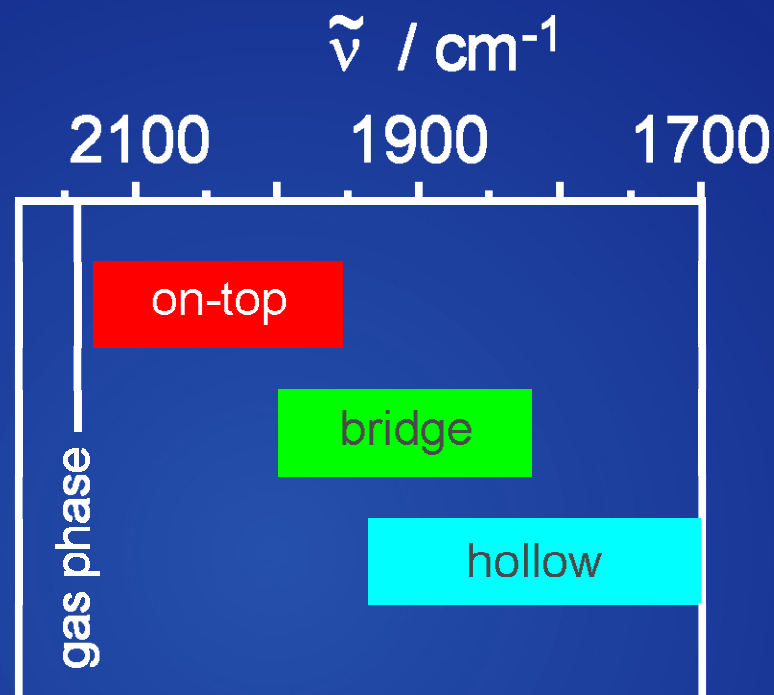
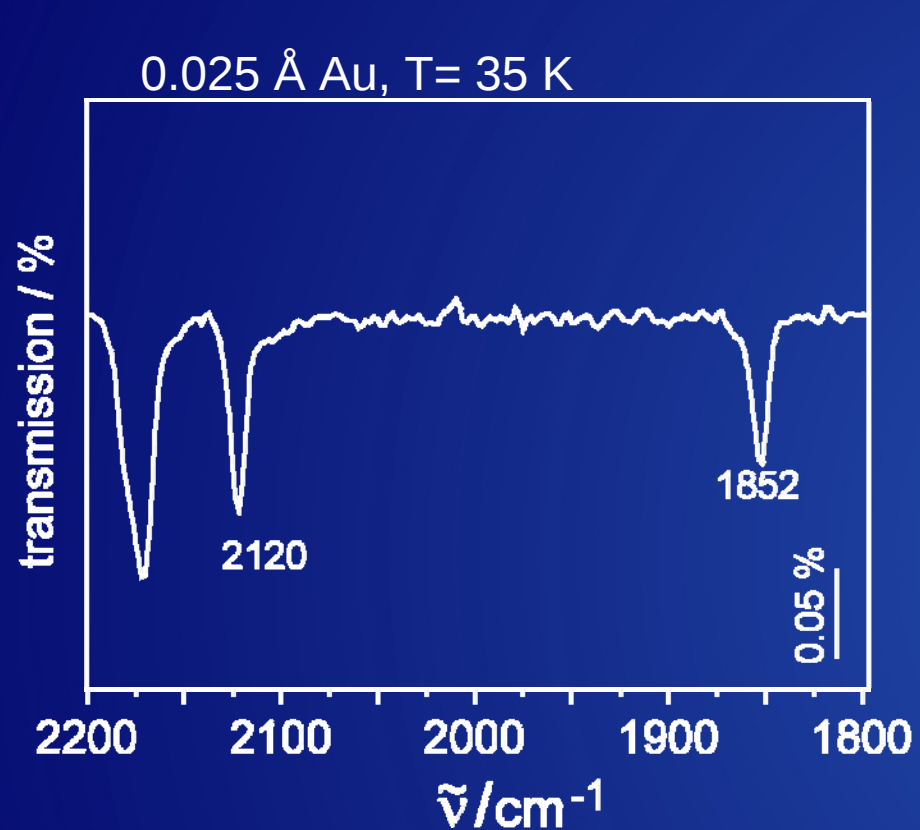
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Au/MgO(001)/Mo(001)

CO IRAS



Max-Planck-Gesellschaft



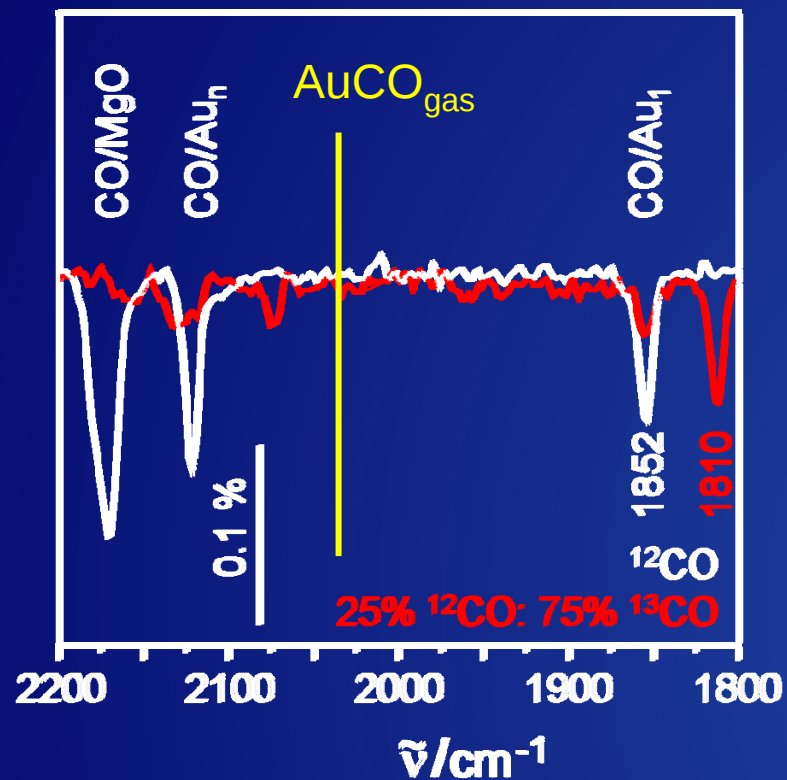
M. Sterrer et al. Angew. Chem. Int. Ed. **45**, 2633 (2006)

Au/MgO(001)/Mo(001)

CO IRAS



Max-Planck-Gesellschaft



red-shift of CO stretching frequency by 187 cm^{-1}
with respect to gas phase AuCO carbonyl

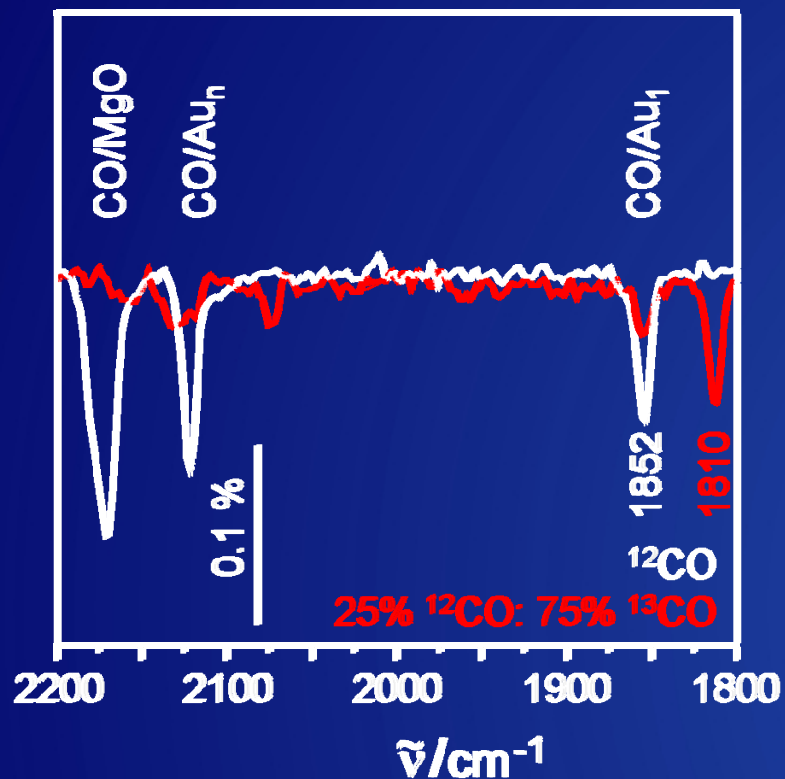
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MgO(001)/Mo(001)

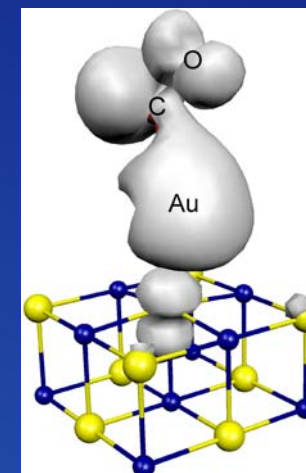
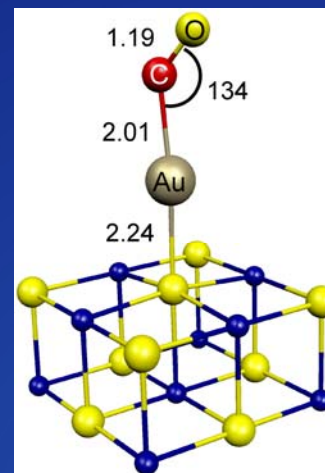
IR of Au atoms



Max-Planck-Gesellschaft



red-shift of CO stretching frequency by
291 cm^{-1} with respect to gas phase



Theory prediction:

- red-shift of 286 cm^{-1} to gas phase
- large electron transfer to Au/CO complex

=> CO induces the charge transfer

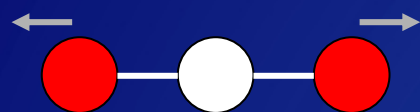
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IRAS

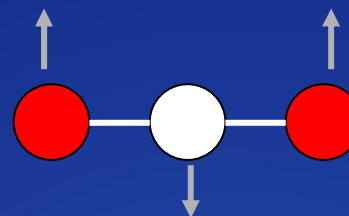
CO₂ on NaCl(001)



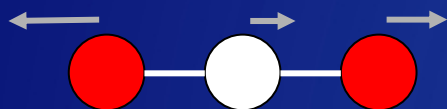
Max-Planck-Gesellschaft



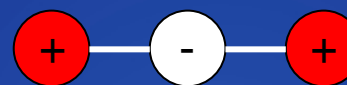
1288 cm⁻¹



667 cm⁻¹



2349 cm⁻¹



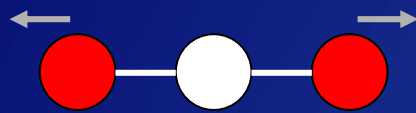
667 cm⁻¹

IRAS

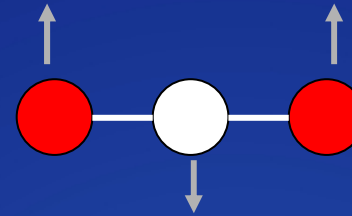
CO₂ on NaCl(001)



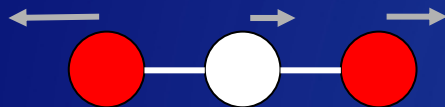
Max-Planck-Gesellschaft



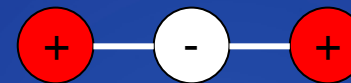
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What is expected?

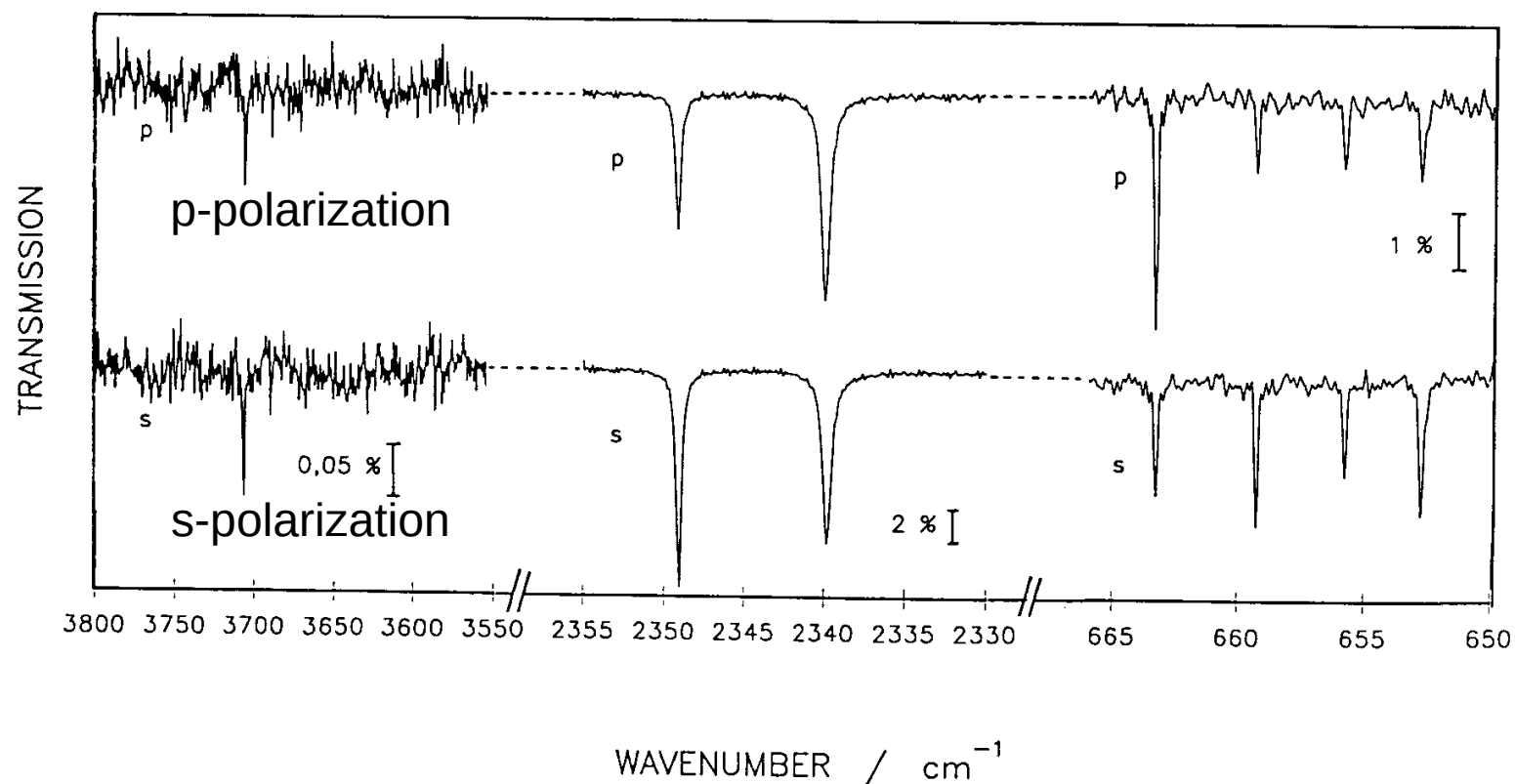


IRAS

CO₂ on NaCl(001)



Max-Planck-Gesellschaft



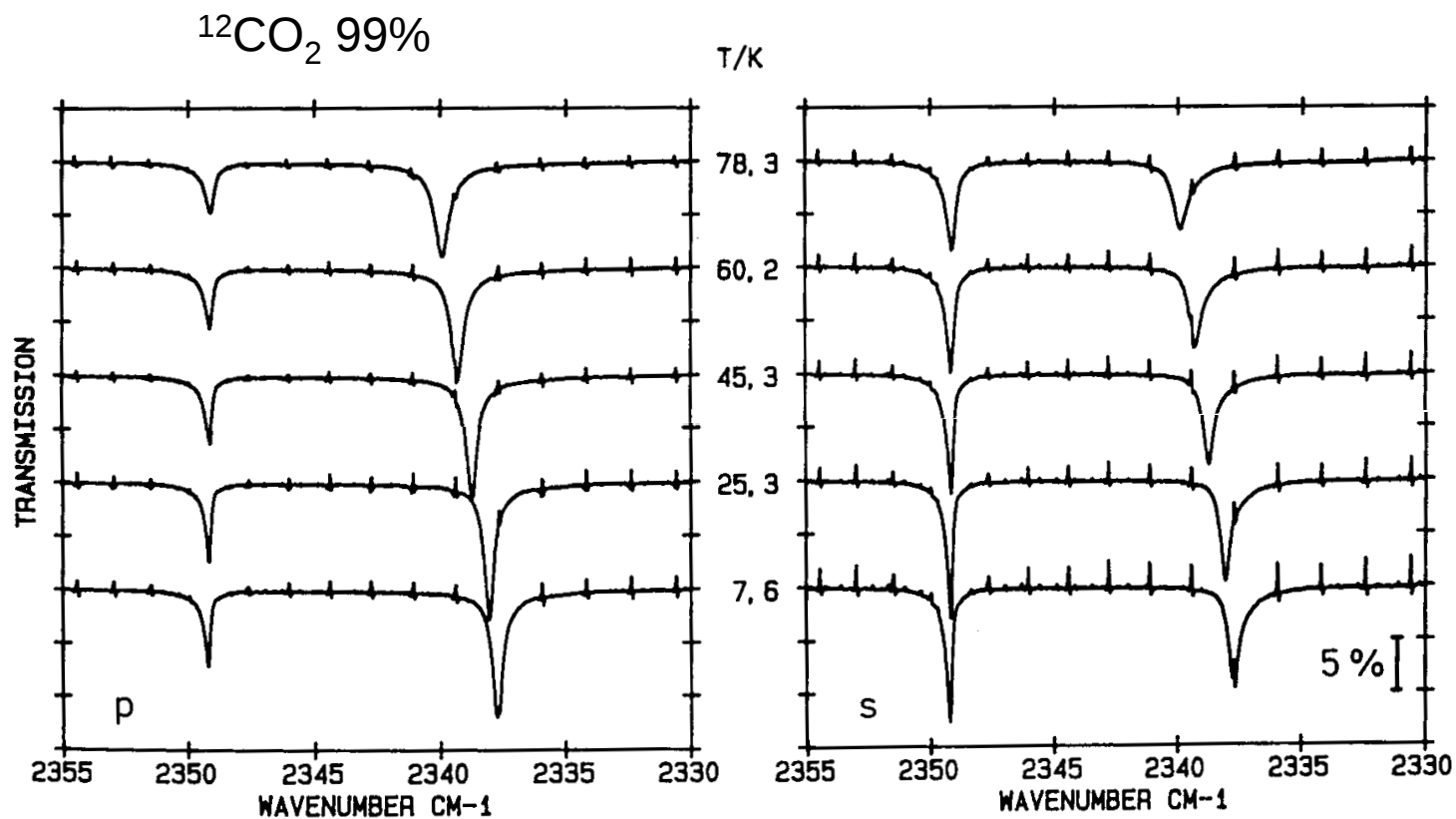
J. Heidberg et al. *J. Electr. Spec. Rel. Phenom.* 64/65, 341 (1993).

IRAS

CO₂ on NaCl(001)



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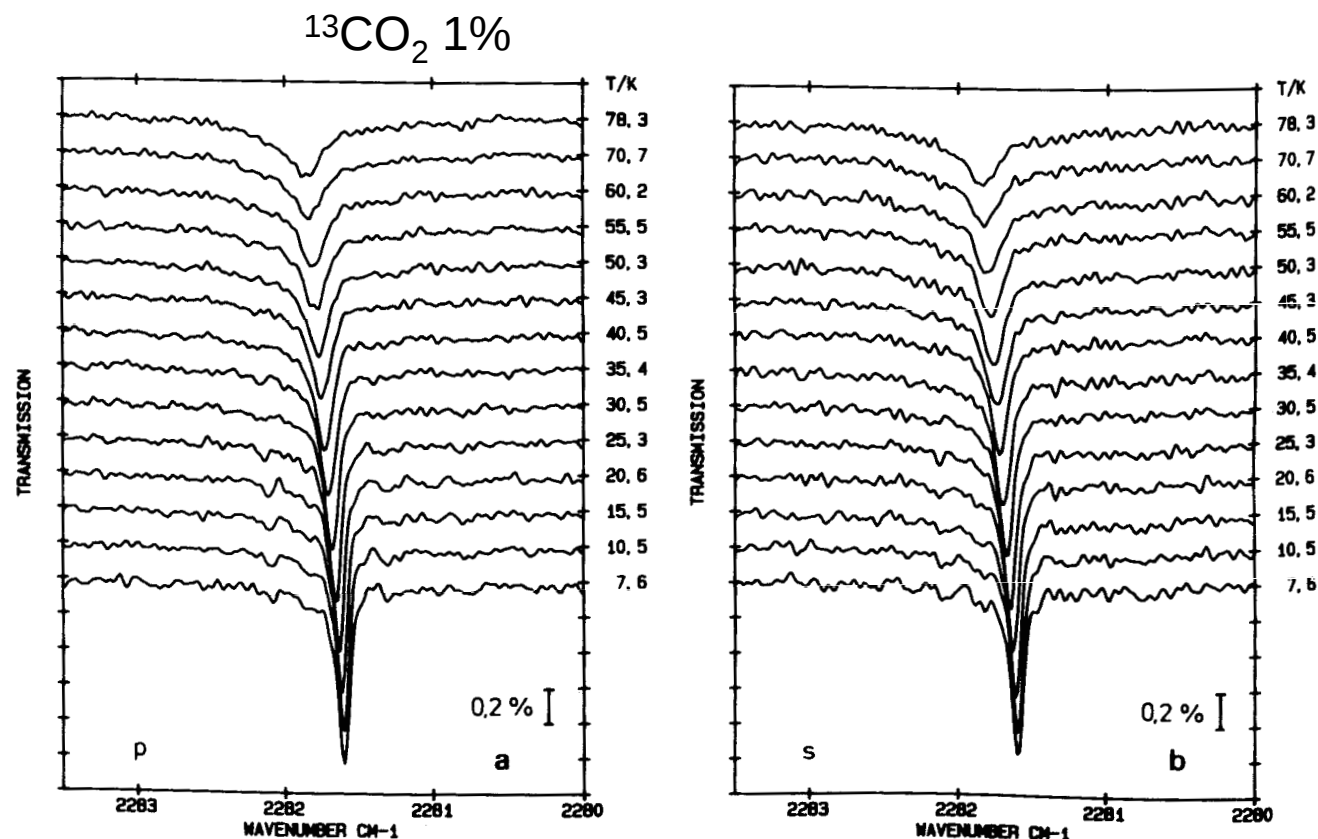
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IRAS

CO₂ on NaCl(001)



Max-Planck-Gesellschaft



- two lines for ¹²CO₂; one line for highly diluted ¹³CO₂
- Peak area constant with T
- same shift for s and p polarization

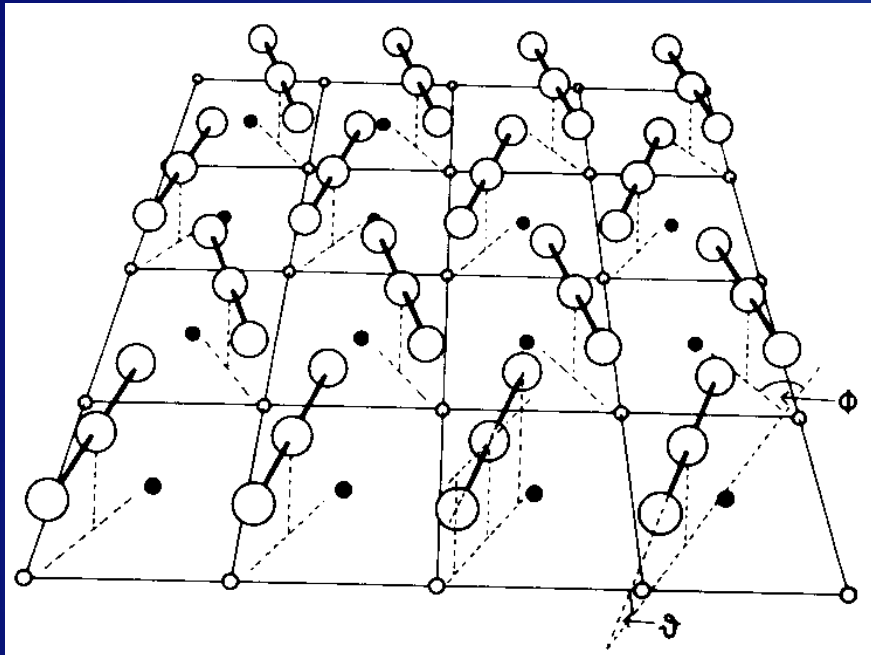
J. Heidberg et al. *J. Electr. Spec. Rel. Phenom.* 64/65, 341 (1993).

