CO₂ Adsorption on CaO(001): Temperature-Programmed Desorption and Infrared Study

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Supporting Information

ABSTRACT: We studied adsorption of CO₂ on well-ordered thin CaO(001) films, prepared on Mo(001) and Pt(001) single-crystal surfaces in ultrahigh vacuum (UHV) conditions, using infrared reflection-absorption spectroscopy (IRAS) and temperature-programmed desorption (TPD). At low coverages, CO₂ adsorbs as monodentate carbonates (CO₃²⁻). TPD spectra showed pseudo-first-order desorption kinetics with a maximum shifting from 500 to 470 K with increasing CO₂ coverage. However, at further increasing exposures, desorption maximum is shifted to the considerably higher temperatures (570 K), although CO₂ uptake remained almost the same. This unusual effect was found to correlate with dissociative adsorption of residual water in the UHV background as observed both by TPD and IRAS. Comparative analysis of spectral evolution on crystalline CaO(001) films and CaO nanoparticles favors the model, where surface hydroxyls only affect adsorption geometry of the carbonates rather than form bicarbonate species. However, hydroxyls show stabilizing effect on CO₂ binding to the CaO surface.

1. INTRODUCTION

Carbon dioxide (CO₂) is an abundant chemical feedstock with wide applications in industry. Driven by growing energy and environmental concerns, capturing/sequestering of CO₂ has recently received enormous attention. In this respect, alkaline-earth oxides such as calcium oxide (CaO) were found as promising materials: they are abundant in nature, low-cost, and exhibit high CO₂ uptake and good thermal stability.

Adsorption and reactivity studies on CaO were primarily performed on powders. Complete adsorption of CO₂ results in the formation of calcium carbonate (CaCO₃), which was found to occur in two kinetic regimes. There has been great interest in understanding the initial stages of adsorption. In an early infrared spectroscopy study, the formation of carbonate (CO₃²⁻) species on a CaO surface was proposed as a monodentate at room temperature, with the appearance of bidentate carbonates at higher temperatures. Substantial efforts to understand the IR spectra of CO₂ adsorption on CaO were made via theory. Monodentate carbonates were predicted to form on terrace, edge, and step sites, whereas bidentate carbonates may only be formed on corner sites.

However, full understanding of chemical reactions on CaO surfaces is still missing that renders fundamental studies on well-defined systems a necessity. Experimental studies employing “surface science” methodology remain scarce, mainly because of the difficulties in the preparation of well-defined CaO surfaces as well as high reactivity of CaO toward ambient gases. A synchrotron-based photoemission spectroscopy study of CO₂ adsorption on vacuum-cleaved single-crystal CaO(001) surfaces showed the formation of a surface carbonate species at pressures above 10⁻⁶ Torr. Molecular beam and temperature-programmed desorption (TPD) experiments, performed on a reduced CaO(001) crystal surface, suggested carbonate decomposition to CO. Note, however, that the surface under study showed no clear diffraction pattern that was assigned to high density of defects, presumably oxygen vacancies. Few adsorption studies primarily using photoelectron spectroscopies were performed on CaO thin films grown on Si(111) and Si(001). However, the prepared films were polycrystalline in nature and poorly defined.

In our own laboratories, well-ordered thin films of CaO(001) were prepared on a Mo(001) substrate. Low-energy electron diffraction (LEED) and scanning tunneling microscopy (STM) studies showed films exhibiting properties virtually identical to the bulk, although adventitious Mo segregation at elevated temperatures to the surface has to be controlled. It was shown that the presence of Mo atoms in the sub-surface region may affect surface reactions, in particular with O₂.