IRAS

CO₂ on NaCl(001)



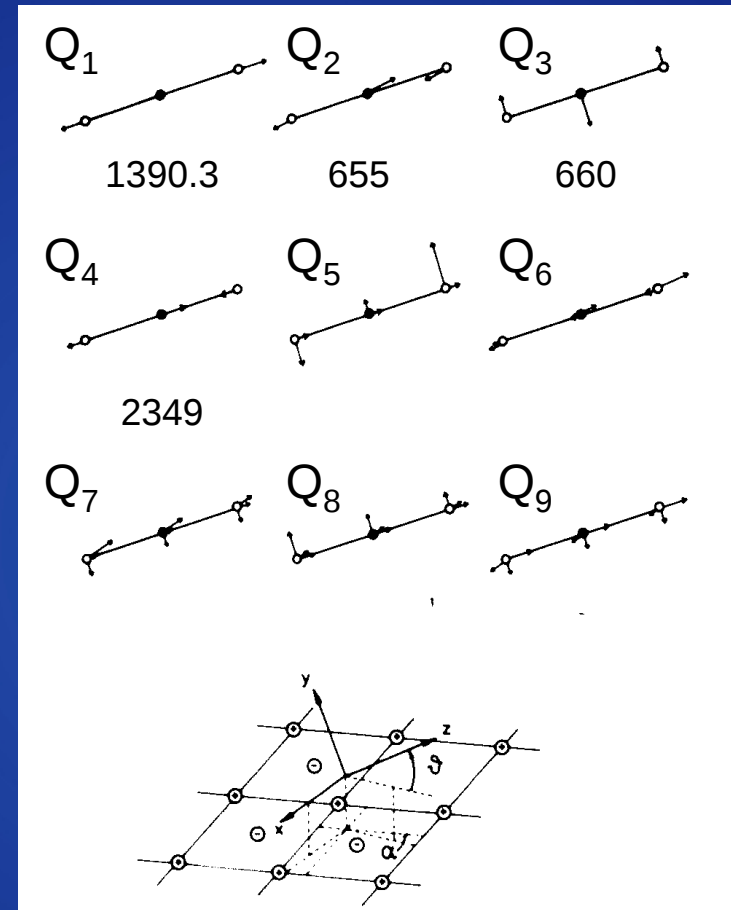
Max-Planck-Gesellschaft



from structural analysis (LEED):

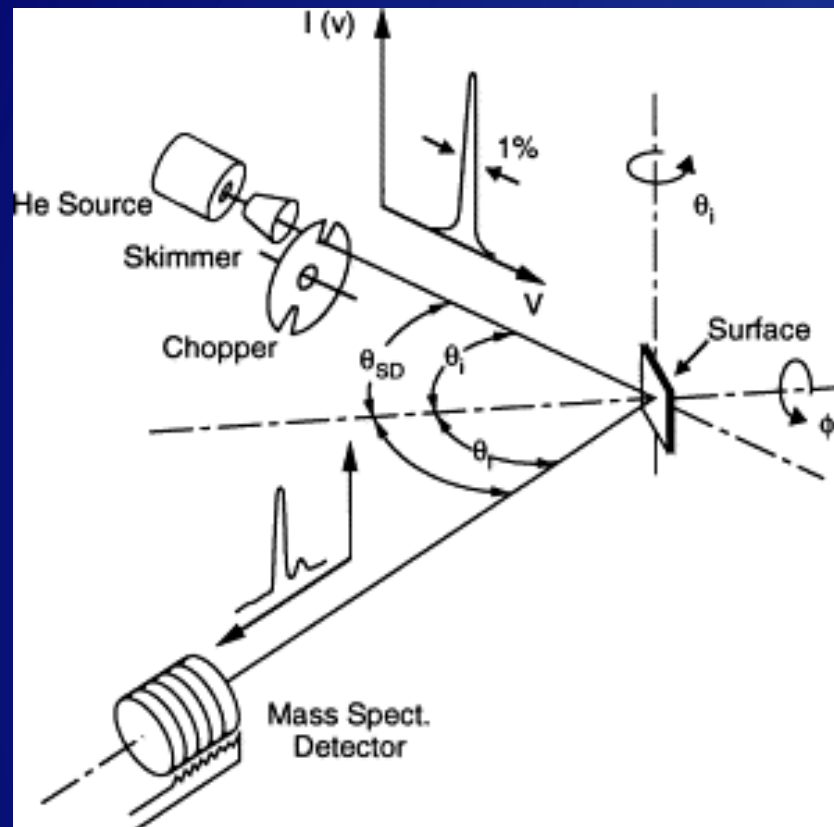
- glide mirror plane (pg symmetry)

J. Heidberg et al. *Surf. Sci.* 269/270, 120 (1992).



Helium atom scattering (HAS)

experimental setup



- energy of the primary beam is ca. 10-40 meV
- inelastic losses by means of a TOF analysis
- Energy resolution: approx 100 μeV
- There is also momentum transfer

For fixed scattering angle of 90° :

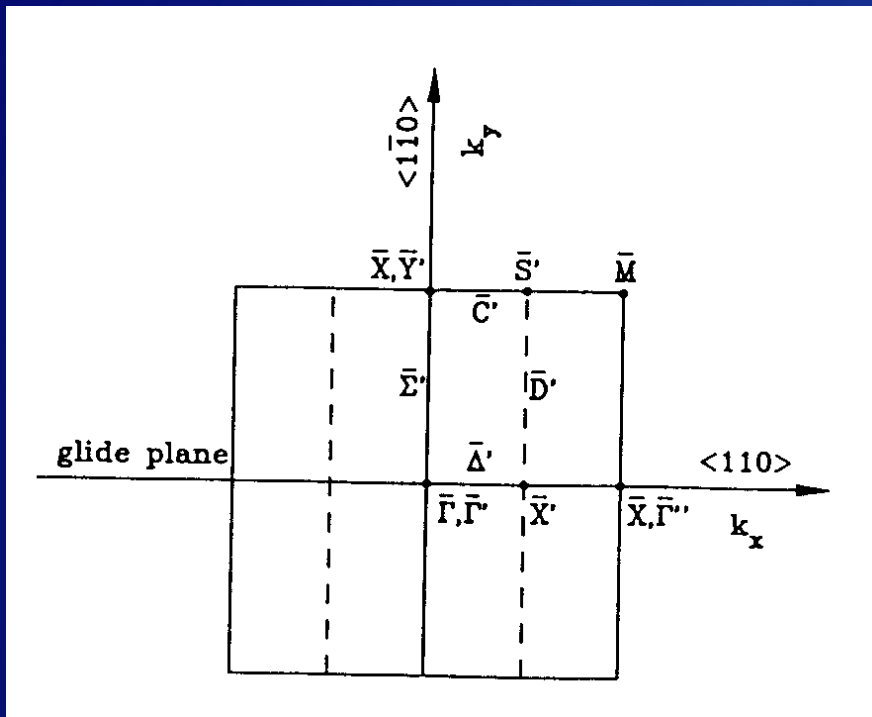
$$\Delta E/E_i = \frac{(\sin(\theta_i) + \Delta k/k_i)^2}{\cos^2(\theta_i)} - 1$$

HAS

CO₂ on NaCl(001)



Max-Planck-Gesellschaft



How many “external” modes to expect?

- two molecules in the unit cell, each five external modes

⇒ 10 modes

additionally: symmetry equivalent islands (along Δ' and Σ')

along Δ' (glide plane) symmetry applies: only five modes

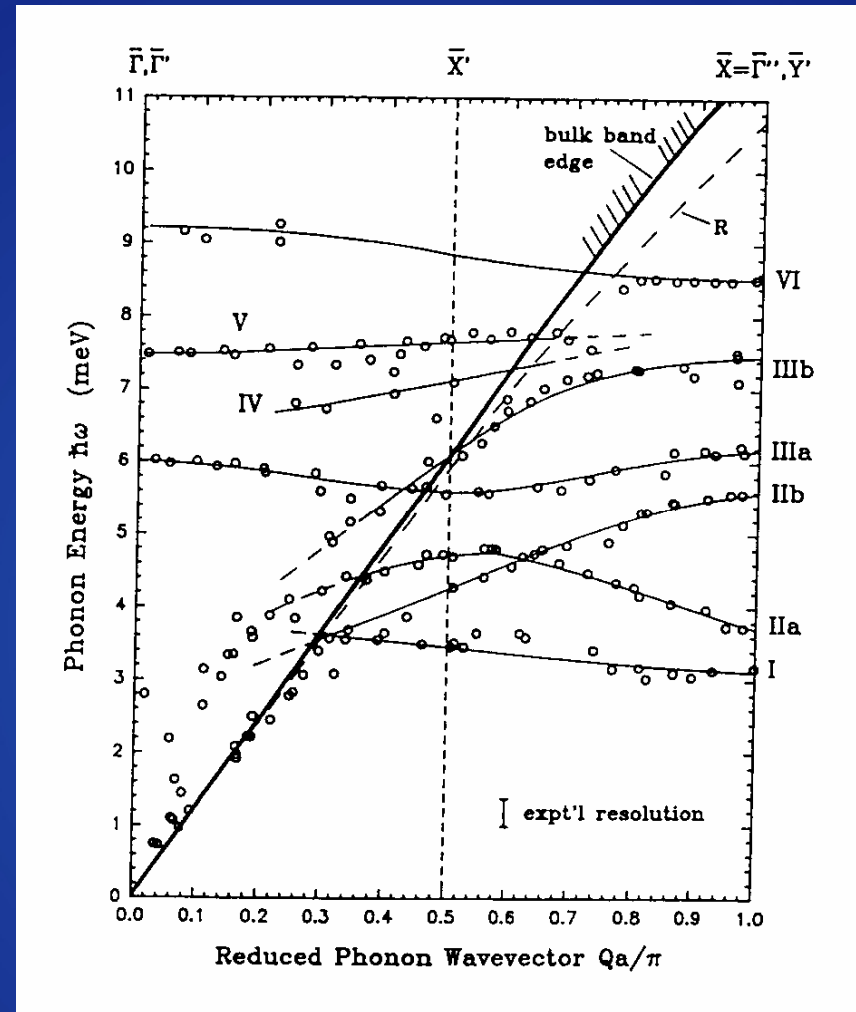
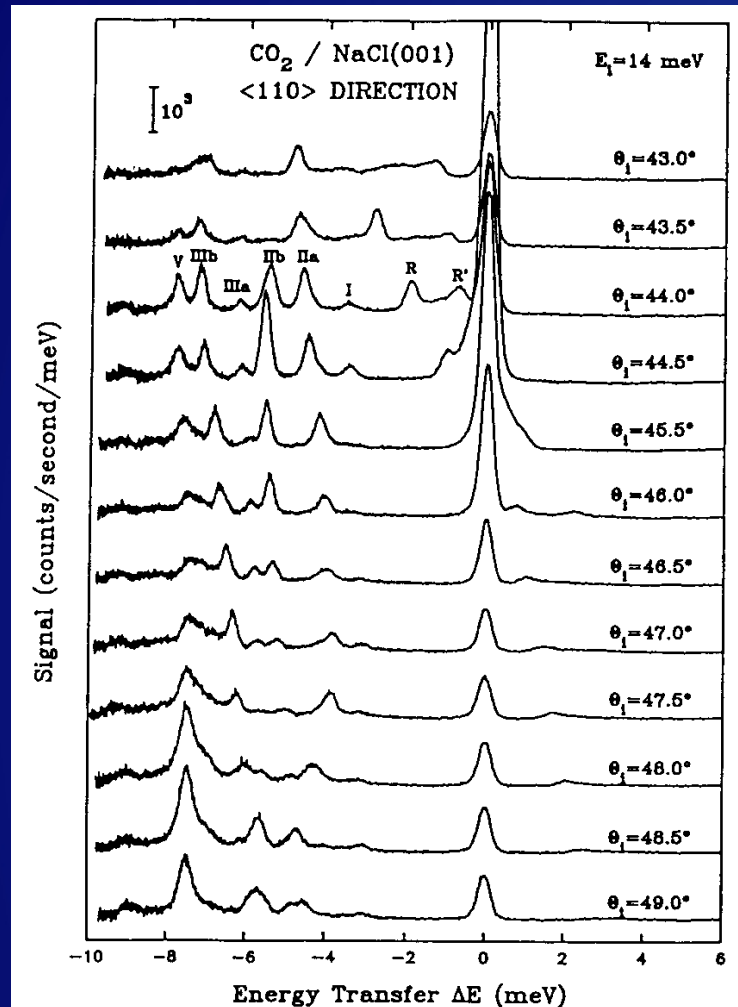
⇒ 15 modes in total

HAS

CO₂ on NaCl(001)



Max-Planck-Gesellschaft



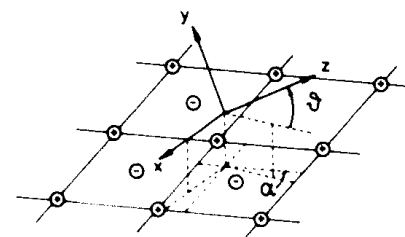
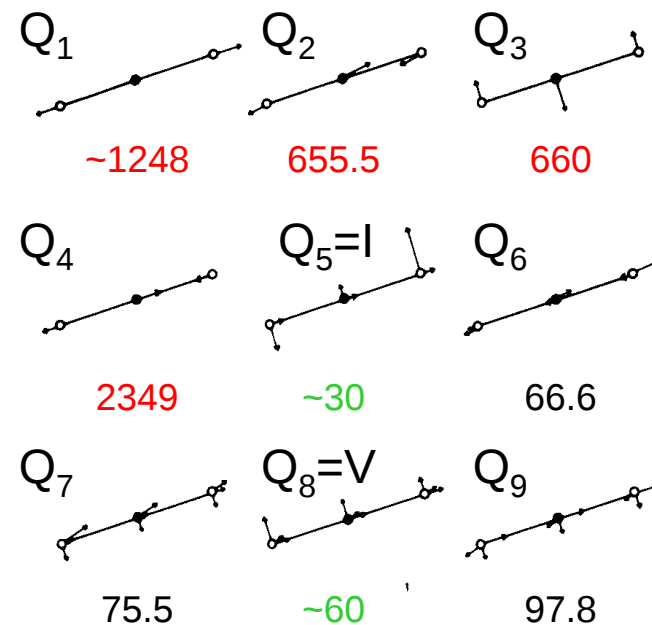
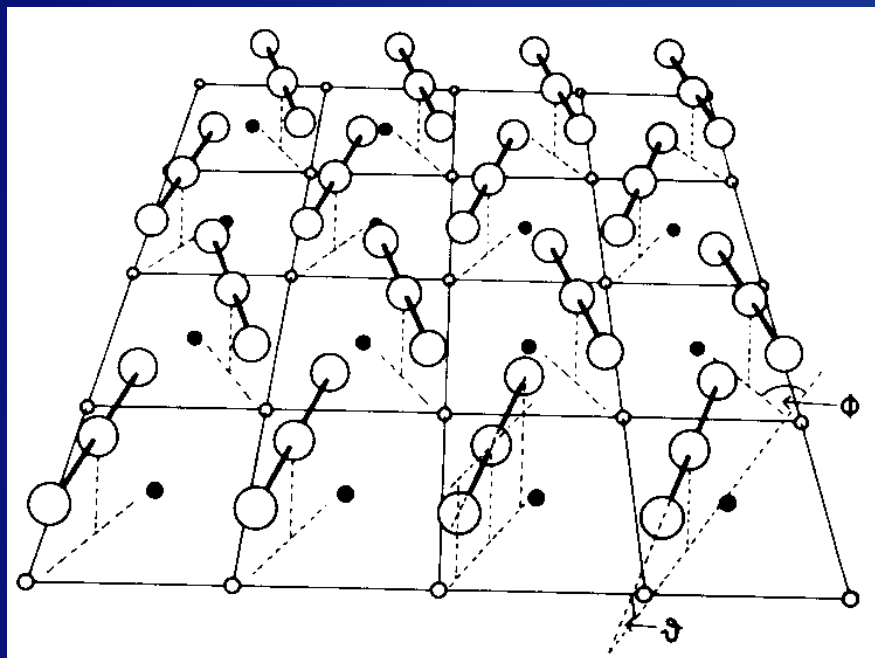
J. Heidberg et al. *J. Electr. Spec. Rel. Phenom.* 64/65, 341 (1993).

IRAS

CO₂ on NaCl(001)



Max-Planck-Gesellschaft



J. Heidberg et al. *Surf. Sci.* 269/270, 120 (1992).

Single molecule vibrations

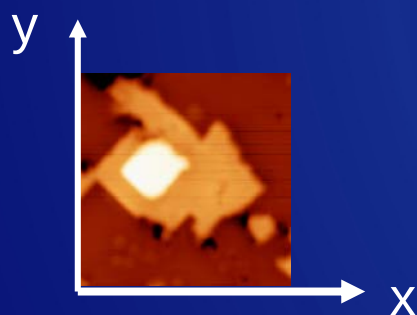
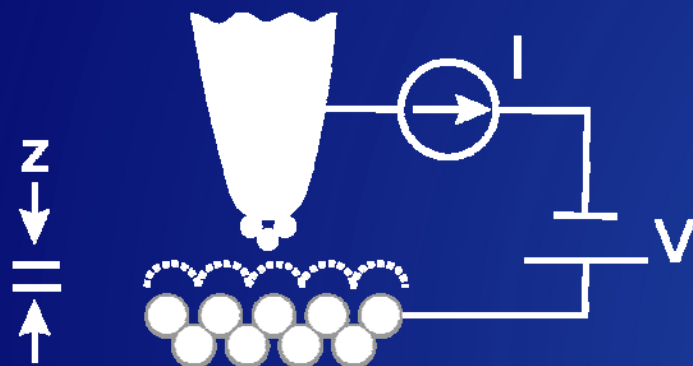
STM/STS



Max-Planck-Gesellschaft

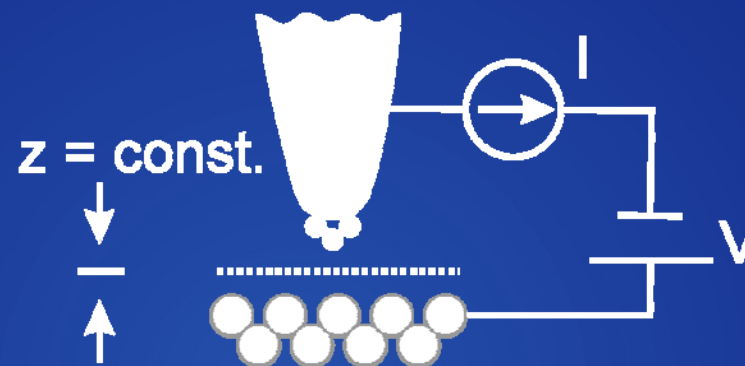
STM

constant current mode

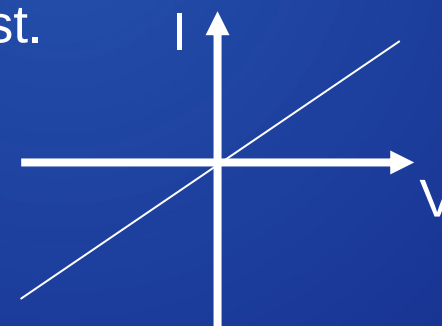


STS

constant height mode



$x, y = \text{const.}$



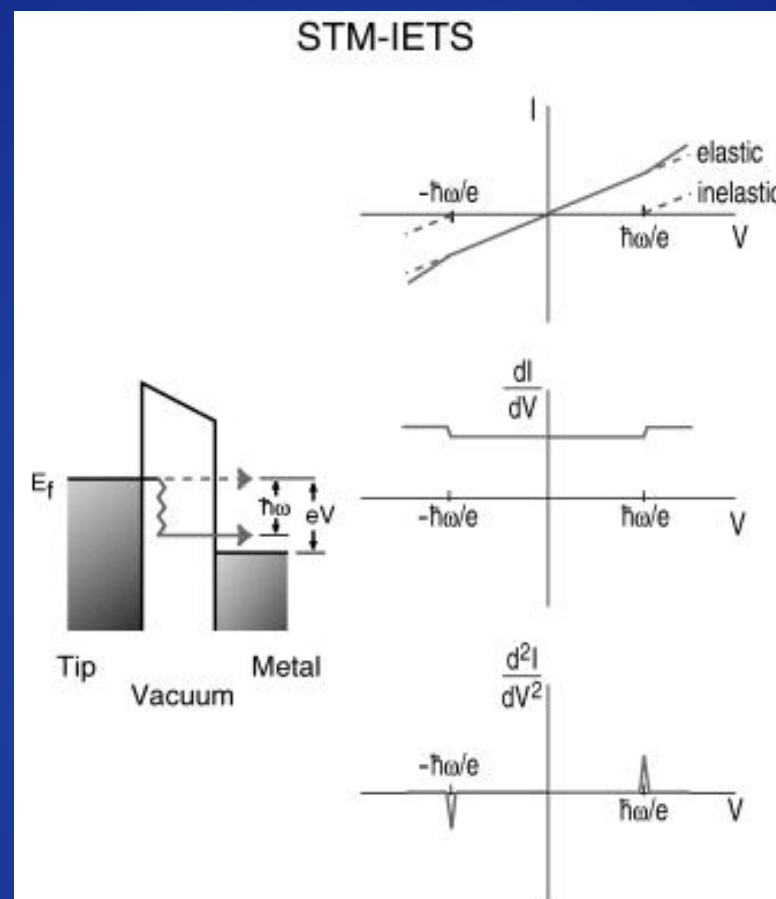
Single molecule vibrations

STM-IETS



Max-Planck-Gesellschaft

threshold spectroscopy



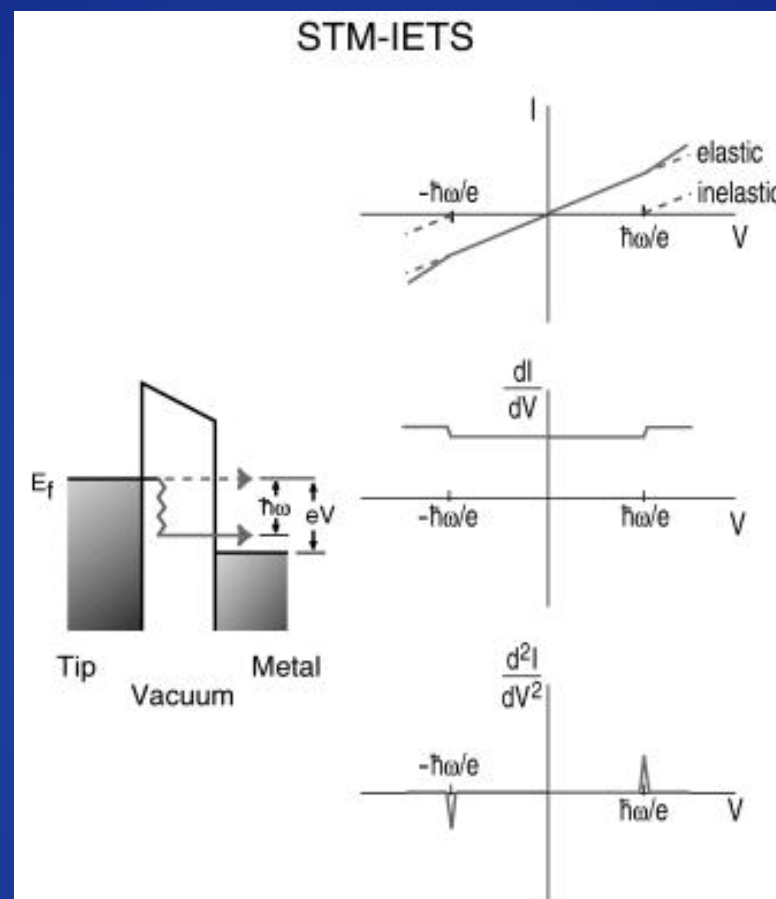
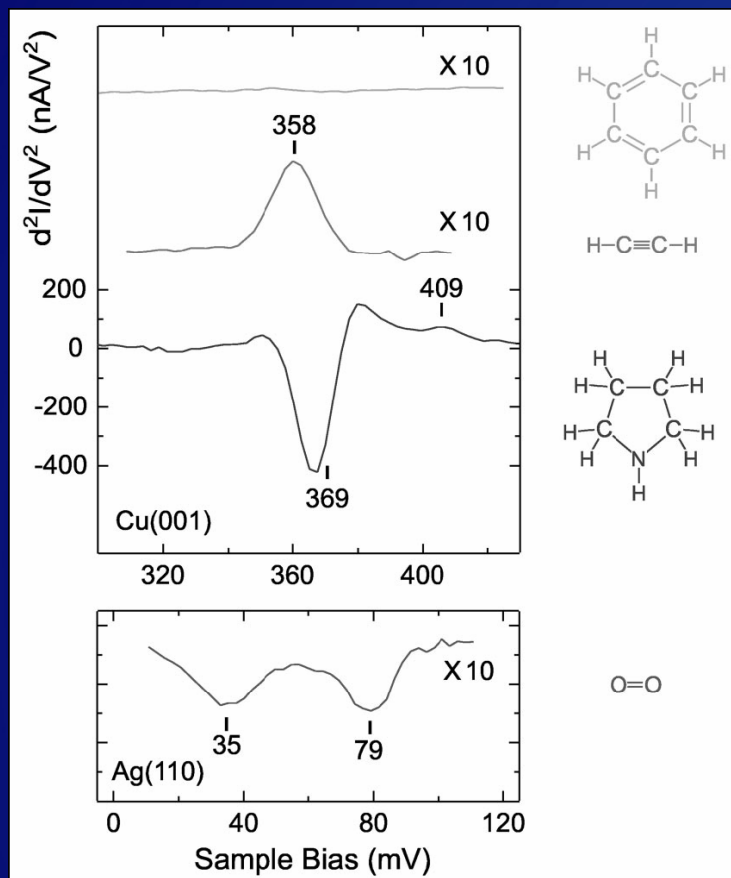
W. Ho J. Chem. Phys. 117, 11033 (2003).

Single molecule vibrations

STM-IETS



Max-Planck-Gesellschaft



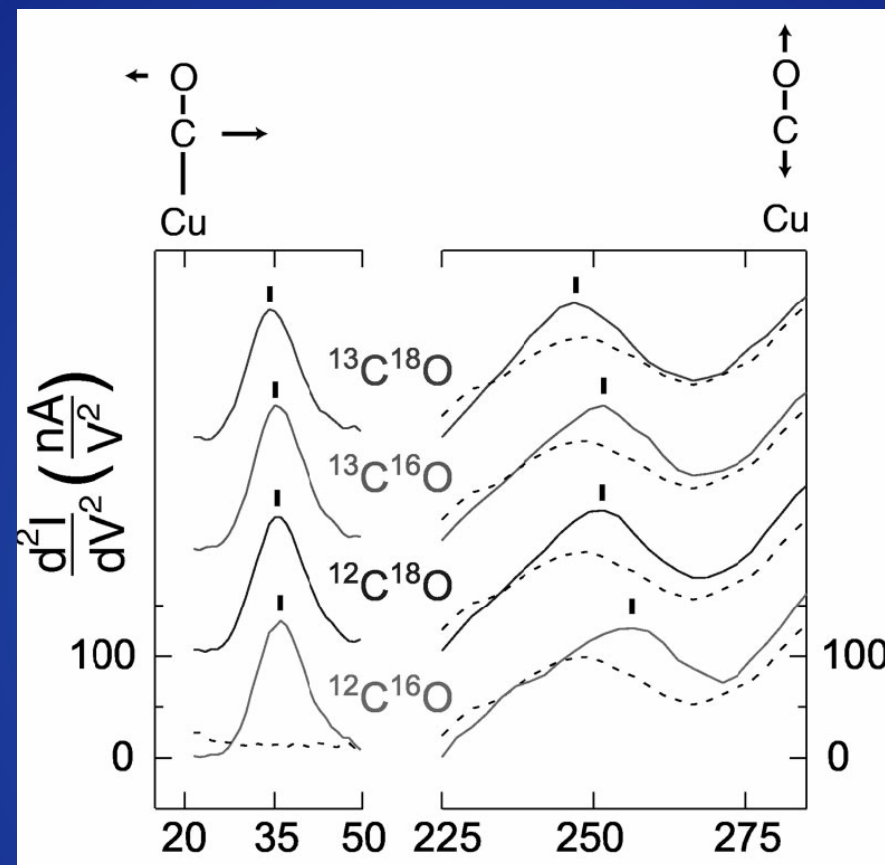
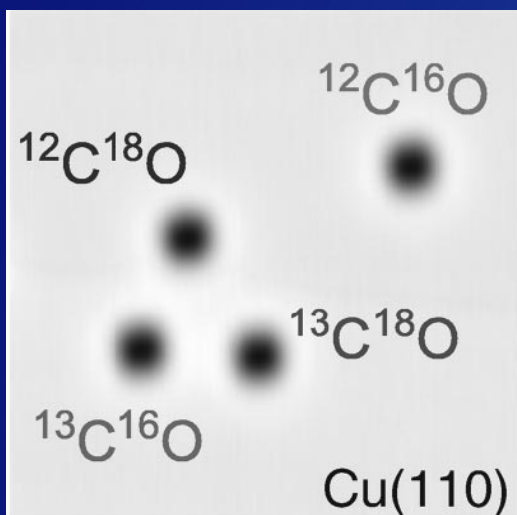
W. Ho J. Chem. Phys. 117, 11033 (2003).

Single molecule vibrations

STM-IETS



Max-Planck-Gesellschaft



W. Ho J. Chem. Phys. 117, 11033 (2003).



Max-Planck-Gesellschaft

vibrational and rotational dynamics of adsorbates on surfaces

Part II

Thomas Risse

Fritz-Haber-Institut der Max-Planck Gesellschaft,
Abteilung Chemische Physik
Faradayweg 4-6
14195 Berlin

Outline



Max-Planck-Gesellschaft

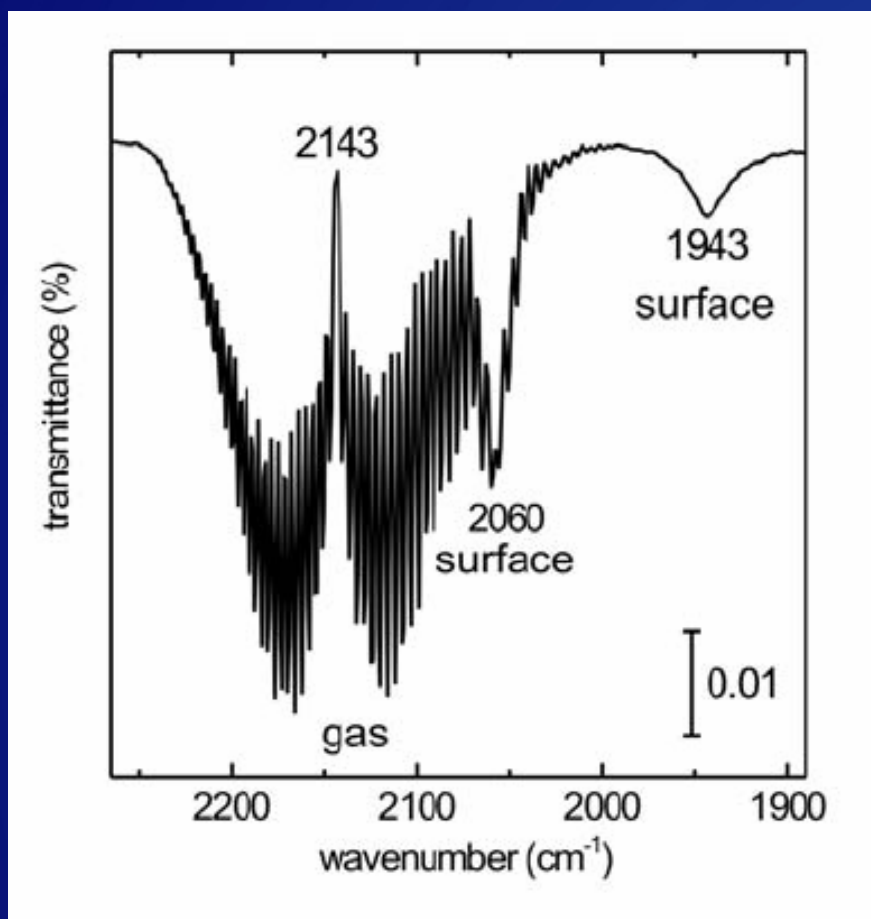
- Introduction
- electron energy loss spectroscopy (EELS)
- infrared spectroscopy
- helium atom scattering (HAS)
- PMIRAS
- SFG
- Raman
- scanning tunneling spectroscopy
- electron spin resonance

IRAS

Ambient pressures



Max-Planck-Gesellschaft



50 mbar CO on
Pd(111) at 300 K

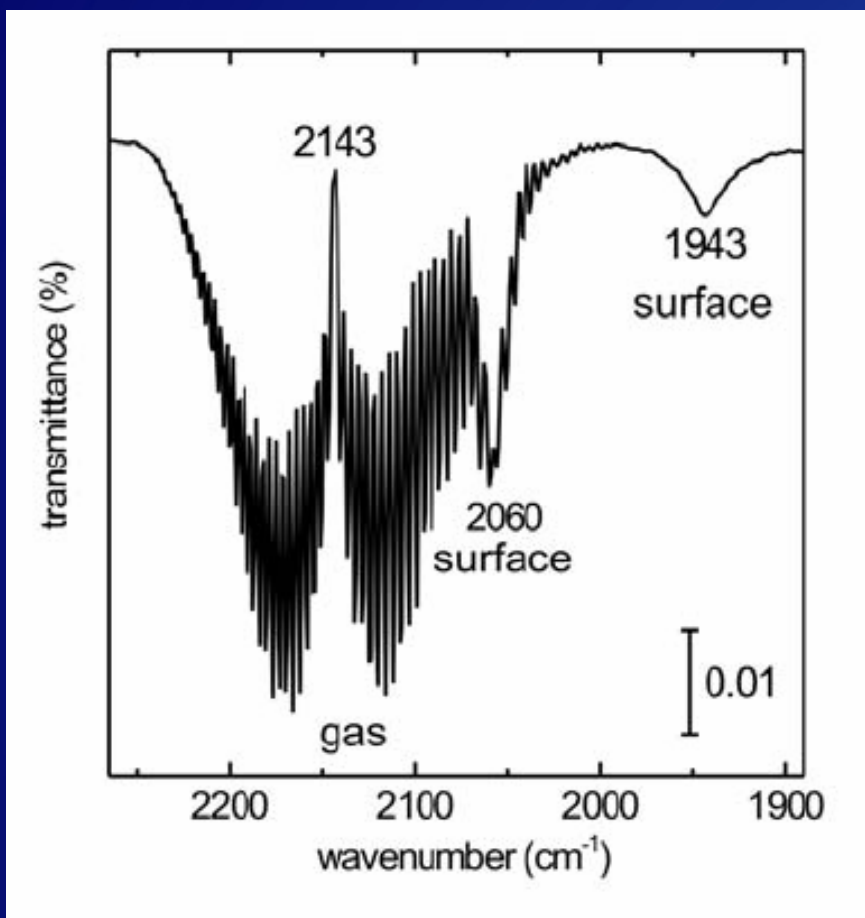
M. Borasio, PhD-thesis, FU-Berlin (2006)

IRAS

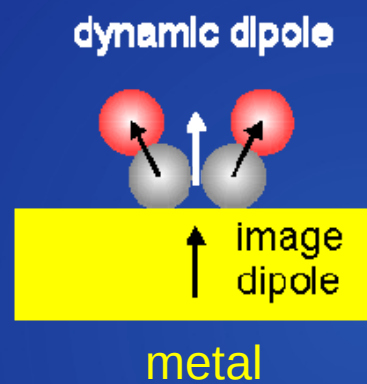
Ambient pressures



Max-Planck-Gesellschaft

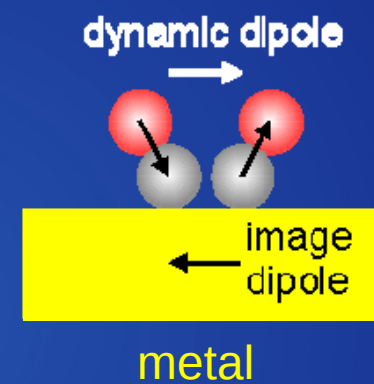


M. Borasio, PhD-thesis, FU-Berlin (2006)



$$\mu_{\text{dyndip}} > 0$$

IR active



$$\mu_{\text{dyndip}} = 0$$

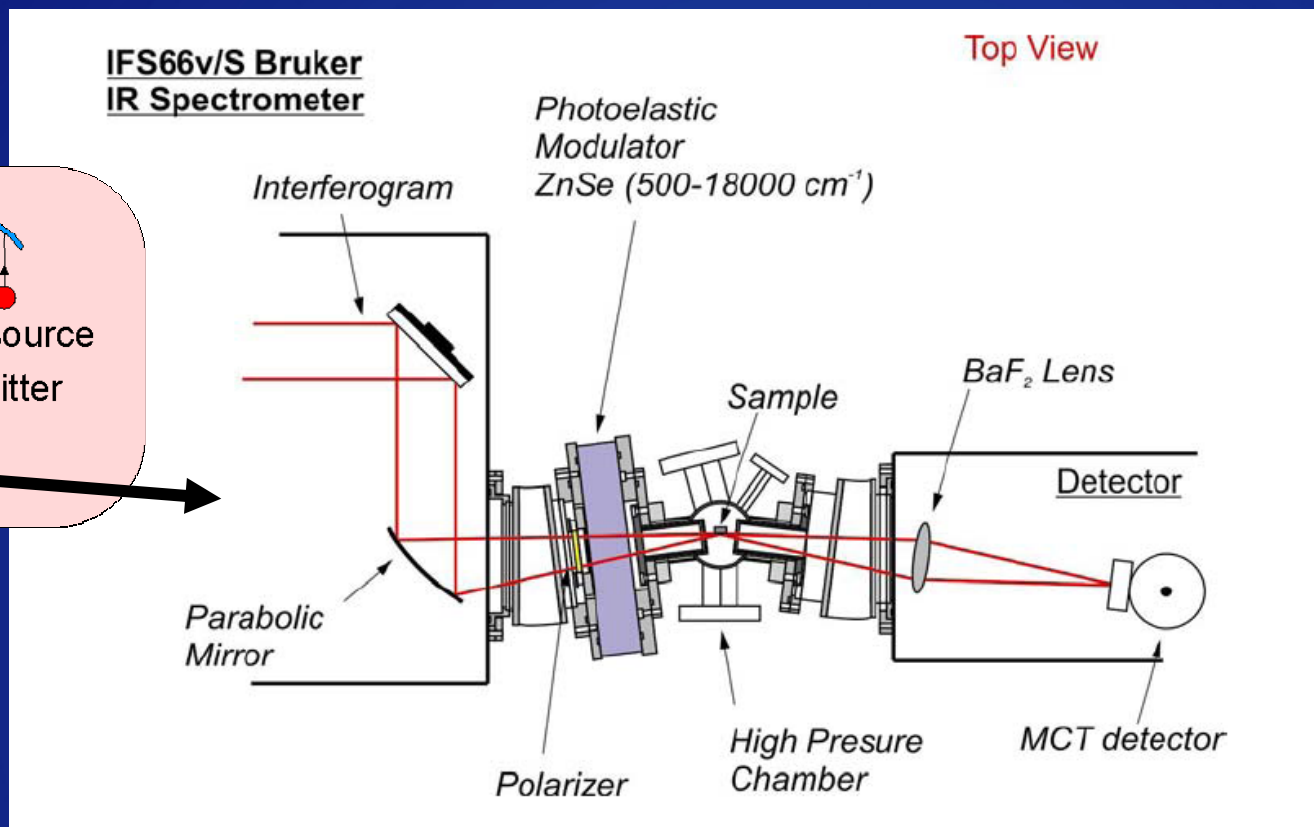
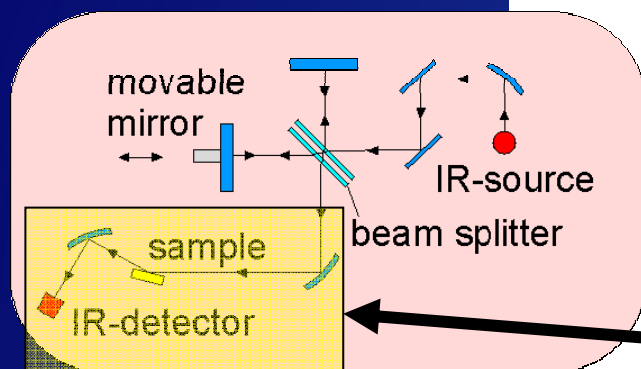
IR inactive

IRAS

PMIRAS



Max-Planck-Gesellschaft



M. Borasio, PhD-thesis, FU-Berlin (2006)

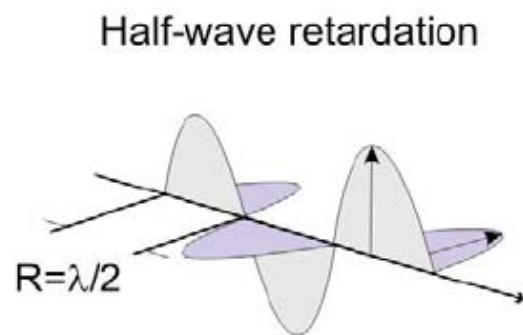
IRAS

PMIRAS

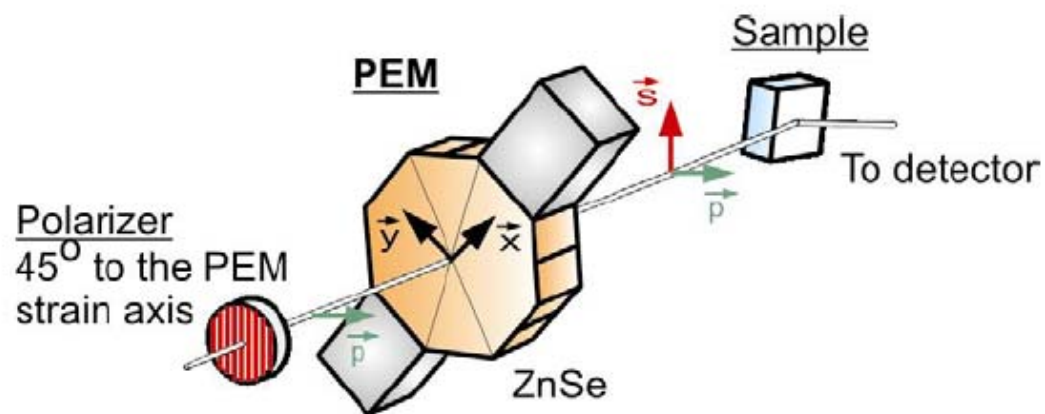


Max-Planck-Gesellschaft

central part: photoacoustic modulator



(a)



(b)

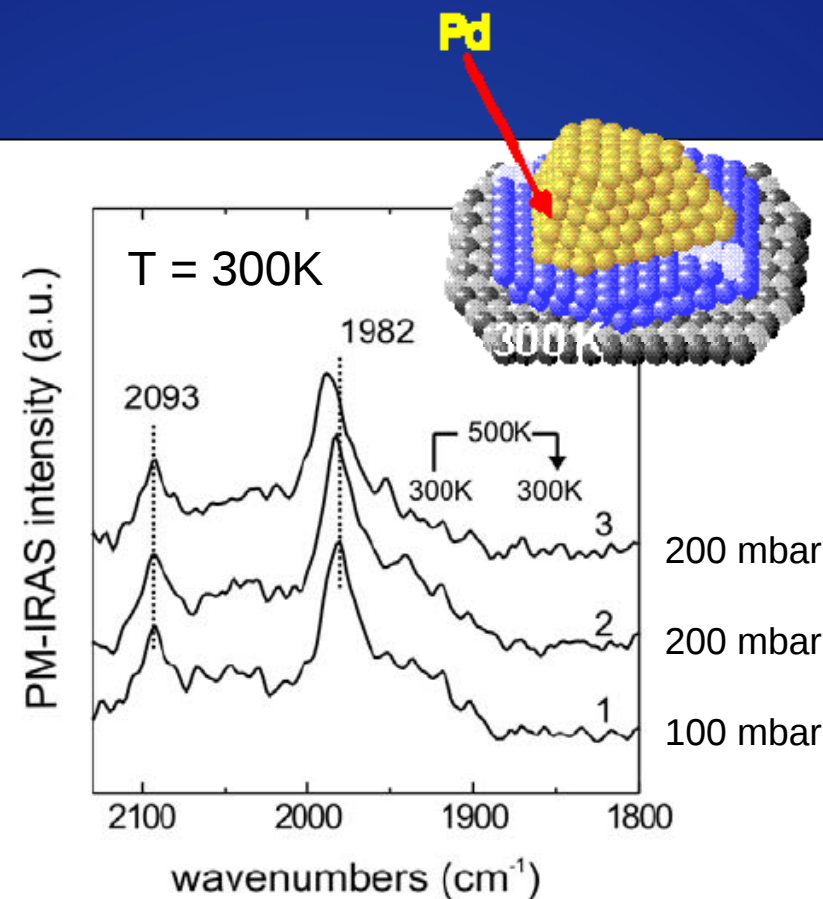
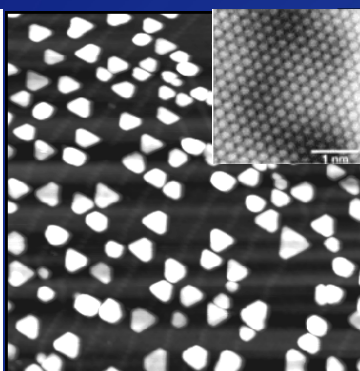
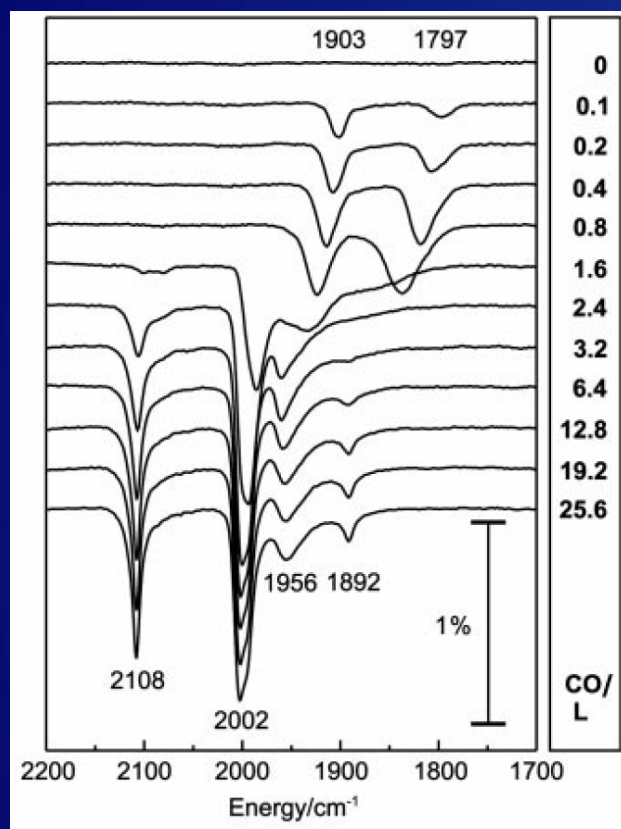
M. Borasio, PhD-thesis, FU-Berlin (2006)

PMIRAS

CO/Pd/Al₂O₃/NiAl(110)



Max-Planck-Gesellschaft



M. Frank and M. Bäumer, PCCP 2, 3723 (2000).

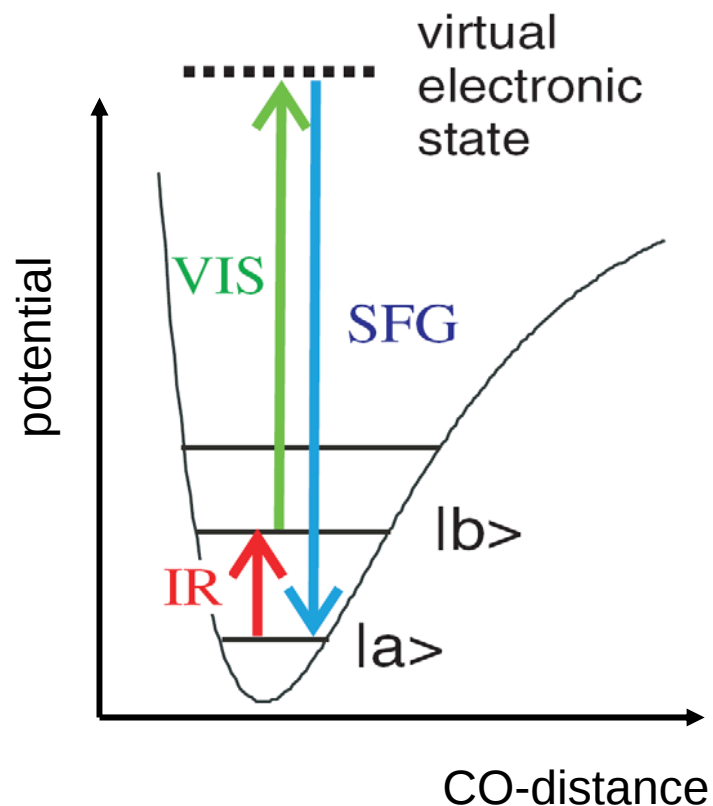
M. Borasio, PhD-thesis, FU-Berlin (2006)

Vibrational spectroscopy

Sum frequency generation



Max-Planck-Gesellschaft



Second order non optical linear process

$$\mathbf{P}_{\omega_{\text{SFG}}}^{(2)} = \chi_{\omega_{\text{SFG}}}^{(2)} \mathbf{E}_{\omega_{\text{IR}}} \mathbf{E}_{\omega_{\text{VIS}}}$$

Surface sensitive method

SFG is symmetry forbidden in isotropic (bulk) media (dipole approximation)

Vibrational spectroscopy

Sum frequency generation

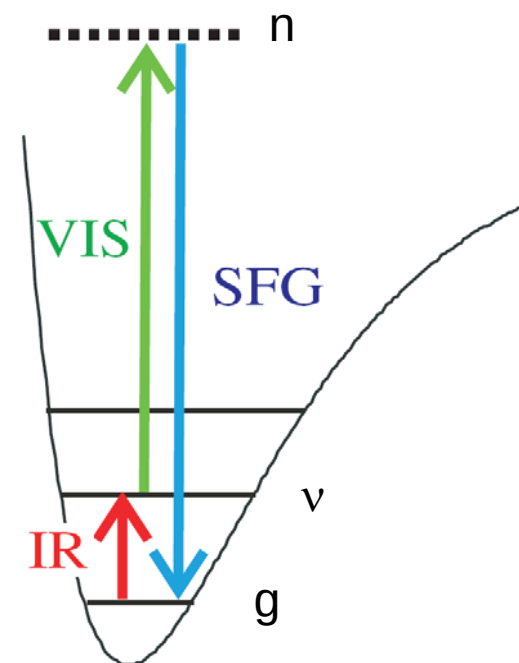


Max-Planck-Gesellschaft

$$\mathbf{P}_{\omega_{\text{SFG}}}^{(2)} = \chi_{\omega_{\text{SFG}}}^{(2)} \mathbf{E}_{\omega_{\text{IR}}} \mathbf{E}_{\omega_{\text{VIS}}}$$

$$\chi_{ijk}^{(2)}(-\omega_{SF}; \omega_{vis}, \omega_{IR}) = \frac{N}{2\epsilon_0 \hbar^2} \sum_n \rho_{gg}^0 \times \left\{ \frac{\mu_{gn}^i \mu_{nv}^j \mu_{vg}^k}{[(\omega_{ng} - \omega_{SF}) - i\gamma_{ng}][(\omega_{vg} - \omega_{IR}) - i\gamma_{vg}]} + \frac{\mu_{gn}^j \mu_{nv}^i \mu_{vg}^k}{[(\omega_{nv} + \omega_{SF}) + i\gamma_{nv}][(\omega_{vg} - \omega_{IR}) - i\gamma_{vg}]} \right\}$$

- molecules in the vibronic ground state g
- IR light in resonance with a vibration



R.W. Boyd *Nonlinear Optics*, Academic Press, New York, 1992.