Infrared reflection-absorption spectroscopy (IRAS) experiments of initial stages of CO₂ adsorption on CaO(001) films showed vibrational modes consistent with surface carbonates. On the basis of theoretical calculations, the IR
bands were assigned to monodentate carbonate species that first adsorb at steps and other low-coordinate sites, followed by surface islanding of monodentate carbonates adsorbed on terraces. In the continuation of our previous studies, we performed TPD and IRAS study of CO₃ adsorption on CaO(001) films in a wide range of coverages and exposures. In addition to films grown on Mo(001), here we report preparation of well-ordered CaO(001) films on a Pt(001) substrate, which is thought (as a noble metal) to be more resistant toward surface segregation. The Pt(001) surface has a much shorter lattice constant (~2.77 Å) than does Mo(001) (~3.147 Å), thus resulting in ~20% lattice mismatch to CaO(001) (~3.402 Å). In principle, comparative study of the two systems would allow one to examine the support effects, if any, on the surface chemistry of thin CaO(001) films.

2. METHODS AND MATERIALS

The experiments were performed in several ultrahigh vacuum (UHV) chambers. The first chamber is equipped with four-grid LEED optics (from Spec), also used for Auger electron spectroscopy (AES) measurements, a differentially pumped quadrupole mass spectrometer (QMS, Hiden HAL 30I) and an STS (Omicron). The Pt(001) crystal (99.99% from MaTeck GmbH) was mounted on the Omicron sample holder. The sample could be heated by electron bombardment from the backside of the crystal using a tungsten filament. The temperature was measured by a chromel–alumel thermocouple spot-welded at the edge of the crystal. CO₂ (Linde, purity 4.5) was dosed by backfilling the chamber. The heating rate in TPD experiments was 3 K s⁻¹.

The second chamber is equipped with LEED/AES (Spec), and differentially pumped QMS (Hiden HAL 30I/3F). The Pt(001) or Mo(001) single crystal (99.99%, both from MaTeck GmbH) was spot-welded to the two Ta wires for resistive heating as well as cooling that is achieved by filling the manipulator rod with liquid nitrogen. The sample temperature was measured by a K-type thermocouple spot-welded to the backside of the crystal. In this setup, CO₂ was dosed by a directional gas doser placed ~0.5 mm from the crystal surface.

The third chamber is equipped with LEED, STM, X-ray photoelectron spectroscopy (XPS), and an IR spectrometer (Bruker IFS 66i). The Mo(001) crystal was mounted on an Omicron sample holder, with the temperature measured by a thermocouple at the edge of the crystal. CO₂ was dosed by backfilling the chamber. The IRA spectra were recorded using a p-polarized light at an 84° grazing angle of incidence (spectral resolution 4 cm⁻¹).

In all chambers, the metal surfaces were cleaned by cycles of Ar⁺ ion sputtering and annealing in UHV at high temperatures. Residual carbon was removed by mild oxidation in 10⁻⁷ mbar O₂ at 700 K and subsequent flash to 1000 K in UHV. The surface cleanliness was checked by LEED and AES (XPS) prior to the Ca deposition performed from a Mo crucible using a commercial evaporator (Omicron EFM 3). Well-ordered CaO(001) films were grown on Mo(001) as described elsewhere.\(^{15,18}\)

3. RESULTS

3.1. Preparation of CaO(001) Films on Pt(001). We first address the preparation of CaO(001) thin films on Pt(001). Basically, we made use of the same approach as reported for a Mo(001) substrate.\(^{15}\) Ca was vapor-deposited in oxygen ambient (5 × 10⁻⁷ mbar O₂) onto the crystal surface kept at 300 K. The resulted CaO overlayer showed no diffraction spots in LEED. Then, the sample was annealed in UHV at elevated temperatures to induce long-range ordering which was monitored by LEED. This typically requires annealing temperatures of 1000 K and above (see Figure S1 in the Supporting Information). Auger spectra revealed no changes in Ca and O peak intensities on stepwise heating from 700 to 1200 K, indicating that annealing solely caused film ordering and does not affect chemical composition (Figure S2 in the Supporting Information).