Vibrational spectroscopy

Sum frequency generation



Max-Planck-Gesellschaft

$$\mathbf{P}_{\omega_{\text{SFG}}}^{(2)} = \chi_{\omega_{\text{SFG}}}^{(2)} \mathbf{E}_{\omega_{\text{IR}}} \mathbf{E}_{\omega_{\text{VIS}}}$$

$$\chi_{ijk}^{(2)}(-\omega_{\text{SF}}; \omega_{\text{vis}}, \omega_{\text{IR}}) = \frac{N \rho_{gg}^0}{2 \epsilon_0 \hbar^2} \frac{\mu_{\nu g}^k}{[(\omega_{\nu g} - \omega_{\text{IR}}) - i \gamma_{\nu g}]} \times \sum_n \left\{ \frac{\mu_{gn}^i \mu_{n\nu}^j}{[(\omega_{ng} - \omega_{\text{SF}}) - i \gamma_{ng}]} + \frac{\mu_{gn}^j \mu_{n\nu}^i}{[(\omega_{n\nu} + \omega_{\text{SF}}) + i \gamma_{n\nu}]} \right\}.$$

$$\chi_{ijk}^{(2)}(-\omega_{\text{SF}}; \omega_{\text{vis}}, \omega_{\text{IR}}) = \frac{N \rho_{gg}^0}{2 \epsilon_0 \hbar} \frac{\mu_{\nu g}^k}{[(\omega_{\nu g} - \omega_{\text{IR}}) - i \gamma_{\nu g}]} \alpha_{g\nu}^{ij}(-\omega_{\text{SF}}; \omega_{\text{vis}})$$

$$\alpha_{g\nu}^{ij}(-\omega_{\text{SF}}; \omega_{\text{vis}}) = \frac{1}{\hbar} \sum_n \left\{ \frac{\mu_{gn}^i \mu_{n\nu}^j}{(\omega_{ng} - \omega_{\text{SF}})} + \frac{\mu_{gn}^j \mu_{n\nu}^i}{(\omega_{ng} + \omega_{\text{vis}})} \right\}.$$

$\alpha_{g\nu}$: 1. order hyperpolarizability (anti Stokes Raman intensity)

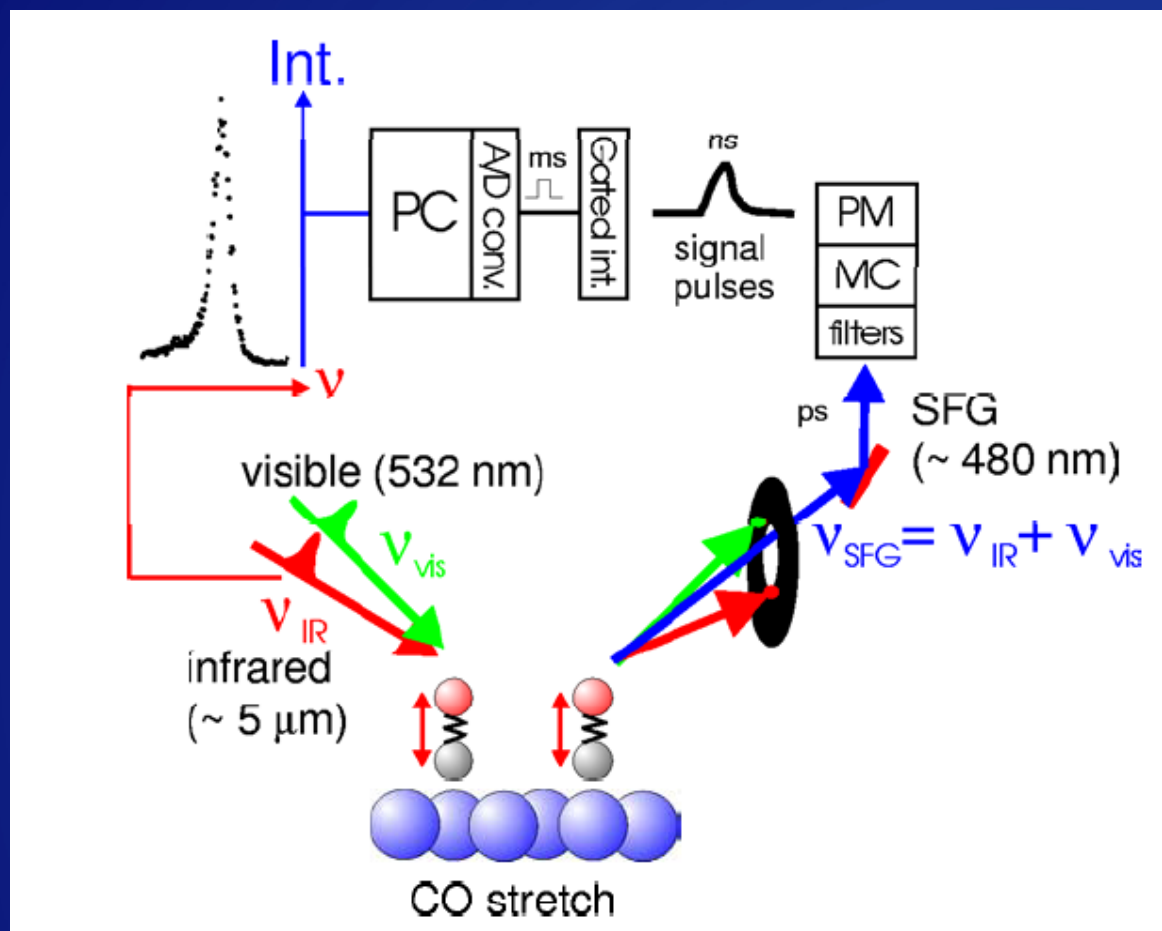
R.W. Boyd *Nonlinear Optics*, Academic Press, New York, 1992.

Vibrational spectroscopy

Sum frequency generation



Max-Planck-Gesellschaft



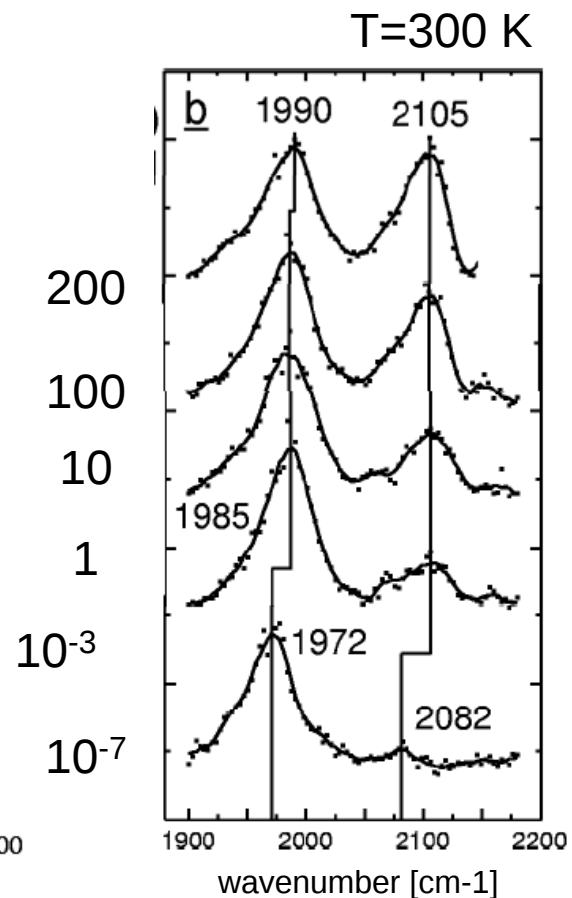
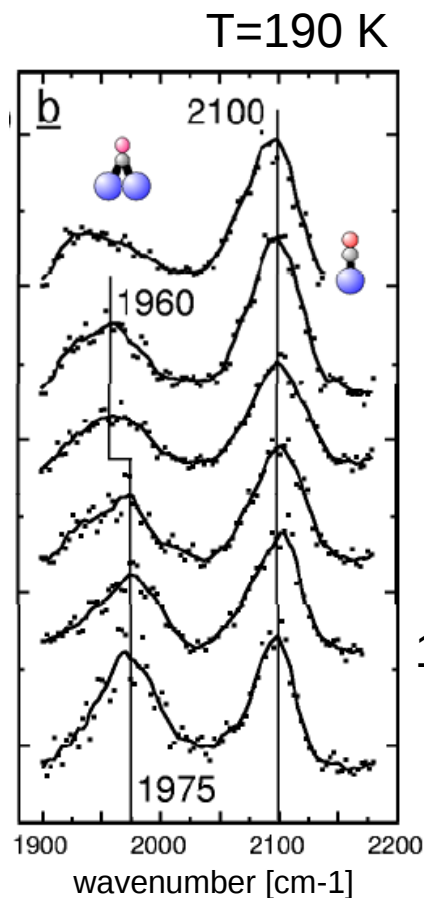
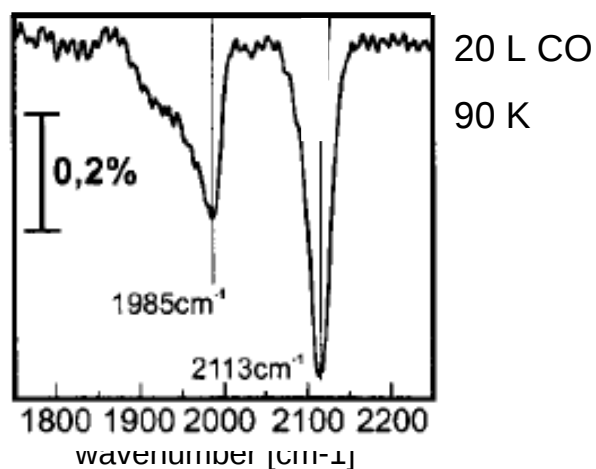
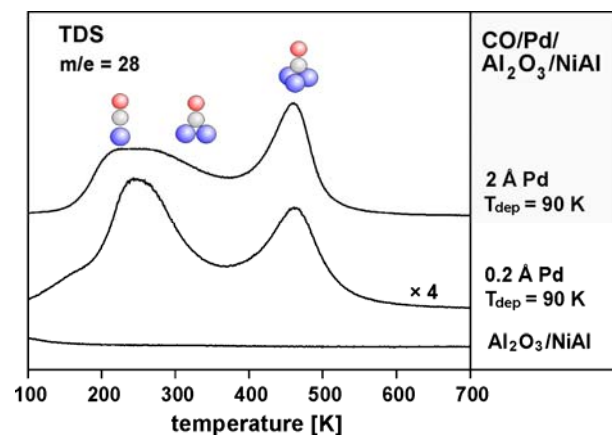
H. Unterhalt, PhD-thesis, FU-Berlin (2002)

SFG

CO/Pd/Al₂O₃/NiAl(110)



Max-Planck-Gesellschaft



K. Wolter et al. Surf. Sci. **399**, 190 (1998).

H. Unterhalt et al. J. Phys. Chem. B **106**, 356 (2002).

M. Bäumer et al. Ber. Bunsenges. Phys. Chem. **99**, 1381 (1995).

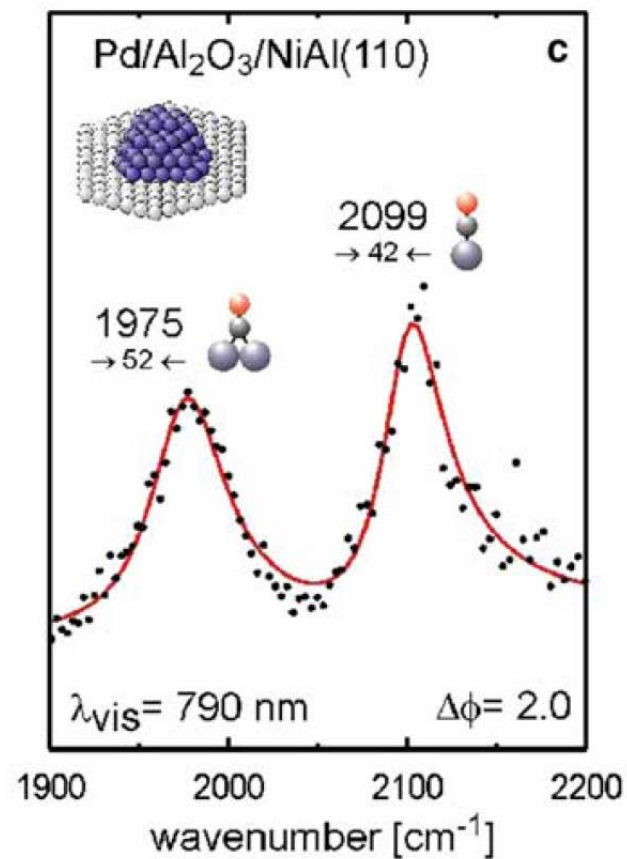
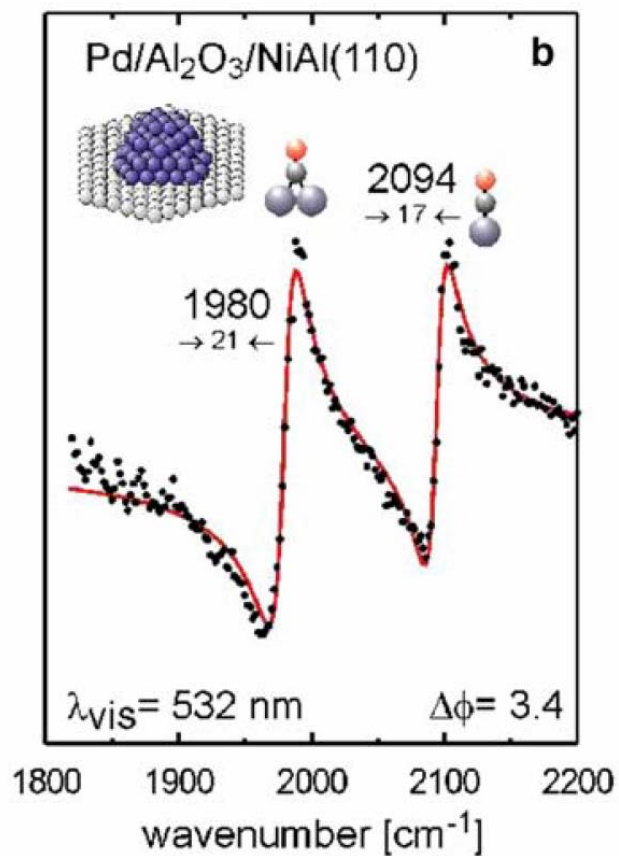
SFG

CO/Pd/Al₂O₃/NiAl(110)



Max-Planck-Gesellschaft

T=190 K



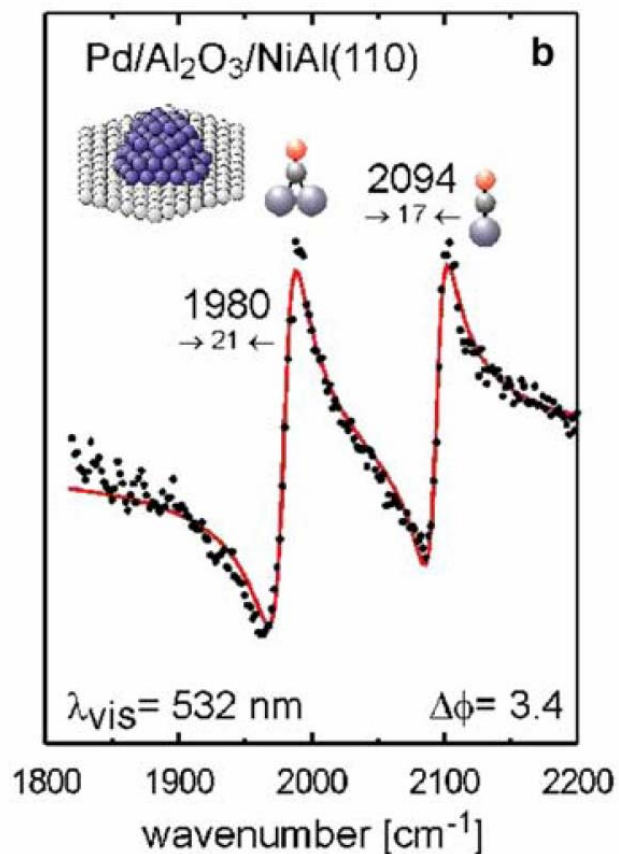
M. Morkel et al. Surf. Sci. **586**, 146 (2005).

SFG

CO/Pd/Al₂O₃/NiAl(110)



Max-Planck-Gesellschaft



$$\chi_s^{(2)} = \chi_{\text{NR}}^{(2)} + \chi_{\text{R}}^{(2)} = A_{\text{NR}} \cdot e^{i\phi_0} + \sum_q \frac{A_q \cdot e^{i\phi_q}}{(\omega_q - \omega_{\text{IR}}) - i\Gamma_q}$$

In case the phase of the non resonant background is constant I_{SFG} does only depend on the phase difference $\Delta\phi_{0q}$:

$$I_{\text{SFG}} \propto \left| A_{\text{NR}} + \sum_q \frac{A_q \cdot e^{i\Delta\phi_{0q}}}{(\omega_q - \omega_{\text{IR}}) - i\Gamma_q} \right|^2 \cdot I_{\text{vis}} \cdot I_{\text{IR}}$$

M. Morkel et al. Surf. Sci. **586**, 146 (2005).

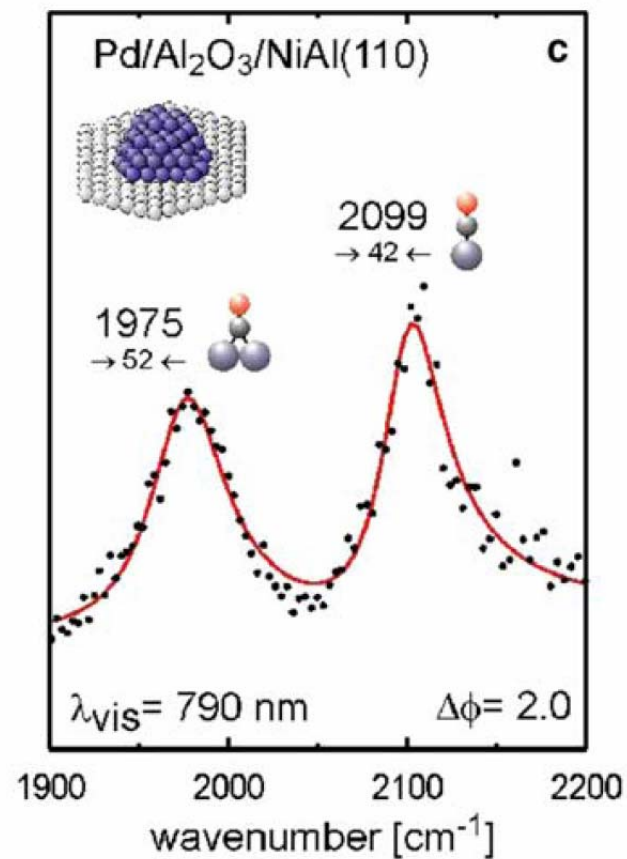
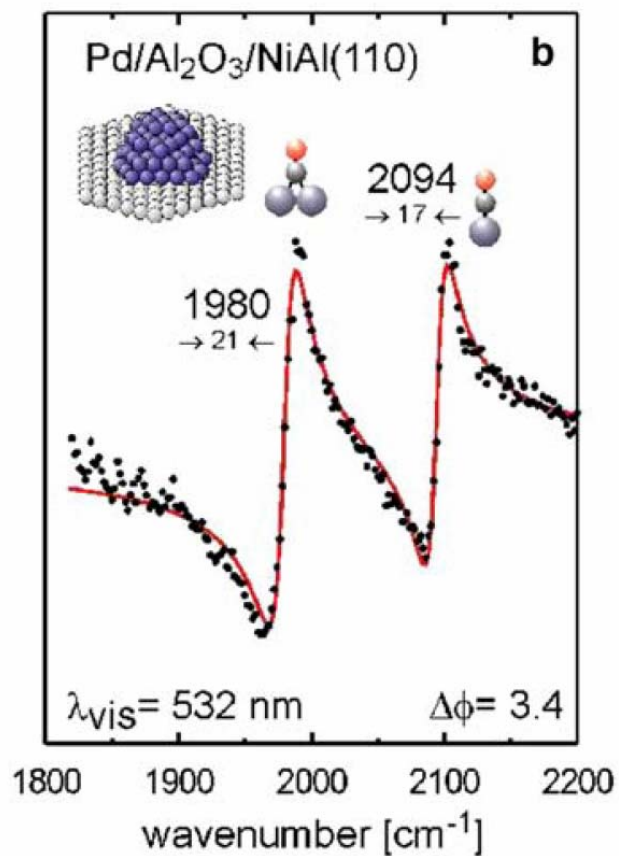
SFG

CO/Pd/Al₂O₃/NiAl(110)



Max-Planck-Gesellschaft

T=190 K



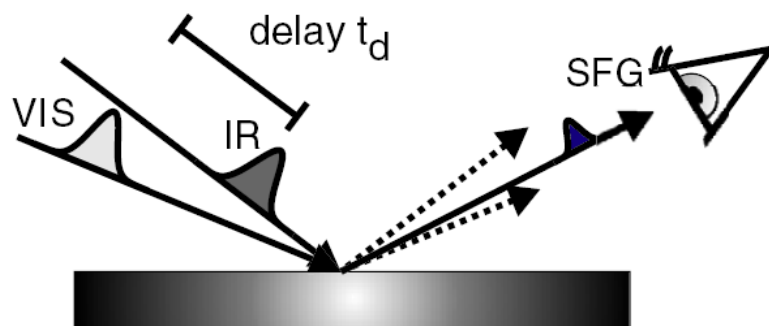
M. Morkel et al. Surf. Sci. **586**, 146 (2005).

SFG

Time resolved measurements



Max-Planck-Gesellschaft



**What happens if both pulses
are fs in duration?**

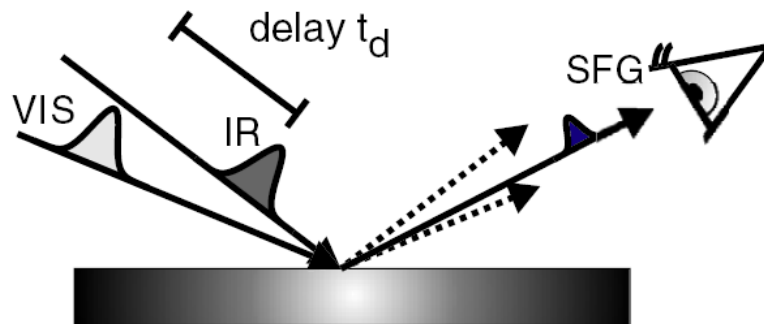
M. Bonn *et al.* J. Phys.: Condens. Mat. **17**, S201 (2005).

SFG

Time resolved measurements



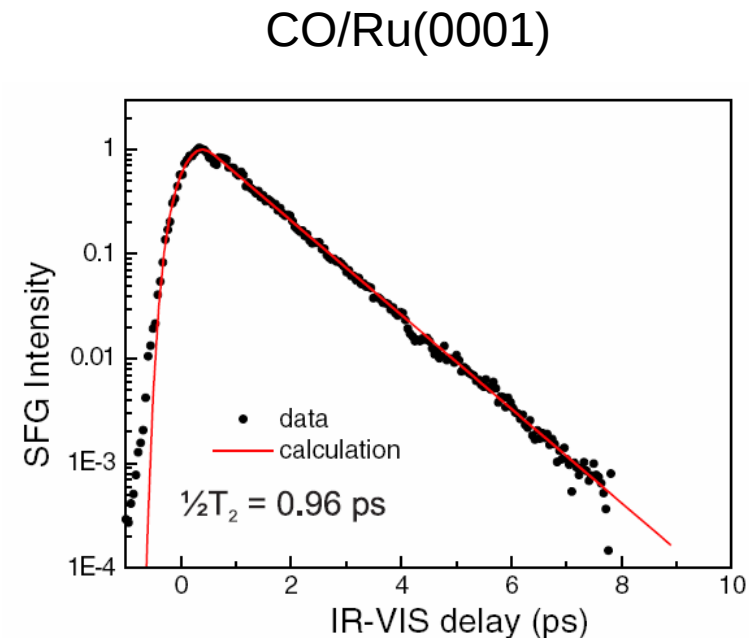
Max-Planck-Gesellschaft



What happens if both pulses are fs in duration?

⇒ Loss of spectral resolution!

For 200 fs spectral width: 165 cm^{-1}



M. Bonn *et al.* J. Phys.: Condens. Mat. **17**, S201 (2005).

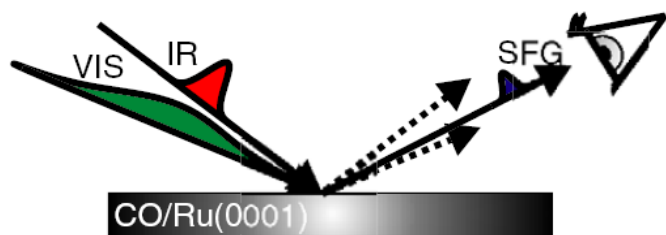
SFG

Time resolved measurements



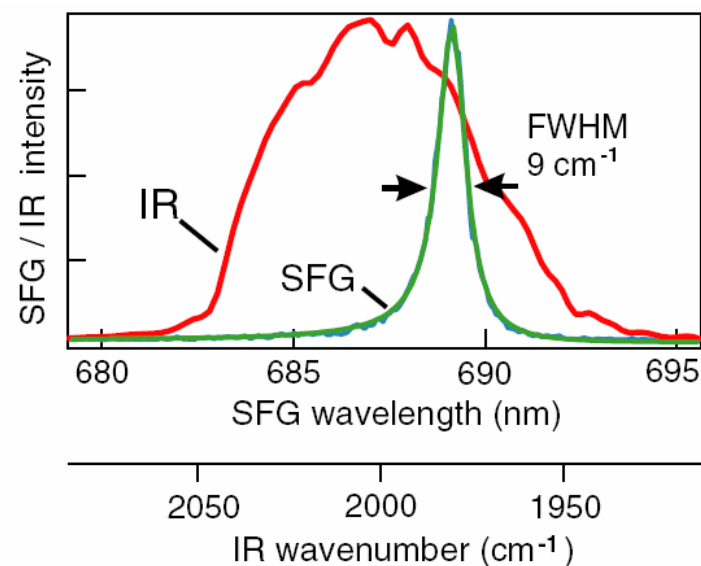
Max-Planck-Gesellschaft

How to measure a frequency domain spectrum?



fs IR pulse: broad spectral width
long VIS pulse (ps; pulse shaper)
(narrow spectral width approx. 7 cm^{-1})

CO/Ru(0001)



M. Bonn *et al.* J. Phys.: Condens. Mat. **17**, S201 (2005).

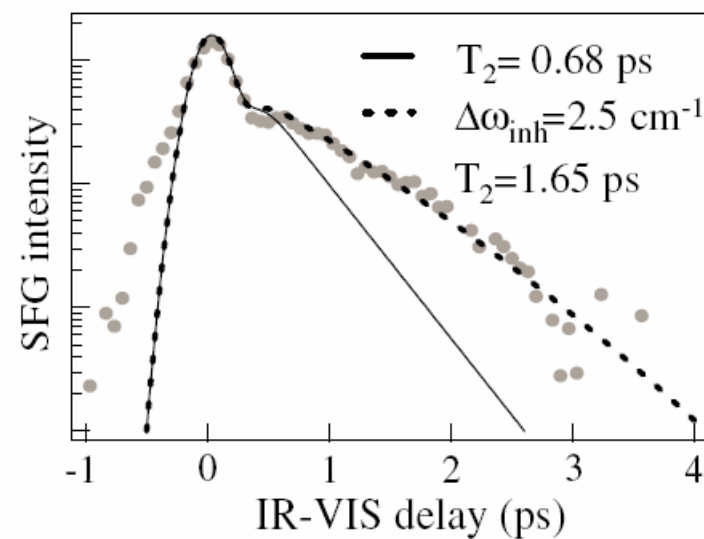
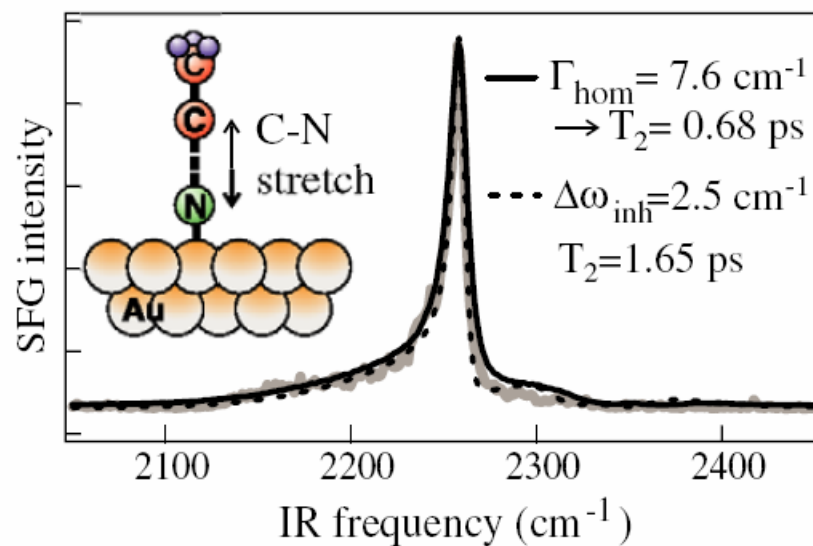
SFG

Time resolved measurements



Max-Planck-Gesellschaft

What method is better suited?



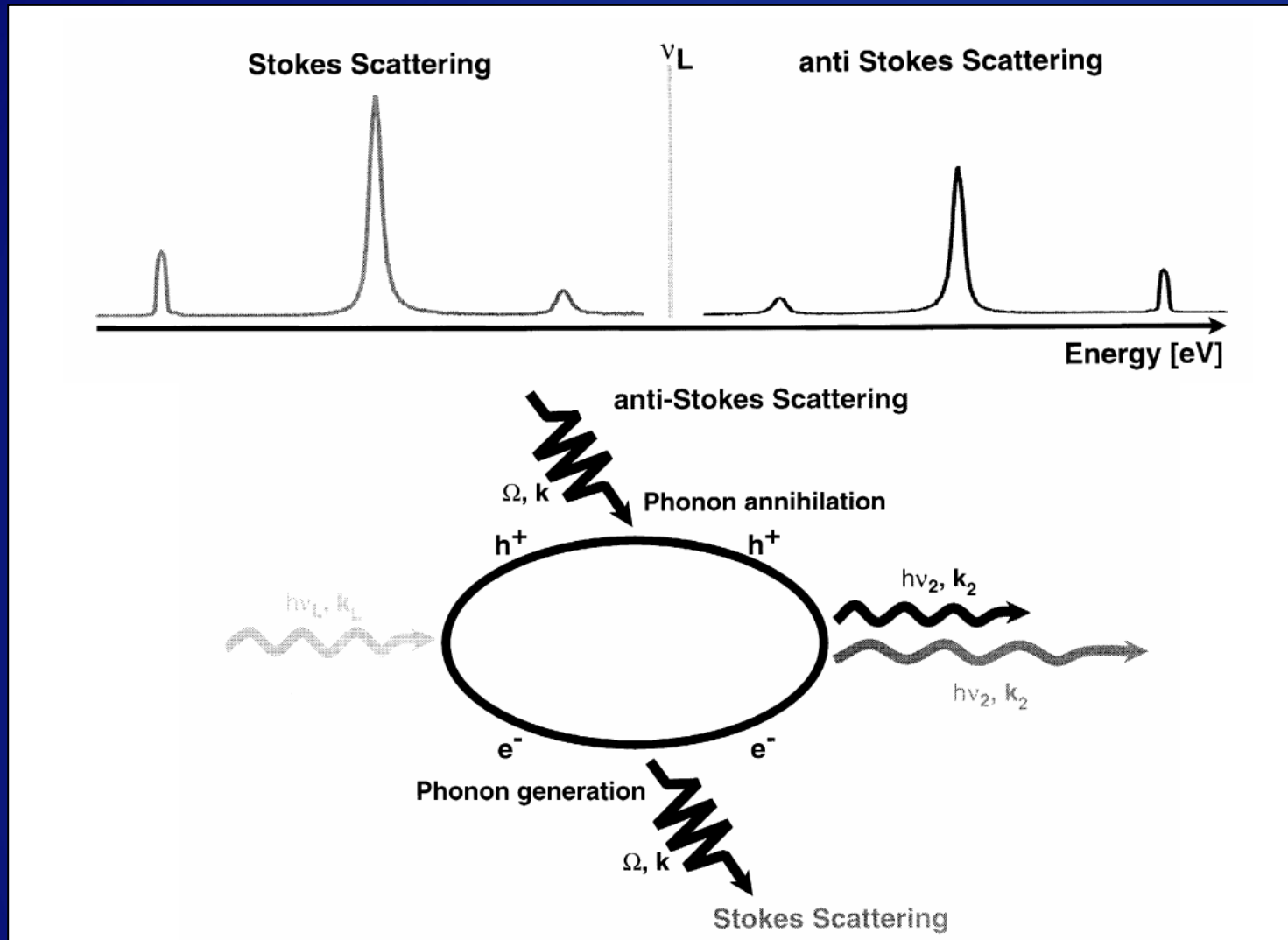
M. Bonn *et al.* J. Phys.: Condens. Mat. **17**, S201 (2005).

Raman spectroscopy

Basic aspects



Max-Planck-Gesellschaft



Raman spectroscopy

Basic aspects



Max-Planck-Gesellschaft

Pro

- wide spectral range 50 – 5000 cm^{-1}
- negligible gas phase scattering
- quartz is a very weak Raman scatterer; cells/windows can be made out of quartz
- can be done also at high temperatures (e.g. 1000 °C) because of detection in the optical regime; no perturbation by blackbody radiation
- catalysis: typical oxide supports (silica, alumina) are weak scatterers

Con

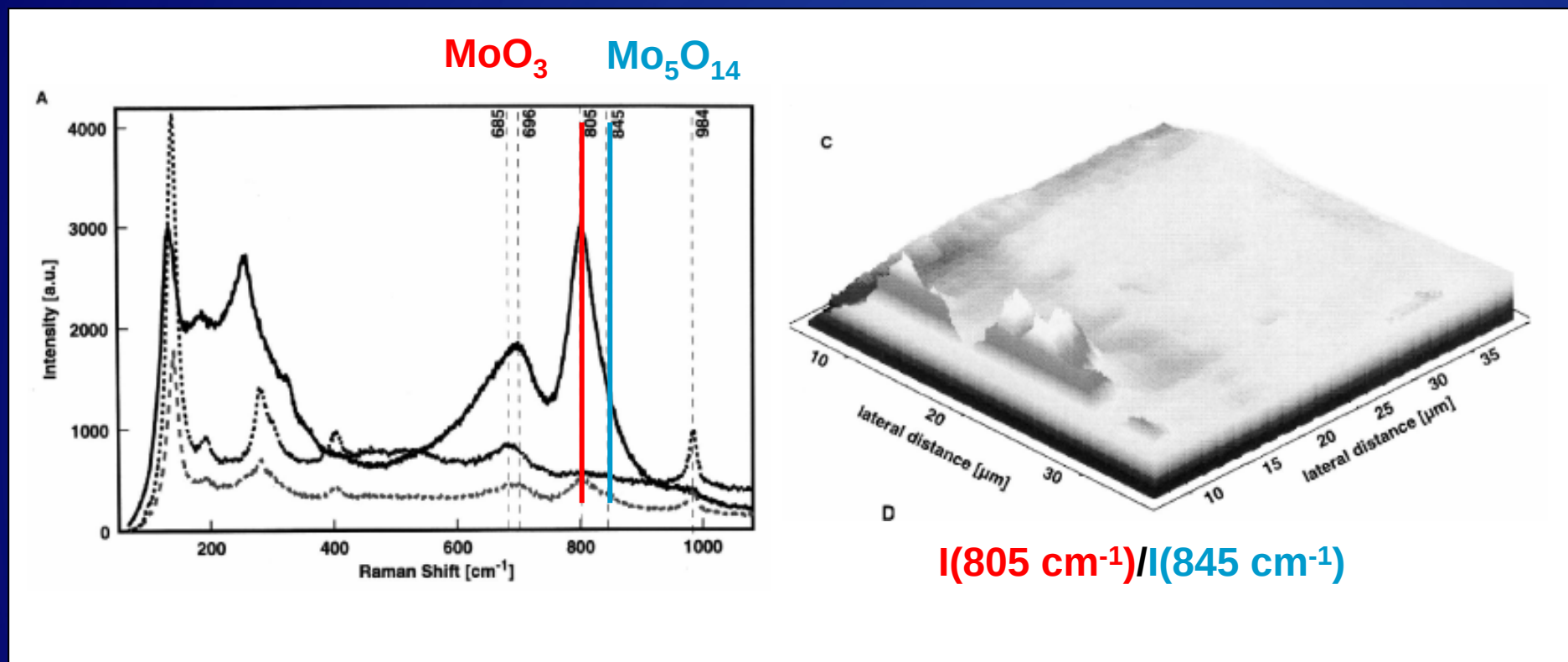
- low intrinsic cross section
- $(d\sigma/d\Omega)_{\text{NRS}} \approx 10^{-28} \text{ cm}^2 \text{ sr}^{-1} (I \propto \nu^4)$
- susceptible to fluorescence (can be up to 10^6 higher; especially coke has a high fluorescence yield)
- heating due to intense lasers
- quantification of Raman intensities is very difficult (even with reference samples, because of possible electronic effects of the substrate)

Raman spectroscopy

Material characterization



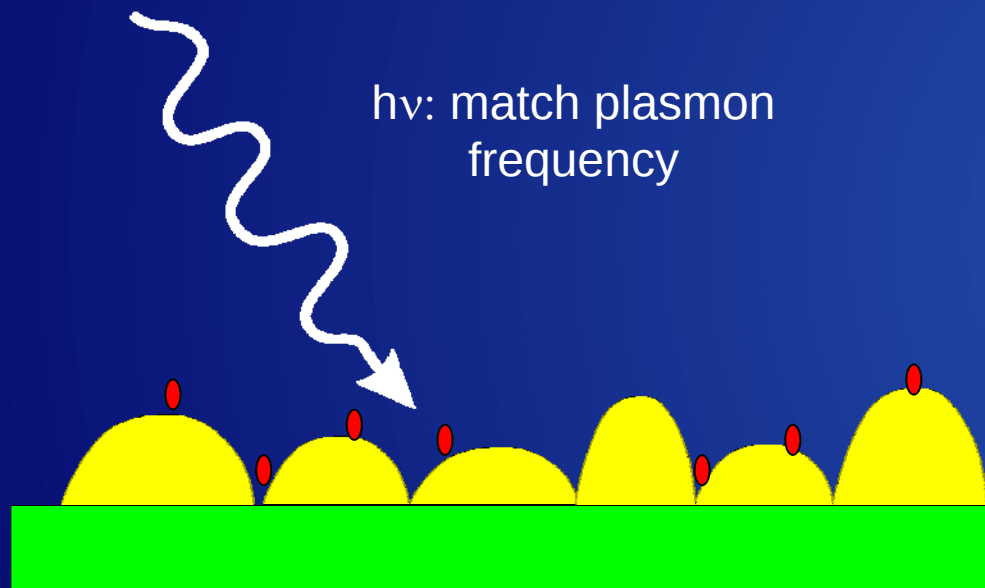
Max-Planck-Gesellschaft



G. Mestl, J. Mol. Catal. A 158, 45 (2000).

Raman spectroscopy

Surface enhancement, SERS



Works for rough surfaces of coinage metals (Cu, Ag, Au)

Idea: excitation of local surface plasmons

Field enhancement by plasmons in particular in between particles (gap states)

Enhancement: average effect approx 10^7 (gap mode may influence only a small number of molecules $\Rightarrow$ local effects even much higher)

Raman spectroscopy

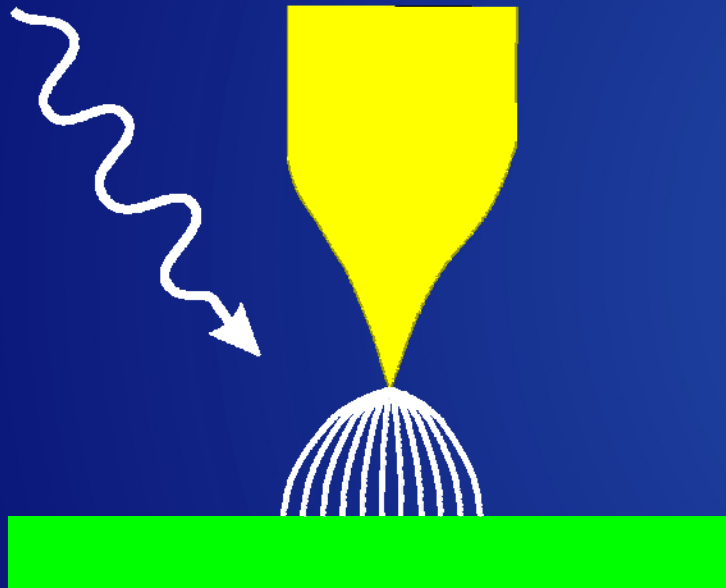
Surface enhancement, TERS

**Idea: using the enhancement due to gap modes and
combine an STM with a Raman spectrometer
(Tip Enhanced Raman Spectroscopy)**

Raman spectroscopy

Surface enhancement, TERS

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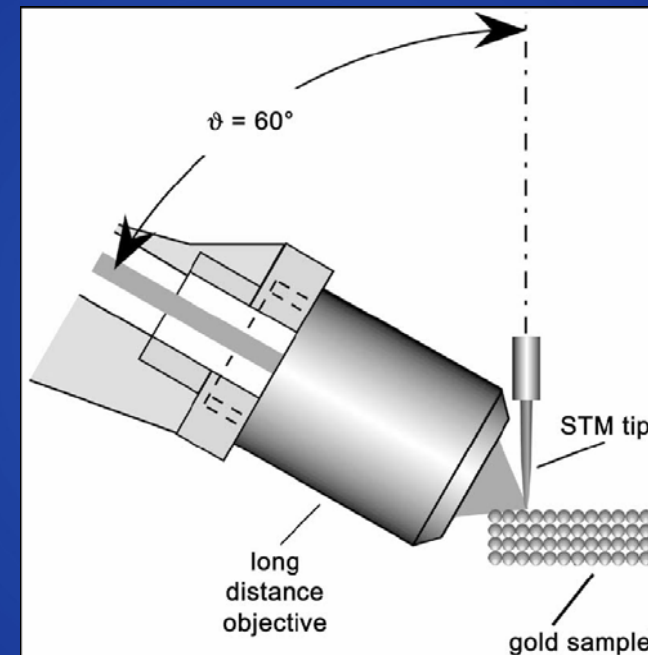
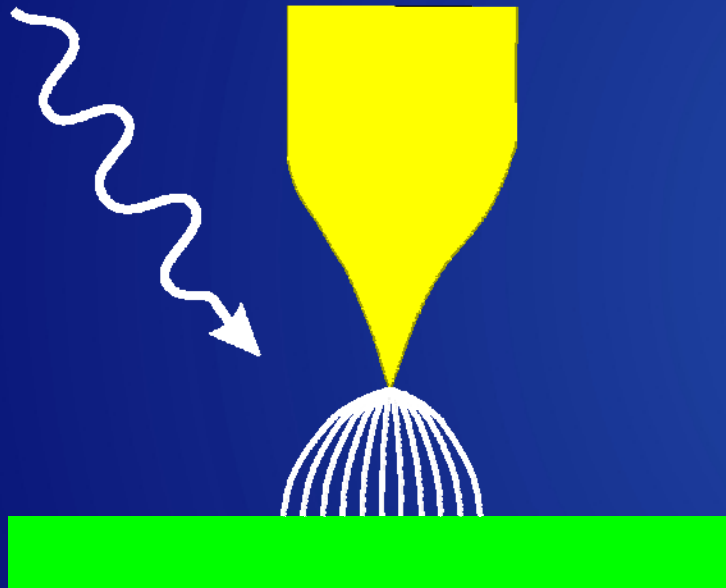


B. Pettinger et al.: Single Molec. **3**, 285 (2002).

Raman spectroscopy

Surface enhancement, TERS

**Idea: using the enhancement due to gap modes and
combine an STM with a Raman spectrometer
(Tip Enhanced Raman Spectroscopy)**



B. Pettinger et al. Phys. Rev. Lett. **92**, 096101 (2004).

Raman spectroscopy

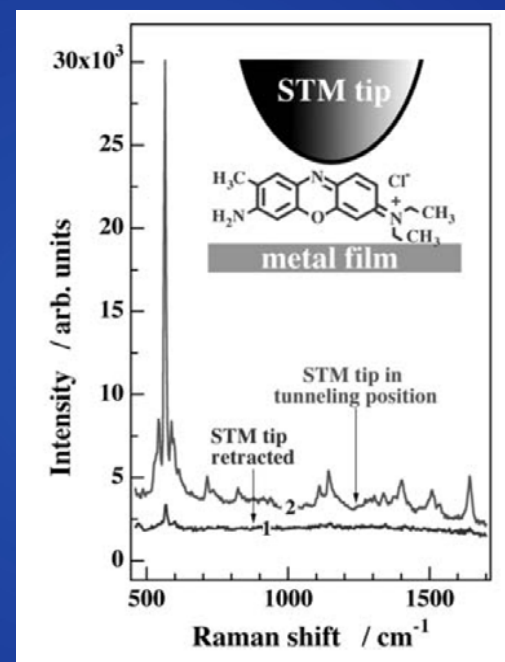
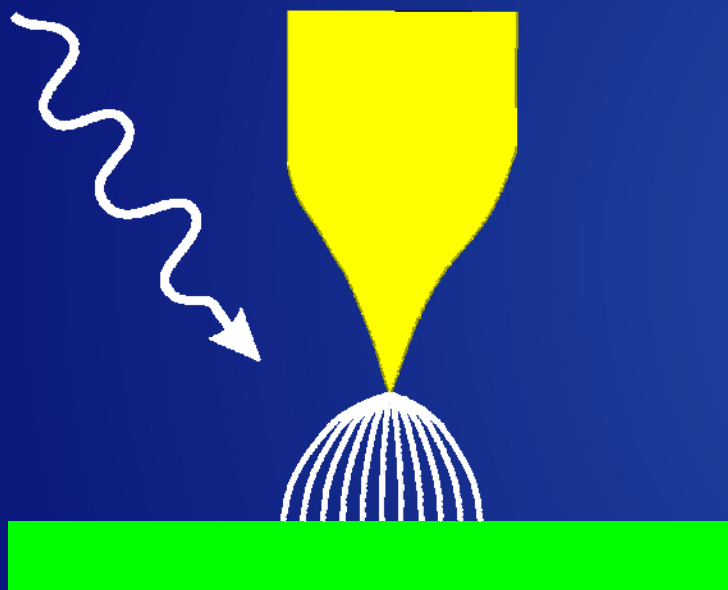
Surface enhancement, TERS



Max-Planck-Gesellschaft

Idea: using the enhancement due to gap modes and
combine an STM with a Raman spectrometer

(Tip Enhanced Raman Spectroscopy)



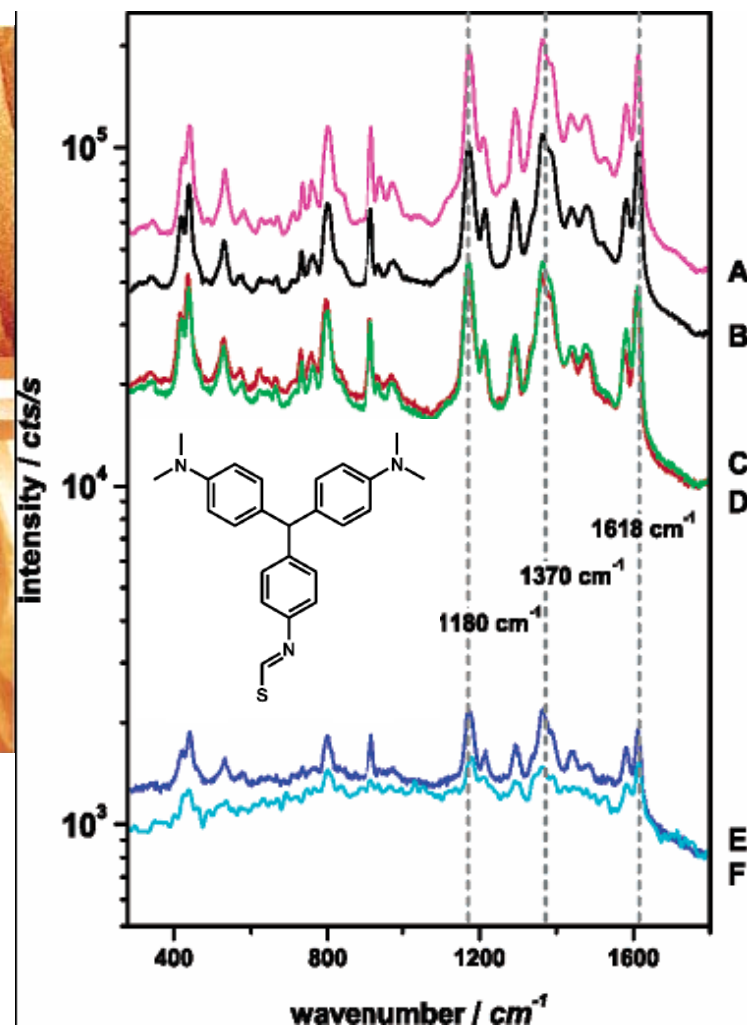
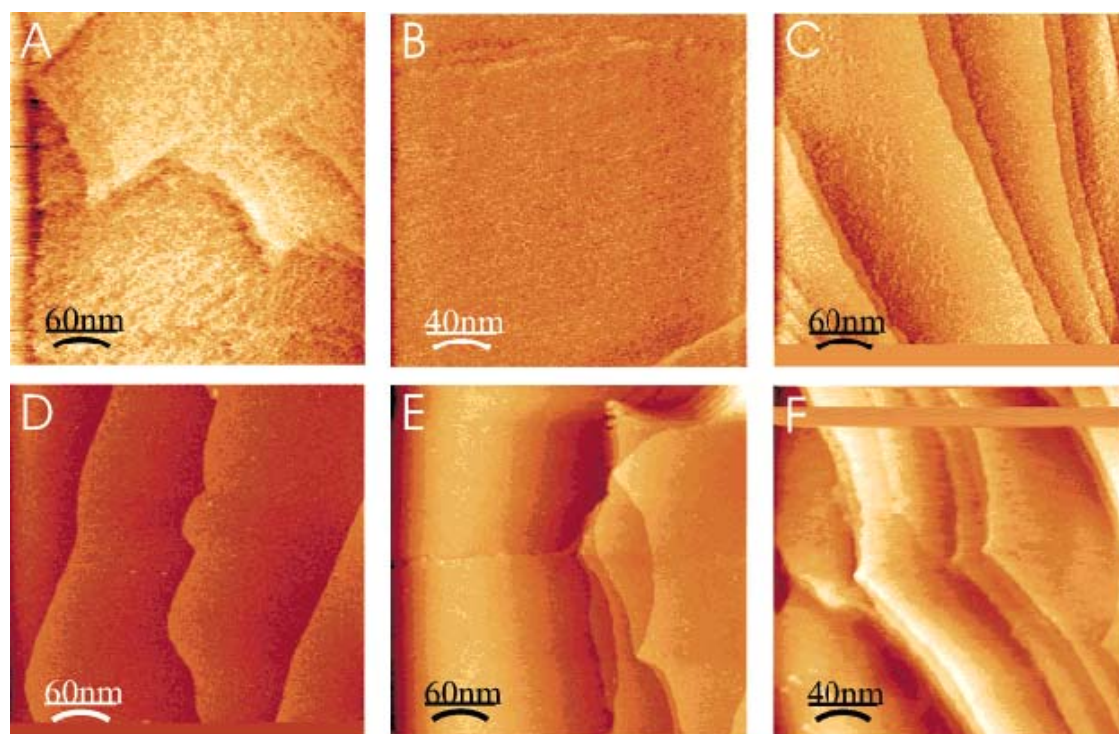
B. Pettinger et al.: Single Molec. **3**, 285 (2002).

Raman spectroscopy

Surface enhancement, TERS



Max-Planck-Gesellschaft



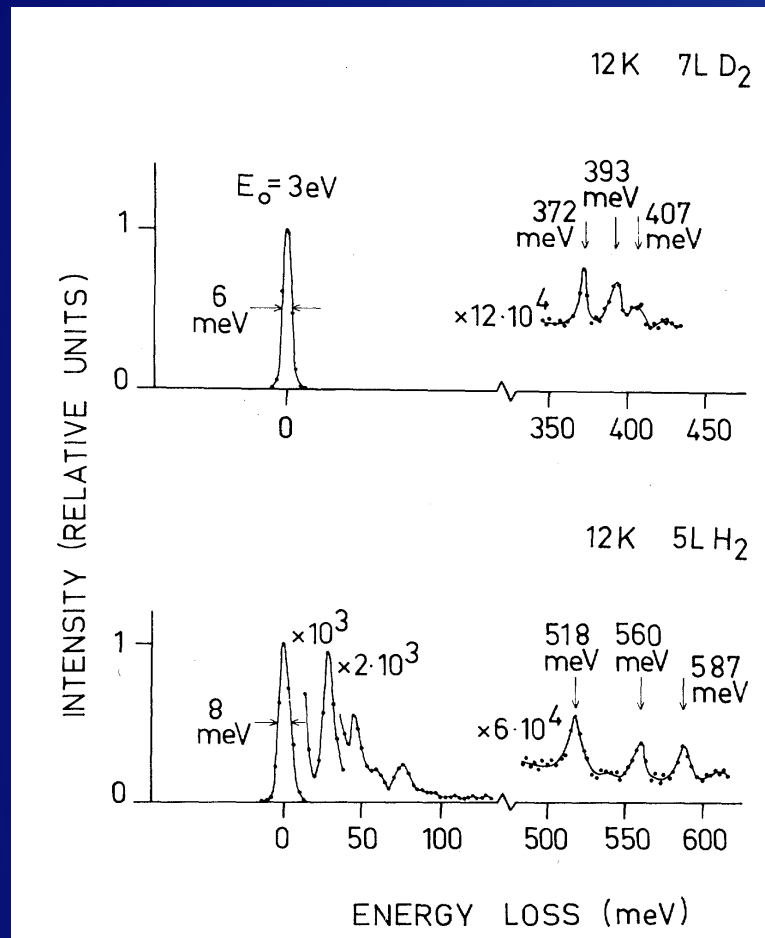
K. F. Domke et al. J. Am. Chem. Soc. **128**, 14721 (2006).

EELS

H₂ on Cu(100)



Max-Planck-Gesellschaft



S. Andersson, J. Harris, Phys. Rev. Lett. 48, 545 (1982).

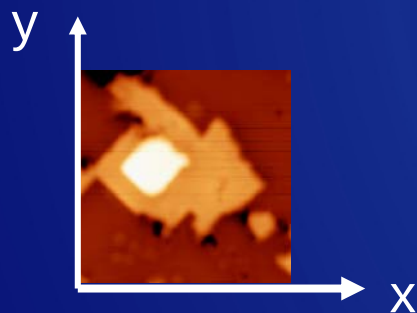
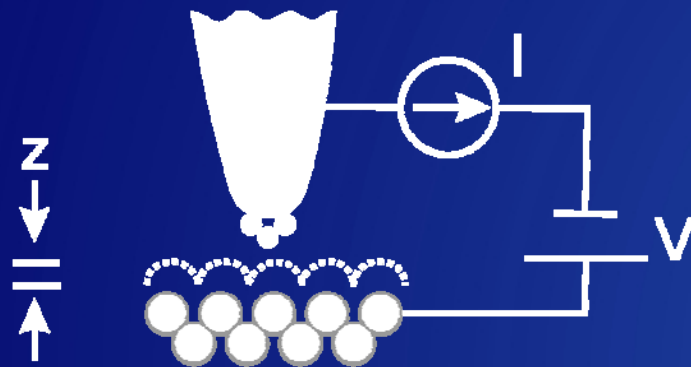
- D₂-stretch at 372 meV
- 21-23 meV is close to gas phase rotation form J(0→2)
- H₂-stretch at 518 meV
- 42 meV is close to gas phase rotation form J(0→2)
- 69 meV is close to gas phase rotation form J(1→3)

STM/STS

modes of operation

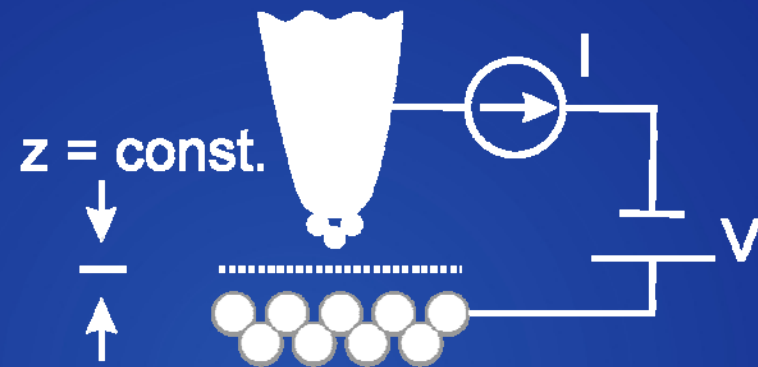
STM

constant current mode

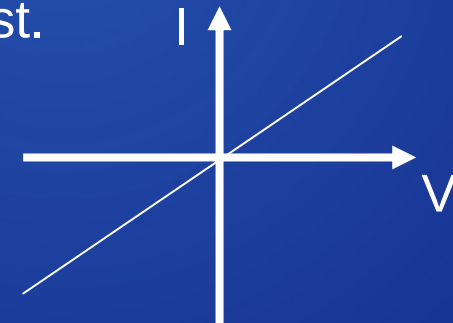


STS

constant height mode



$x, y = \text{const.}$



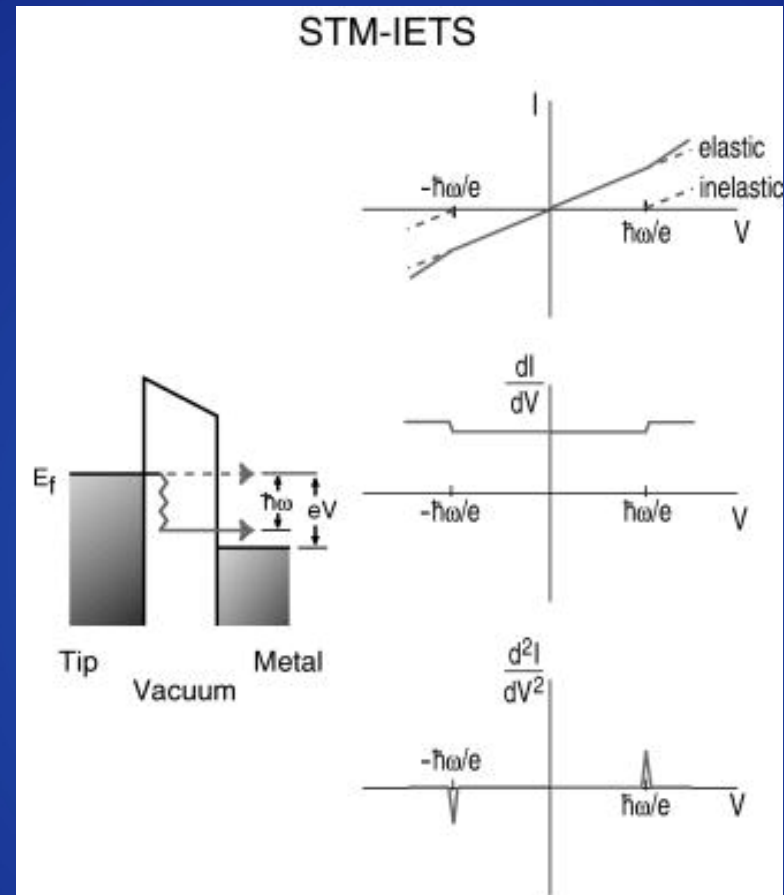
STS

inelastic tunneling processes



Max-Planck-Gesellschaft

threshold spectroscopy



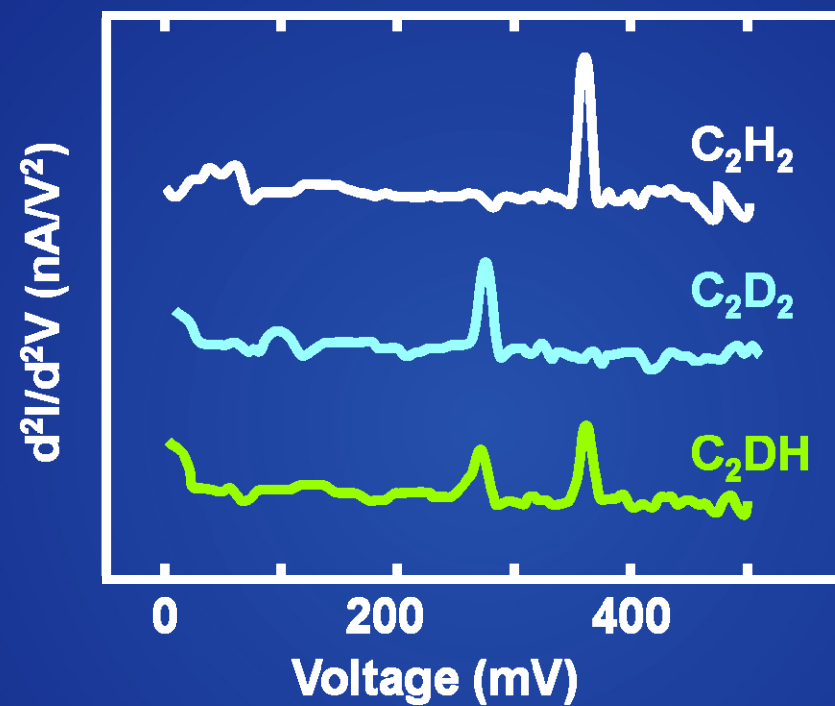
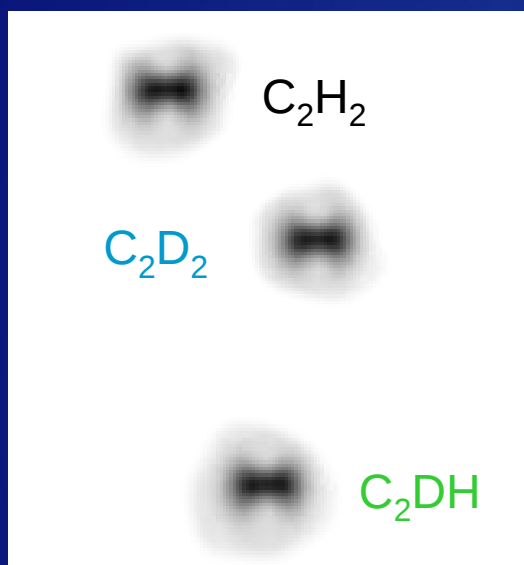
W. Ho J. Chem. Phys. 117, 11033 (2003).

STS

acetylene/Cu(001)



Max-Planck-Gesellschaft



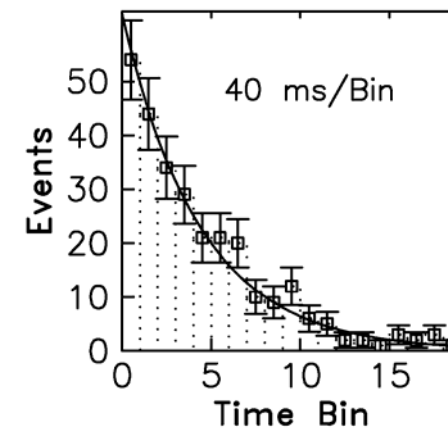
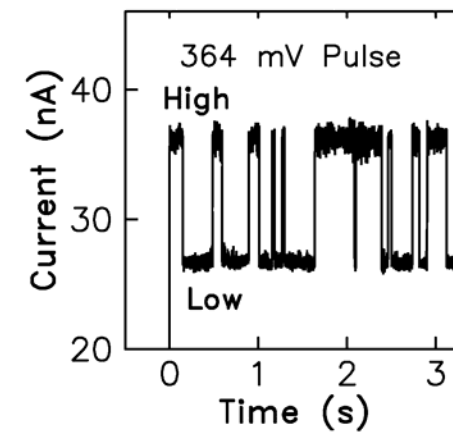
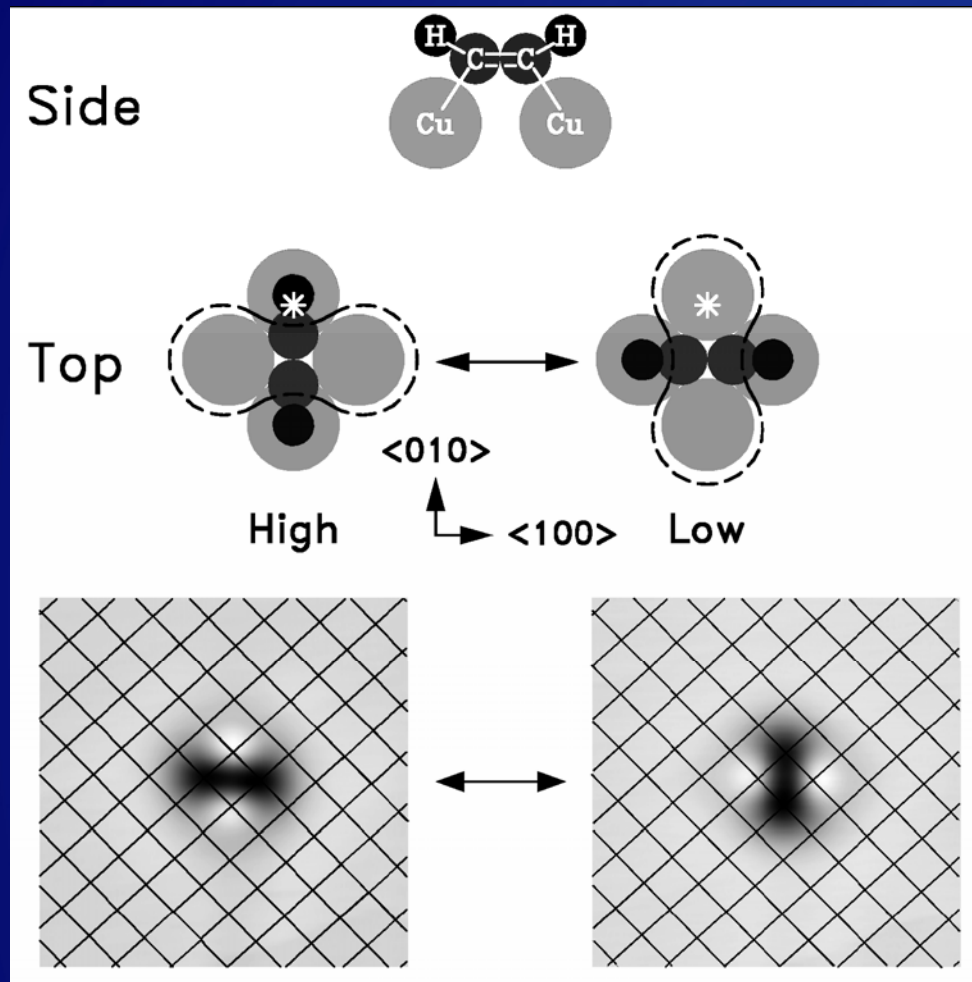
W. Ho J. Chem. Phys. 117, 11033 (2003).

STS

acetylene/Cu(001)



Max-Planck-Gesellschaft



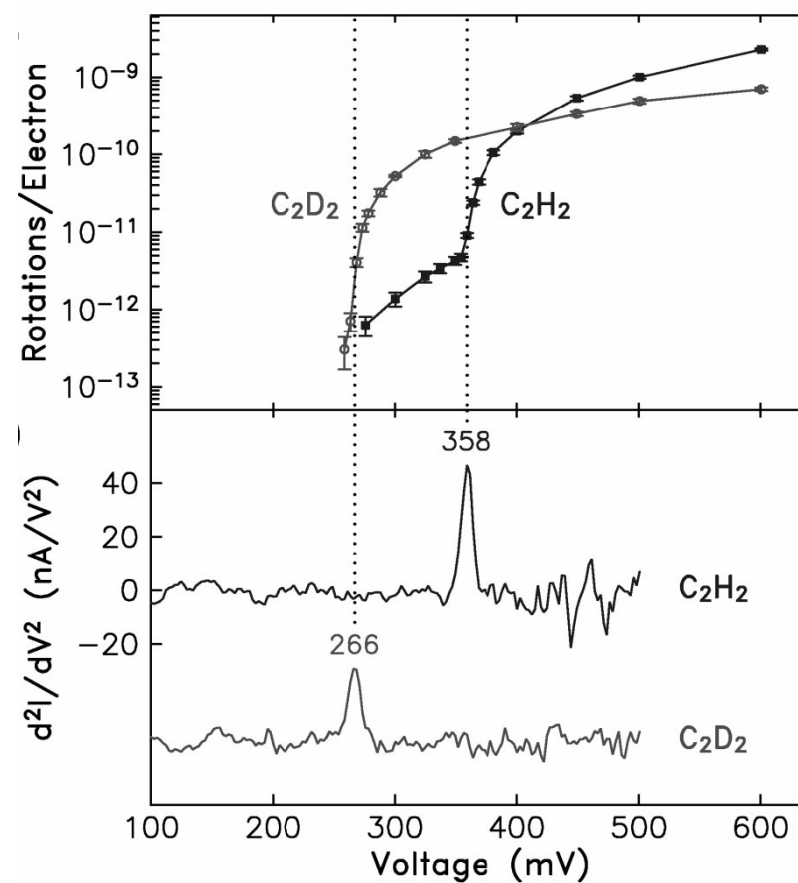
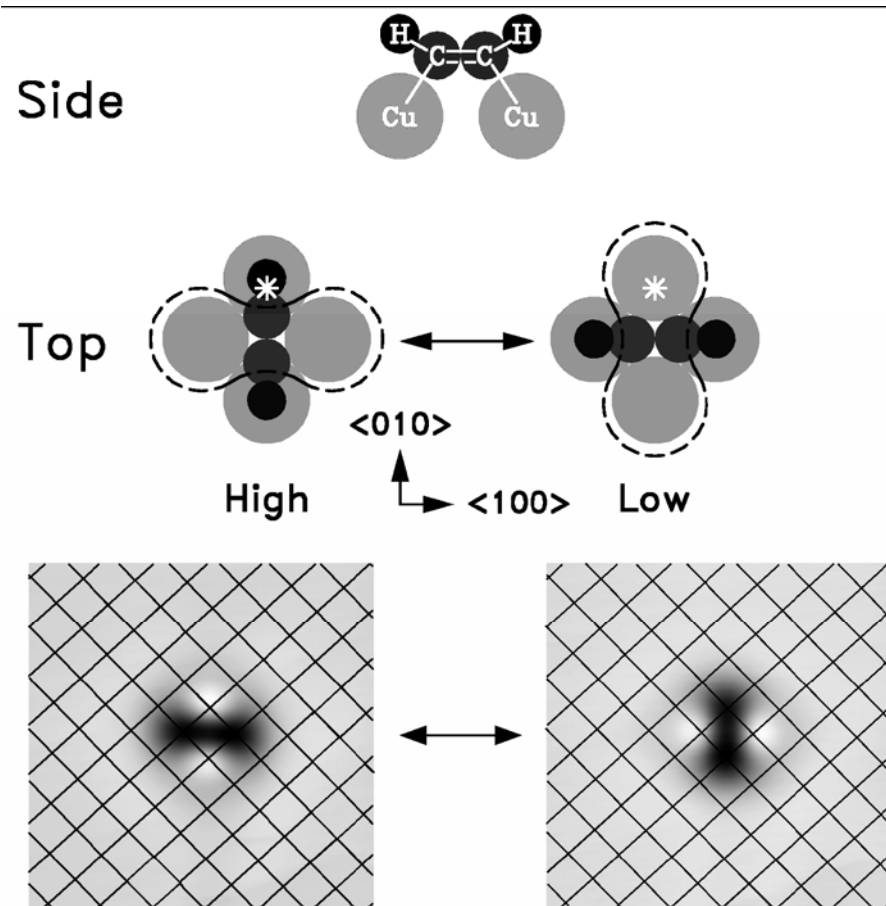
B.C. Stipe et al., Phys. Rev. Lett. 81, 1263 (1998)

STS

acetylene/Cu(001)



Max-Planck-Gesellschaft



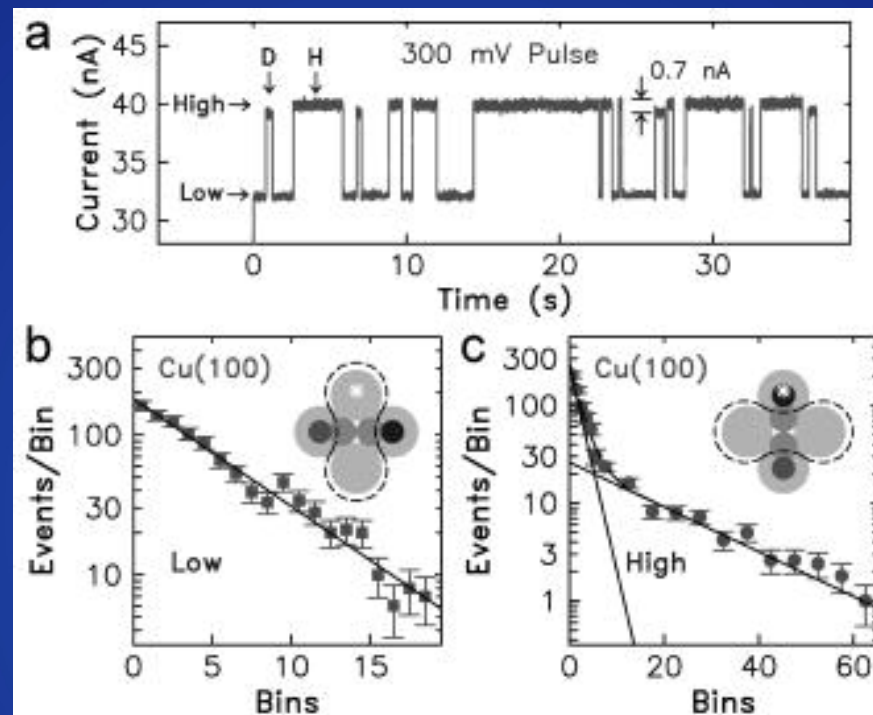
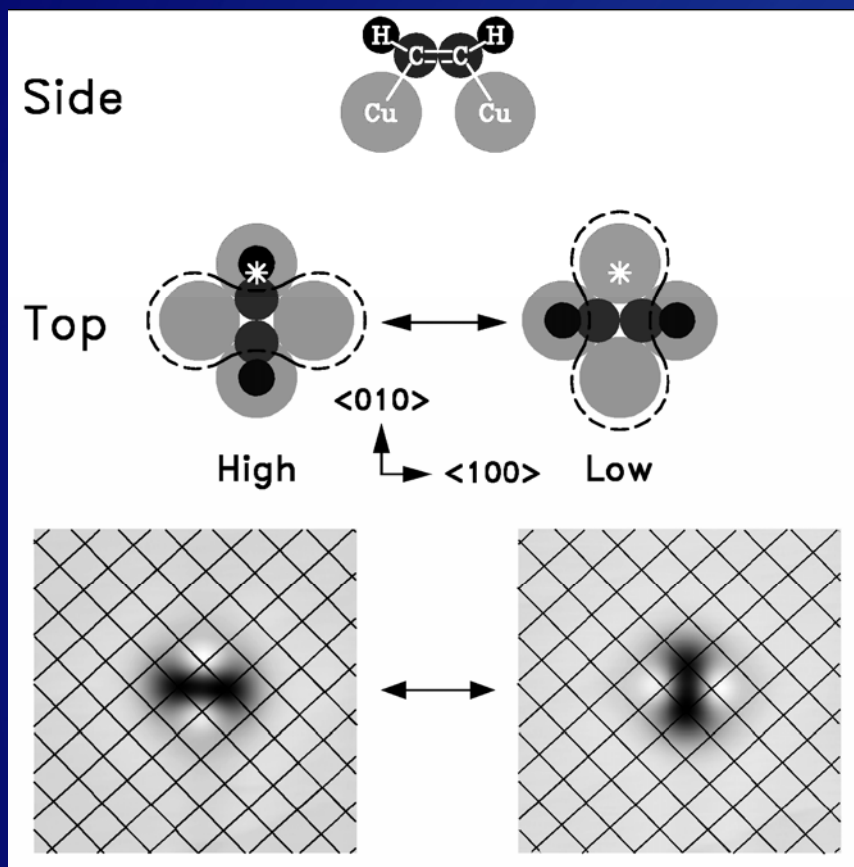
B.C. Stipe et al., Phys. Rev. Lett. 81, 1263 (1998)

STS

acetylene/Cu(001)



Max-Planck-Gesellschaft



B.C. Stipe et al., Phys. Rev. Lett. 81, 1263 (1998)

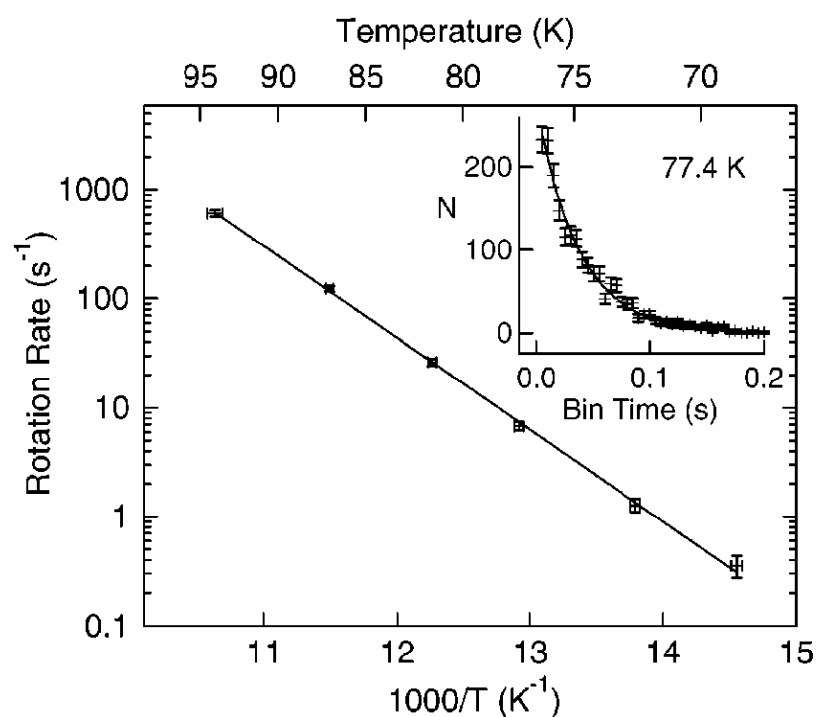
STS

acetylene/Cu(001)



Max-Planck-Gesellschaft

temperature induced rotation



$$E_a = 169 \text{ meV}; \nu = 10^{11.8} \text{ s}^{-1}$$

L. J. Lauhon and W. Ho, J. Chem. Phys. 111, 5633 (1999)

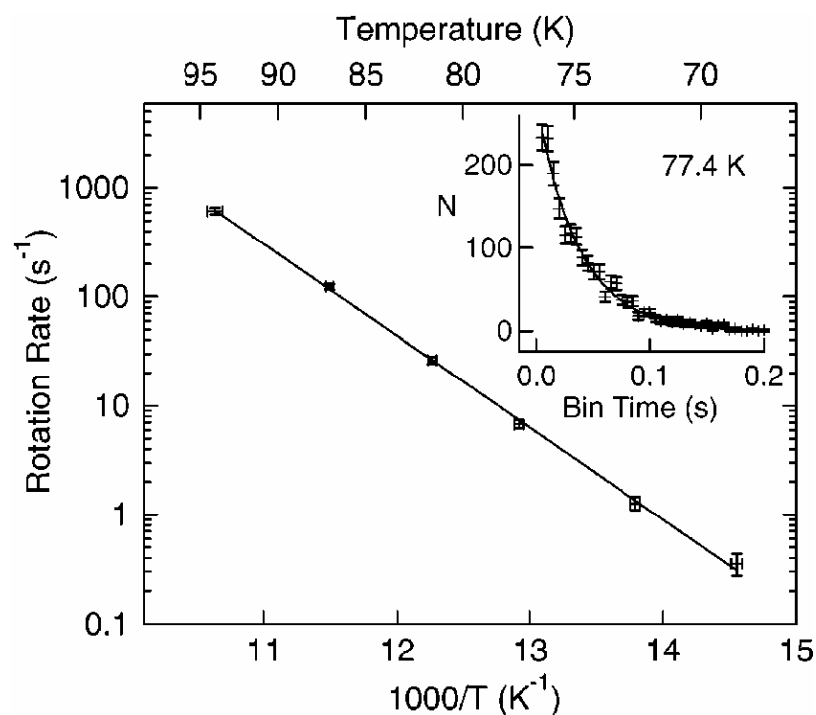
STS

acetylene/Cu(001)



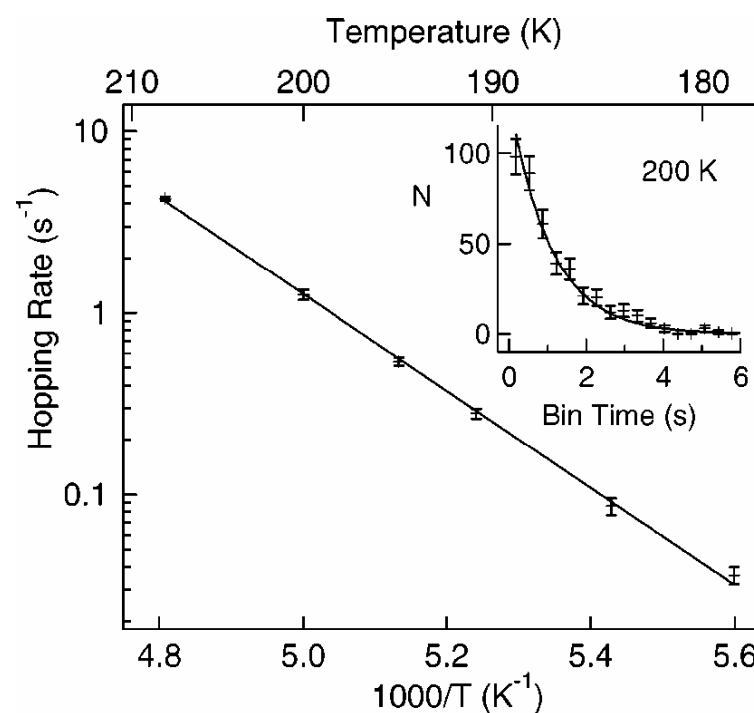
Max-Planck-Gesellschaft

temperature induced rotation



$$E_a = 169 \text{ meV}; \nu = 10^{11.8} \text{ s}^{-1}$$

temperature induced translation



$$E_a = 560 \text{ meV}; \nu = 10^{13.6} \text{ s}^{-1}$$

L. J. Lauhon and W. Ho, J. Chem. Phys. 111, 5633 (1999)

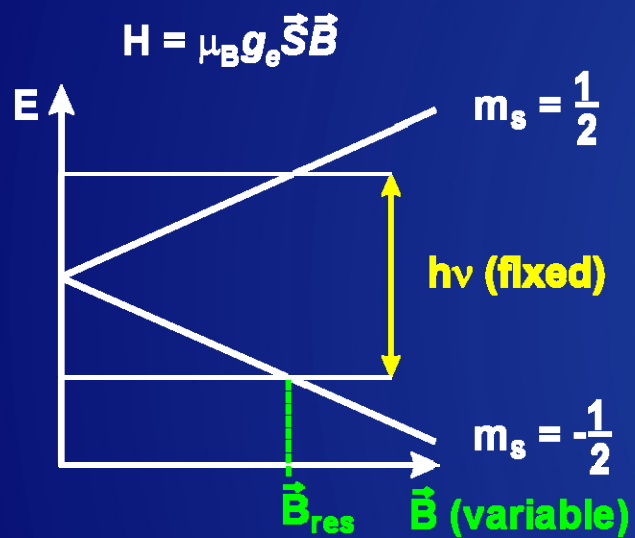
EPR spectroscopy

basic aspects



Max-Planck-Gesellschaft

free electron



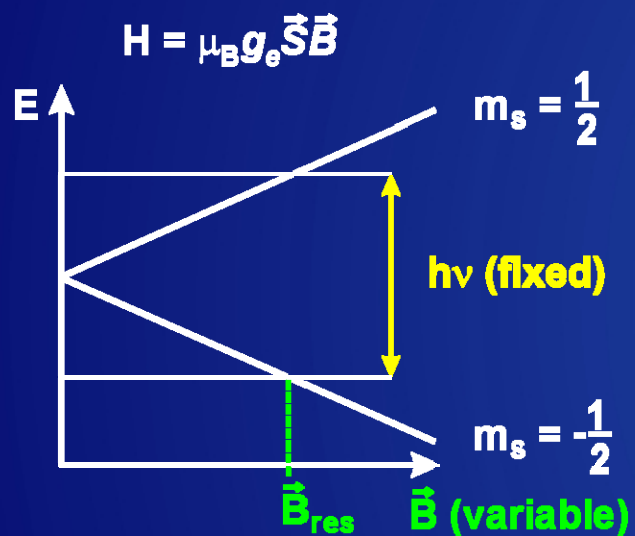
EPR spectroscopy

basic aspects

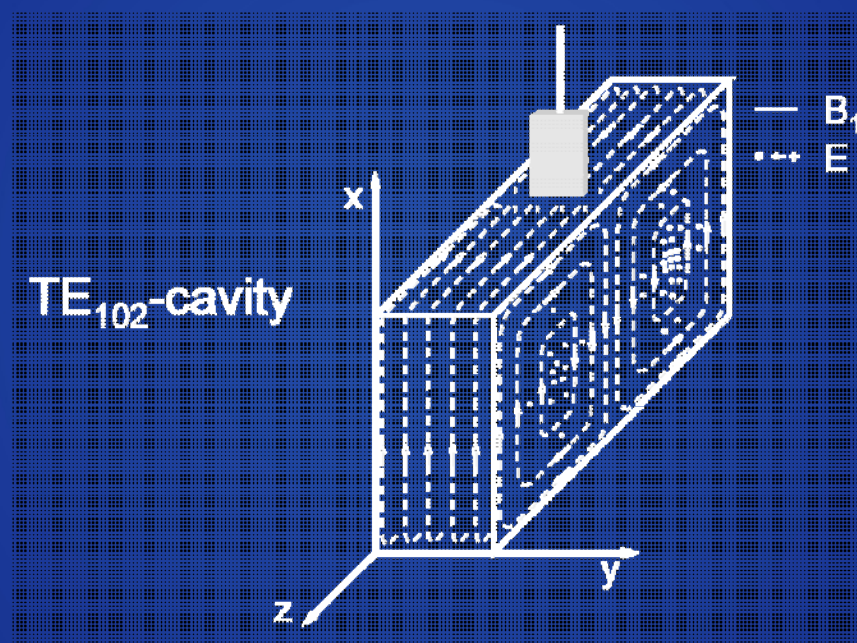


Max-Planck-Gesellschaft

free electron



field distribution in a microwave resonator



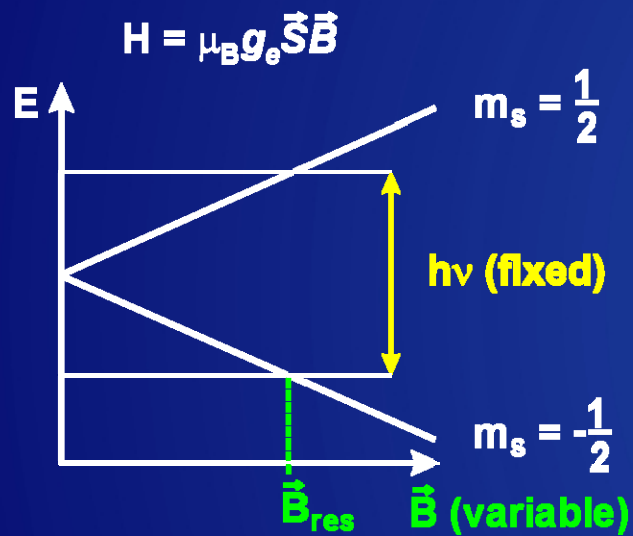
EPR spectroscopy

basic aspects

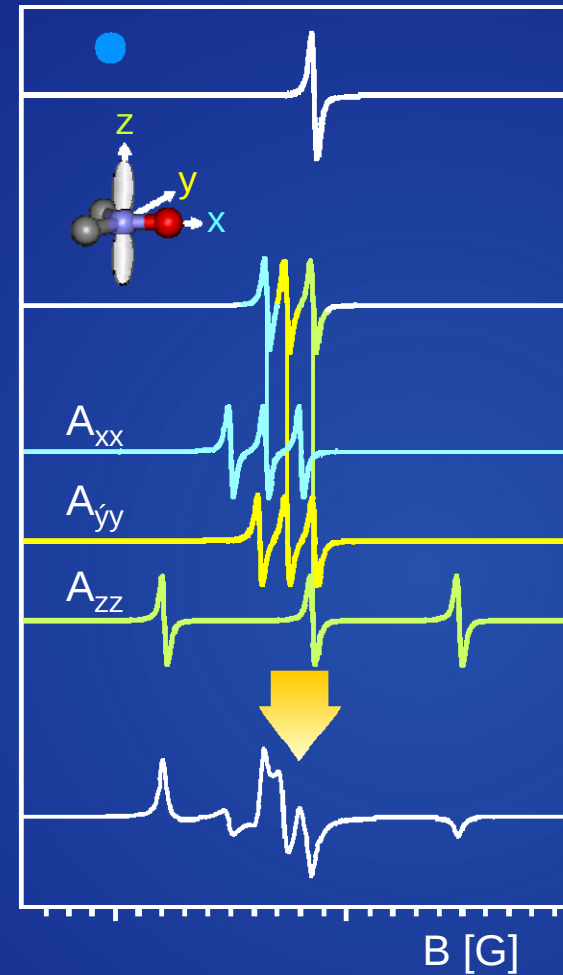


Max-Planck-Gesellschaft

free electron



rigid limit



free electron

g-tensor

$B \parallel x$

$B \parallel y$

$B \parallel z$

hyperfine interaction

$I(^{14}\text{N}) = 1$

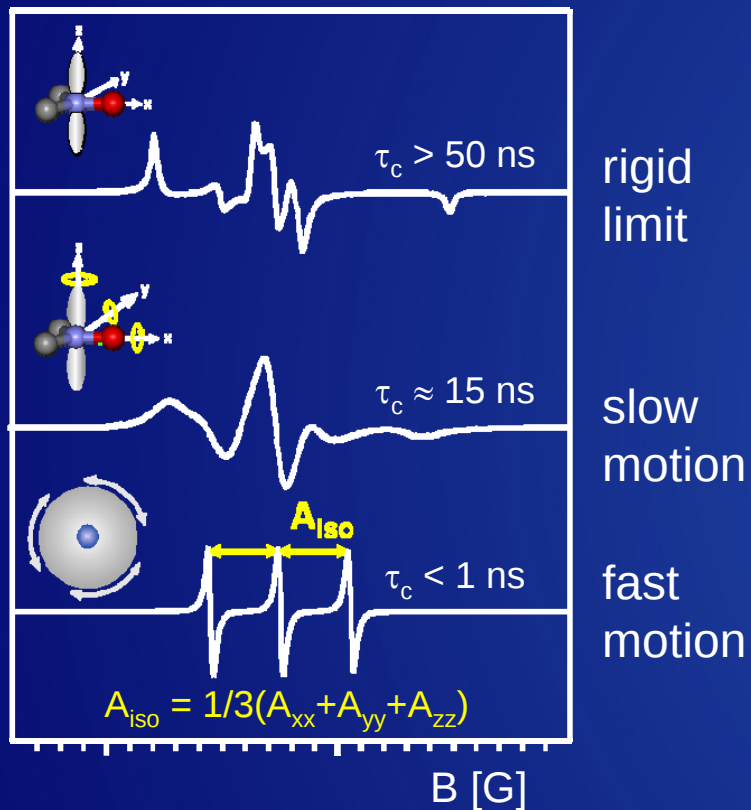
powder spectrum

EPR spectroscopy

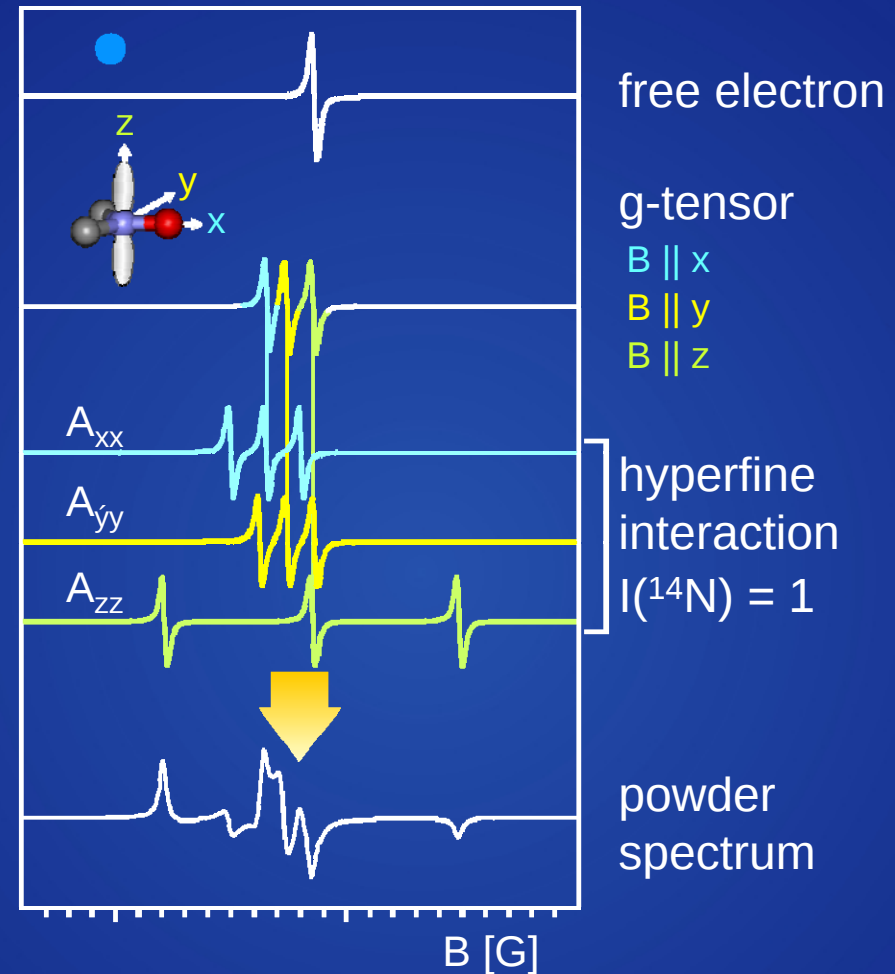
effect of rotational motion

rotational motion

isotropic rotation



rigid limit



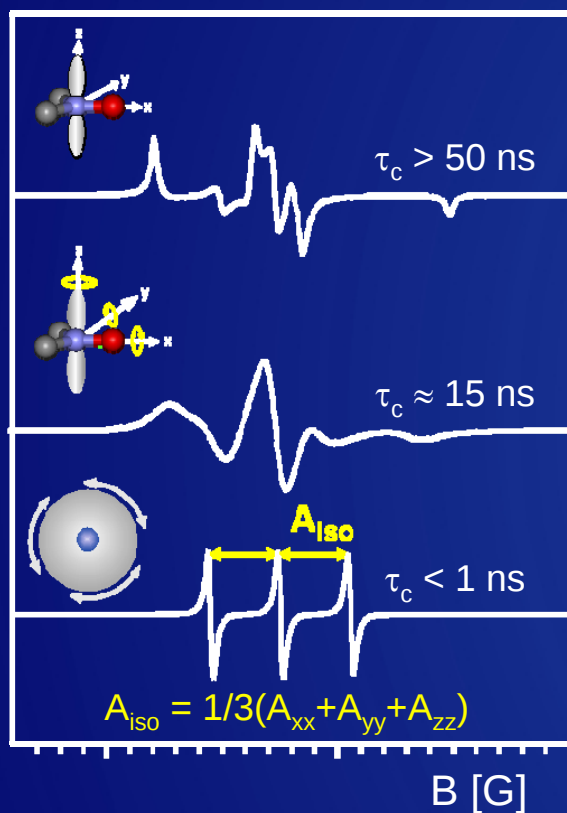


Max-Planck-Gesellschaft

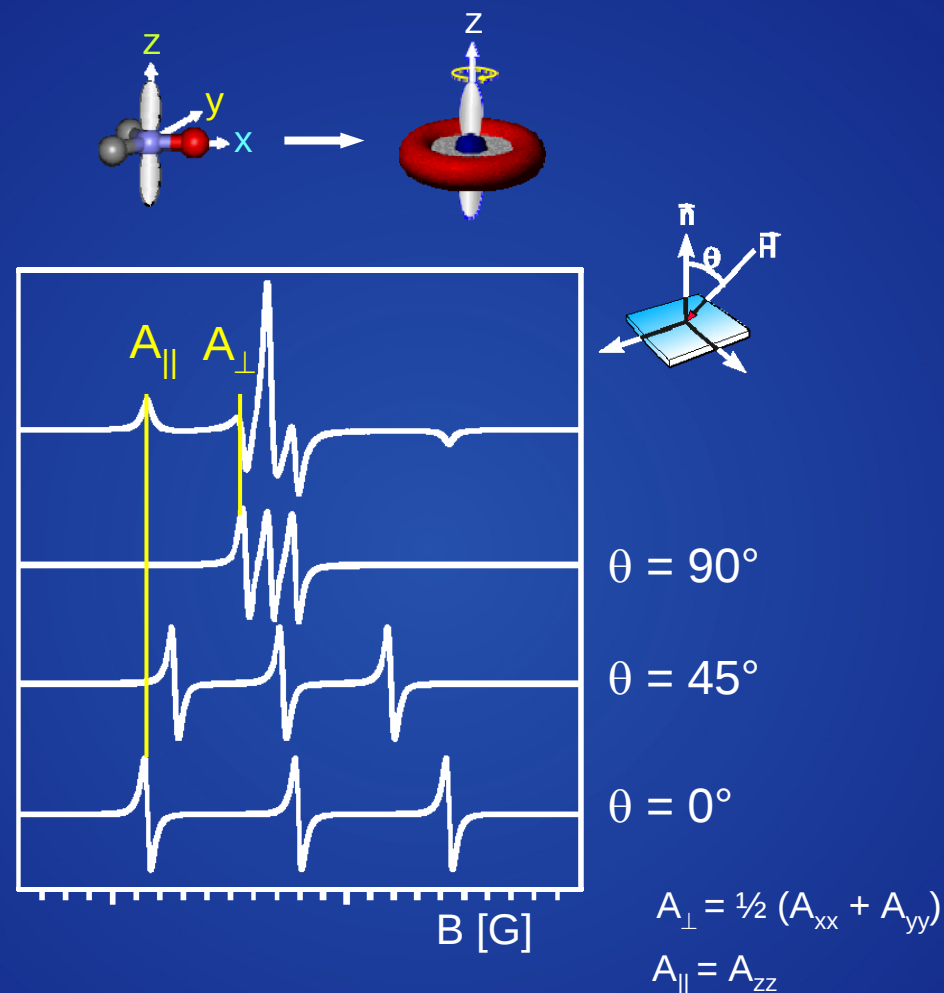
EPR spectroscopy

effects of rotational motion

isotropic rotation



fast motion around z-axis



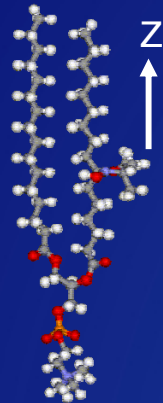
Introduction

Why ESR spectroscopy?

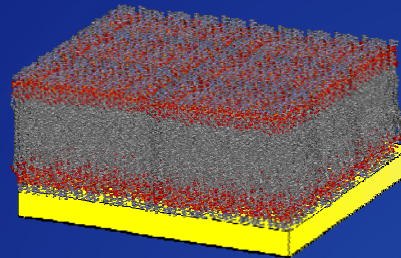


Max-Planck-Gesellschaft

(7-doxyl-PSPC)
spin labeled lipid

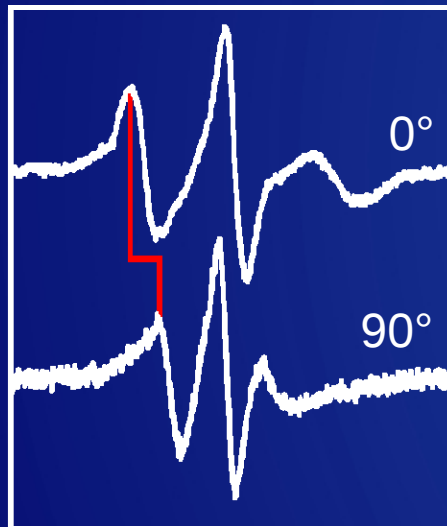
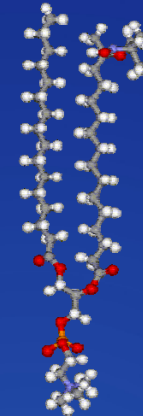


fluid lipid bilayer

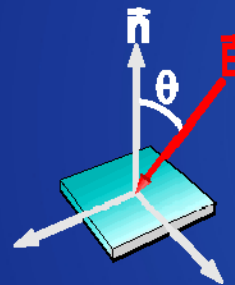


quartz surface

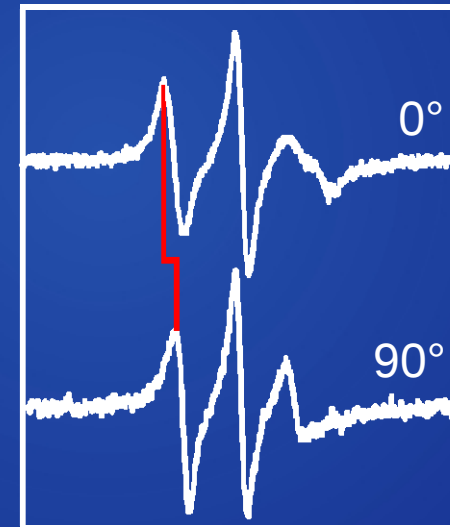
(14-doxyl-PSPC)
spin labeled lipid



B →



Increase of the rotational
dynamics along the
alkyl chain

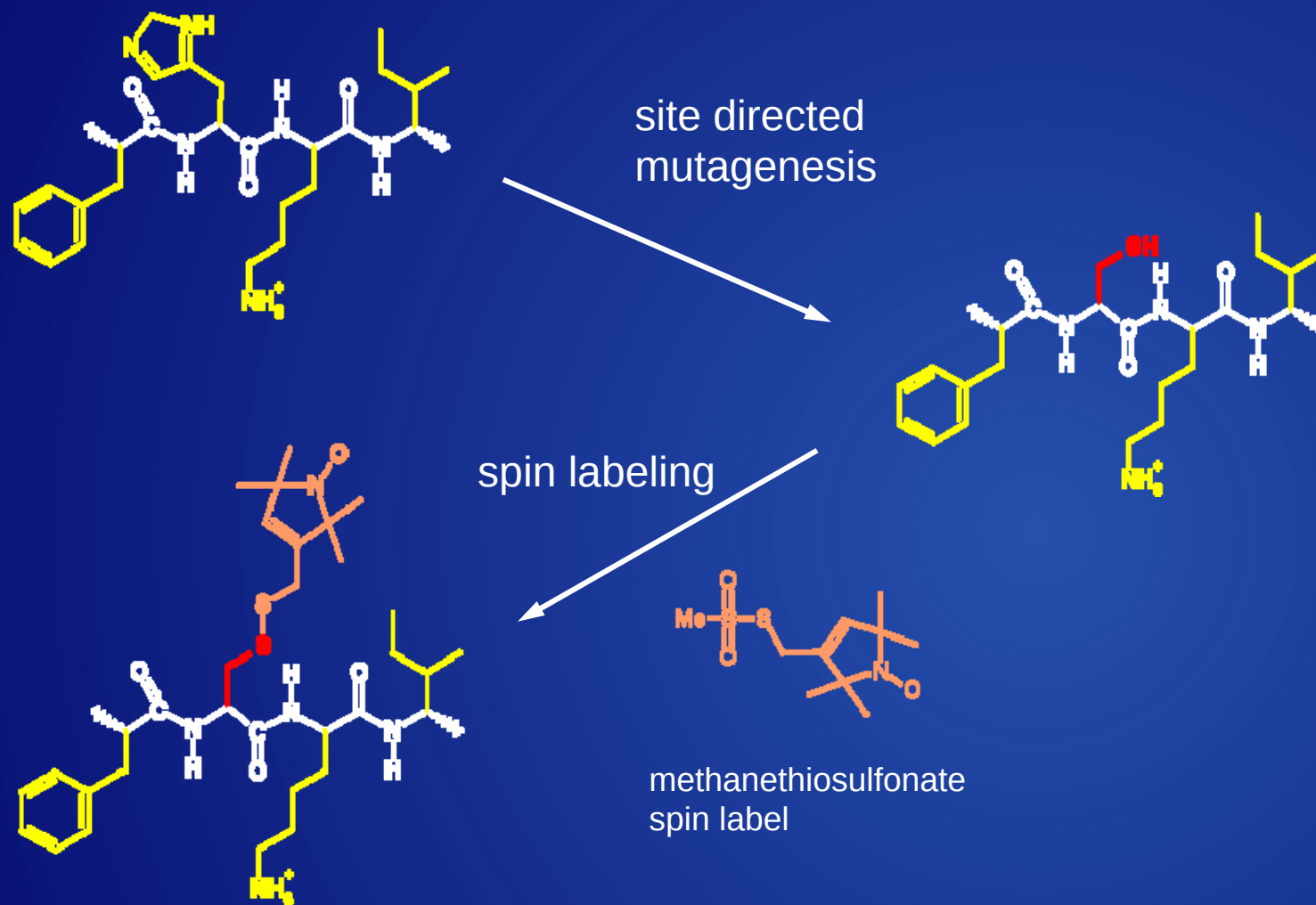


B →

Site Directed Spin Labeling (SDSL) strategy



Max-Planck-Gesellschaft



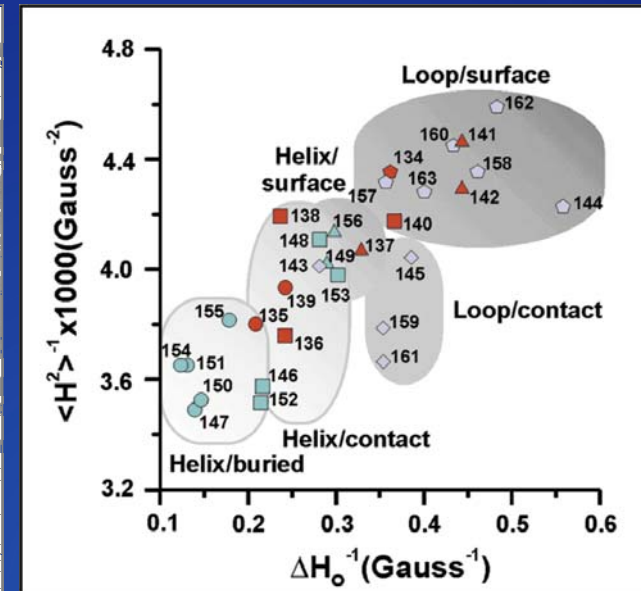
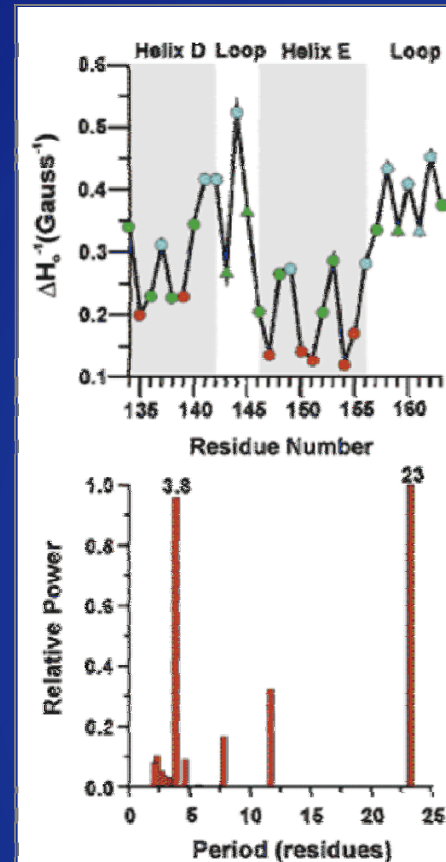
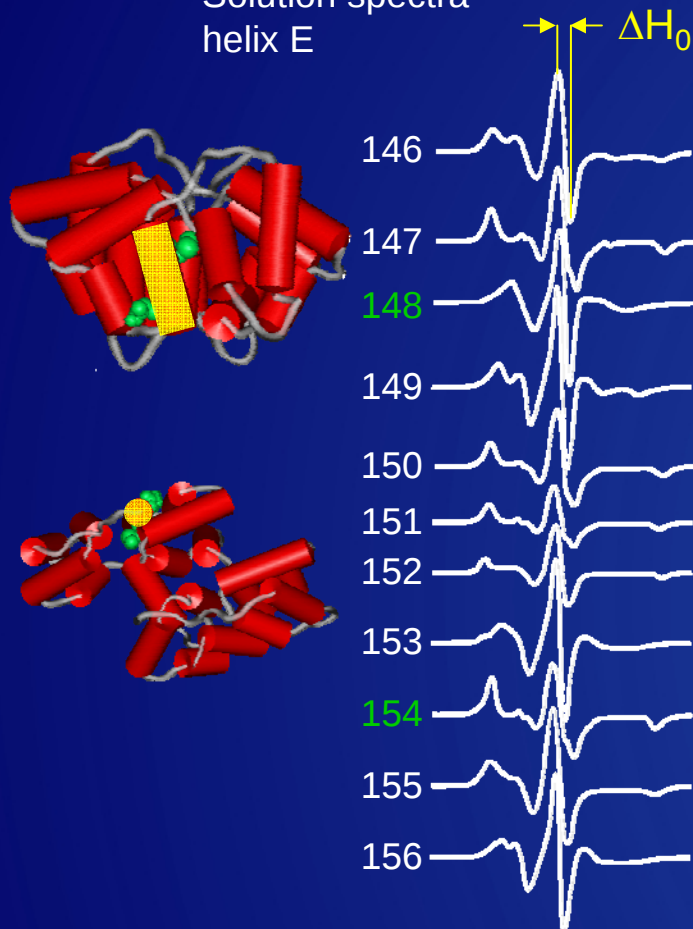
EPR Annexin 12

secondary structure from line shape analysis



Max-Planck-Gesellschaft

Solution spectra
helix E



EPR data: Isas, J. M et al. Biochemistry **41**, 1464 (2002).

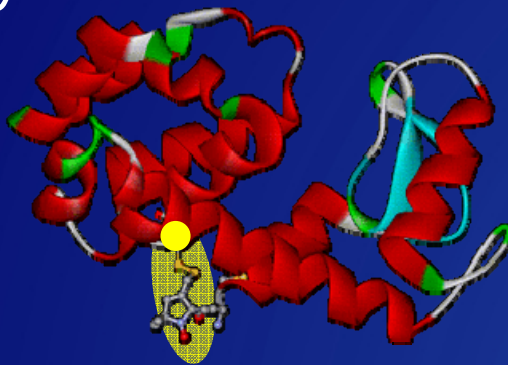
EPR T4 Lysozyme on Surface

Angular Dependence



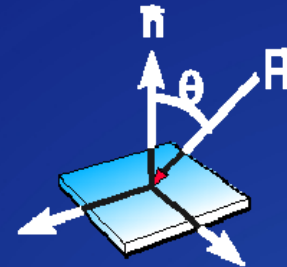
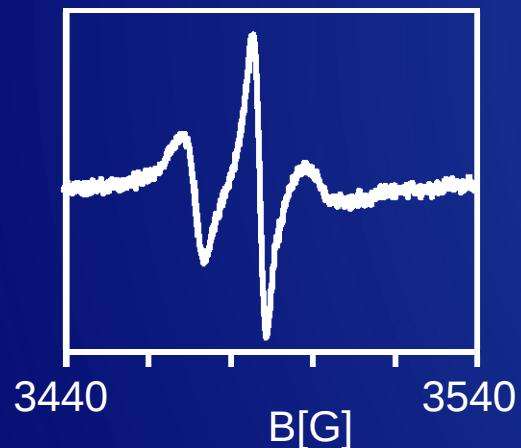
Max-Planck-Gesellschaft

76 Cys

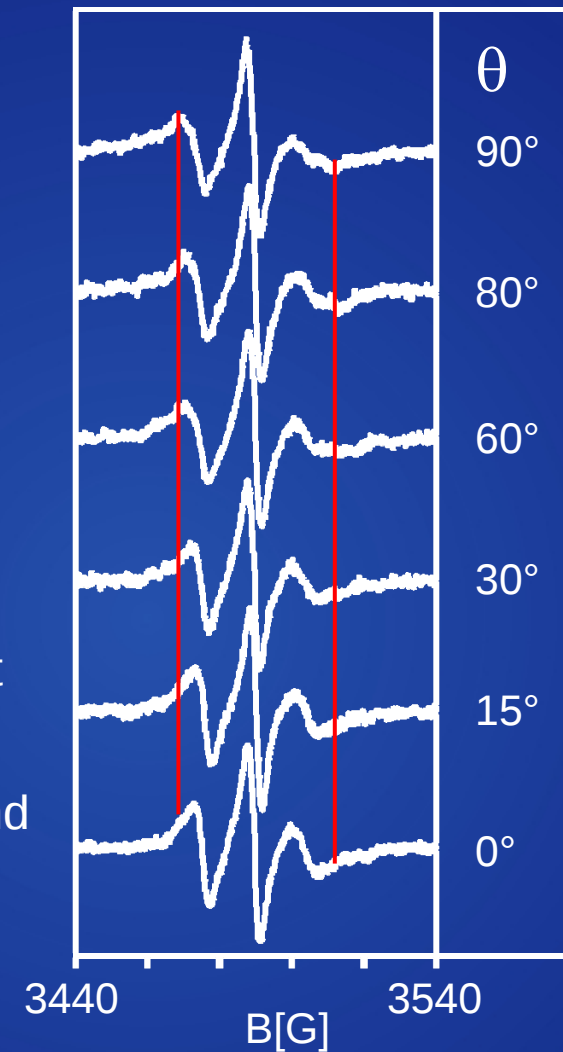


MTSSL

solution



- Angular dependent EPR spectra
- Intensity correspond to 80% of a dense protein layer



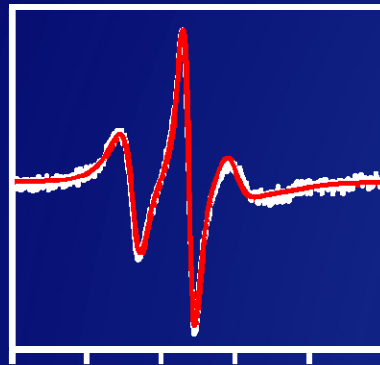
Simulations

Surface – Results: 76 Cys



Max-Planck-Gesellschaft

Dynamic model
parameters from solution



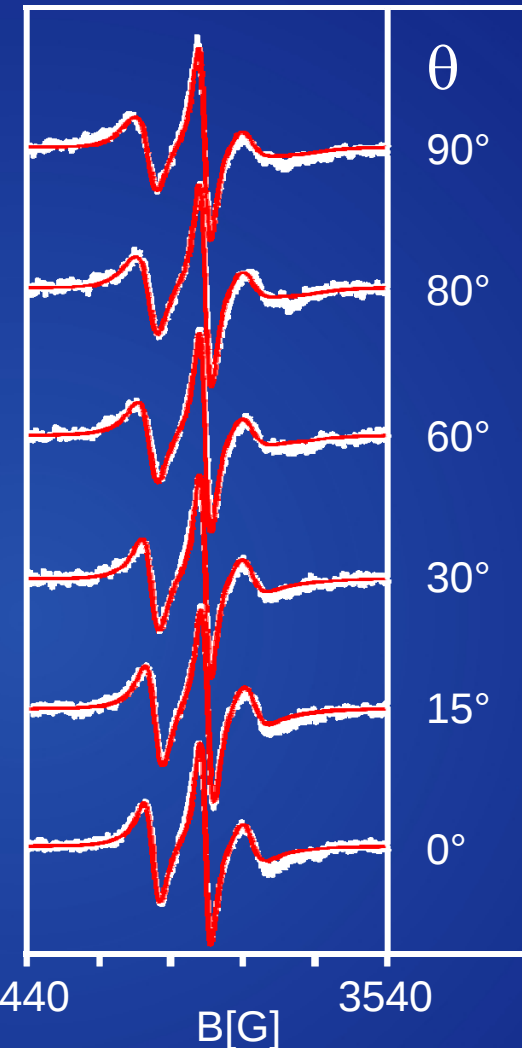
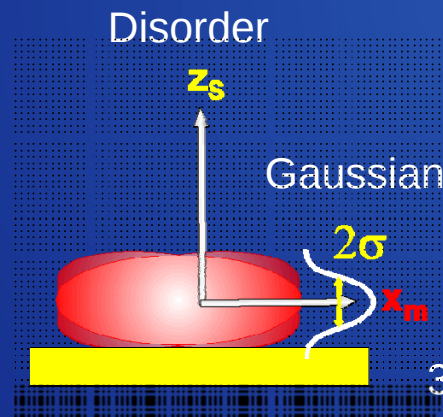
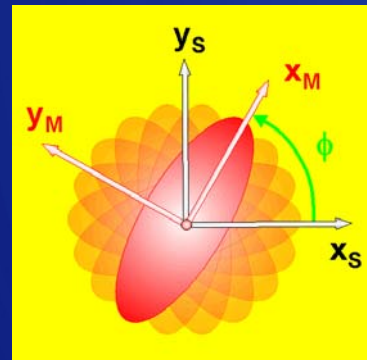
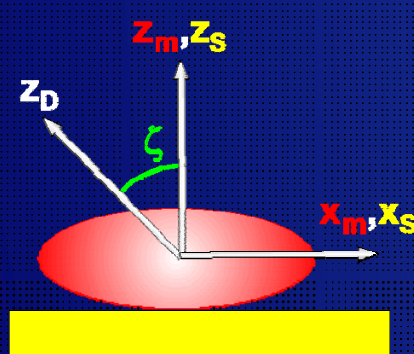
3440 B[G] 3540

Surface model

- tilt angle of director (ζ)
- azimuthal average (ϕ)
- Gaussian disorder profile (σ)

Result of global fit T4L 76:

- $\zeta = 90^\circ$, $\sigma = 15^\circ$

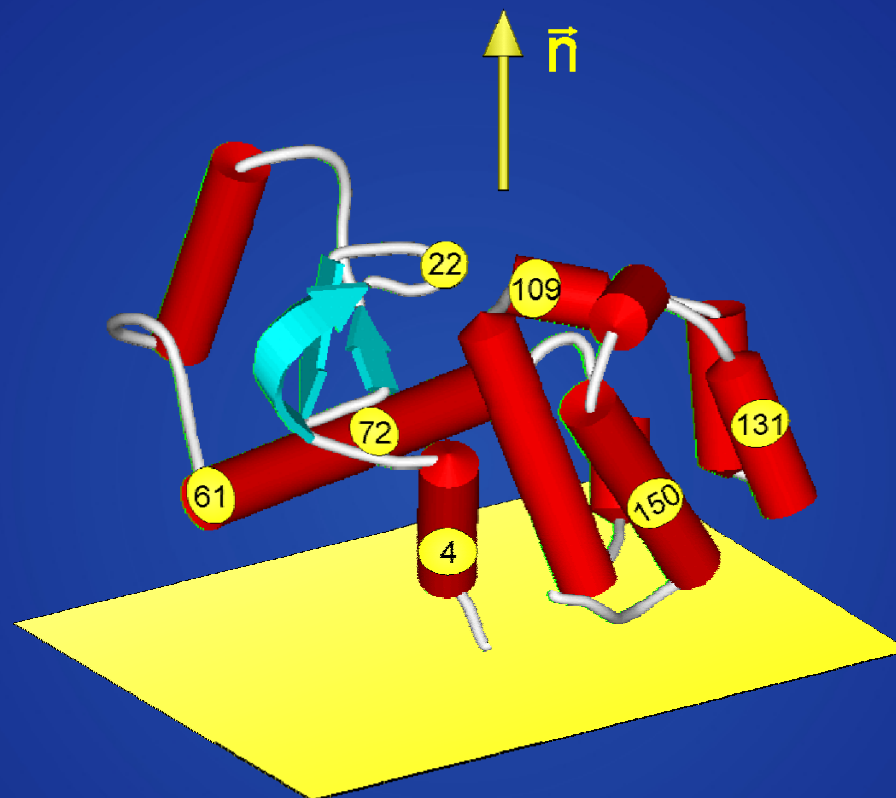
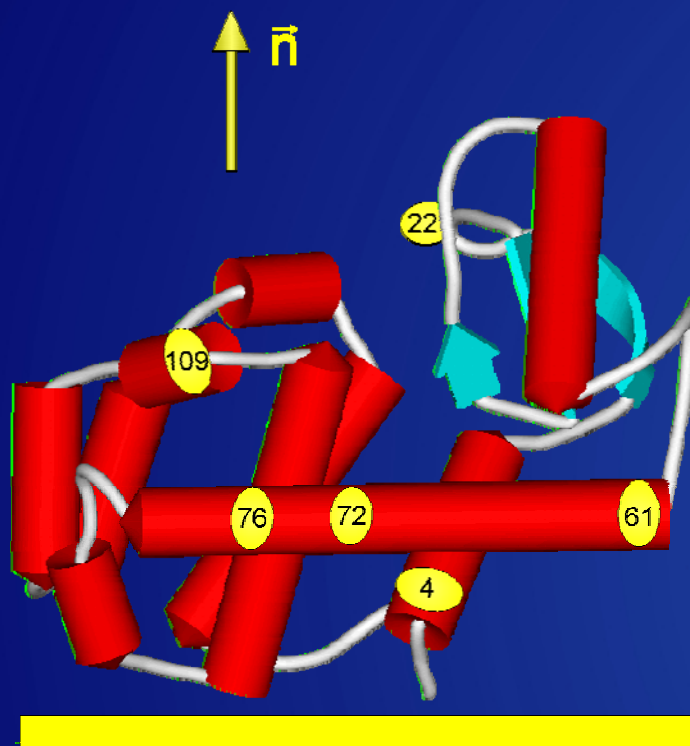


T4Lysozyme on Model Surface

orientation of the molecule on the surface



Max-Planck-Gesellschaft



Conclusions



Max-Planck-Gesellschaft

Vibrations

- IRAS
optimal range from 700 – 3500 cm^{-1} ; restricted to dipole allowed excitations (metal surface selection rule)
- HAS
powerful for low energy vibration such as frustrated translations or rotations
energy resolution approx. 0.1 - 0.5 cm^{-1}
- EELS
complementary to IR; lower frequencies; non dipole transitions possible.
- Raman
comparatively low sensitivity; wide spectral range (50 – 5000 cm^{-1}); enhanced surface sensitivity by plasmon excitation (combination with STM allows resolution down to a few nm)
- SFG
intrinsic surface sensitive; background free detection; suitable for time resolution down to the fs-regime.

Conclusions



Max-Planck-Gesellschaft

Rotations:

- EELS

does only work for very light molecules with large rotational quanta (H_2 , D_2)

- STM,STS

single molecule resolution; rotation (coupling to vibrations. Becomes complex for more complicated molecules.

- EPR

can provide information on the rotational dynamics in the ns-regime; can be used e.g. to characterize large molecules.