Figure 1a,b compares LEED patterns of the well-known “hex”-reconstructed Pt(001) surface and the CaO film, partially covering the metal surface. The diffraction spots of CaO are aligned with those of Pt(001)-(1 × 1), thus suggesting epitaxial growth of CaO(001) on Pt(001) despite of a large (~20%) lattice mismatch. Using diffraction spots of Pt(001)-(1 × 1) as a reference, we obtained 3.35 Å for the lattice constant of the prepared CaO(001) films, which is in fairly good agreement with 3.40 Å of CaO(001). However, on further UHV annealing of this sub-monolayer film to 1200 K, the oxide-related spots disappeared and only the Pt(001)-(1 × 1) diffraction spots remained, indicating film decomposition either via sublimation or dissolution into the Pt crystal. Indeed, preparation of the clean Pt(001) surface after film preparations required numerous sputtering-annealing cycles before LEED patterns start to show “hex”-reconstruction of Pt(001), which is characteristic of a clean surface.

Figure 1c–f shows LEED patterns and STM images of relatively “thick” CaO(001) films annealed at 1000 and 1200 K, respectively. Titration of the Pt atoms by CO revealed <3% of the sample area that might be uncovered. Clearly, UHV annealing at high temperatures leads to a better film ordering as the diffraction spots become considerably sharper. The film morphology studied by STM bears close similarity to the films grown on Mo(001).\(^{16,18}\) In both cases, relatively large, atomically flat terraces having step edges primarily running along the main crystallographic orientations of CaO(001) are observed. Note that there are also substantial amounts of screw bands.
dislocations in the prepared films. Atomic size depressions about 1 Å in depth (see inset in Figure 1f) randomly distributed on terraces can tentatively be assigned to the point-like defects. It appears that the morphology of the CaO(001) films is almost independent on the lattice constant of a metal substrate underneath [2.77 Å on Pt(001) vs 3.15 Å on Mo(001)].

Large-scale STM images showed CaO(001) films uniformly covering a metal substrate. However, some “holes” in the films were observed showing a hexagonal superstructure with an ∼10 Å periodicity. Although atomic resolution was not achieved, such areas could readily be scanned with a low tunneling bias (0.5 V), whereas the regular film surface typically required much higher biases, in the range of 3-4 V because CaO is an electric insulator. These findings allow us to assign such domains to an ultrathin CaO$_2$ layer. Yet, the atomic structure remains unknown, such areas were minority.

### 3.2. TPD Results: Low Coverage

Figure 2a shows TPD spectra recorded on the CaO(001)/Pt(001) films as a function of CO$_2$ exposure up to 0.3 L (1 L = 10$^{-6}$ Torr × s). In these experiments, CO$_2$ was dosed by backfilling the UHV chamber to 5 × 10$^{-9}$ mbar. In order to minimize morphological changes, which may be induced by numerous TPD runs to elevated temperatures, the samples were only heated to 800 K. The original LEED pattern remains almost unchanged after tens of adsorption/desorption cycles, thus indicating no film restructuring in our conditions. As metallic Pt does not adsorb CO$_2$ at room temperature, all desorption features in TPD spectra must be attributed to adsorption on CaO.

Several desorption states can be identified in spectra, which are labeled in Greek starting from the highest desorption temperature. The $\alpha$ peak rapidly saturates and may even be present by reaction with residual CO$_2$ during sample cooling. We assigned the $\alpha$ state to adsorption on yet poorly defined defects. (Interestingly, this peak was not observed on CaO(001) films grown on Mo(001), see below.) The spectra are dominated by the $\gamma$ peak which gradually shifts from 496 to 469 K with increasing coverage. Two other desorption states ($\beta$ centered at 390 K, and $\delta$—at 560 K) appear as the shoulders to the $\gamma$ peak. The peak areas of the $\beta$, $\gamma$, and $\delta$ states obtained by spectral deconvolution are plotted in the inset in Figure 2 as a function of CO$_2$ exposure together with the total uptake, which are all increasing in this coverage regime. Comparison of CO$_2$ desorption spectra measured on the same film sequentially annealed to 1100, 1150, and 1200 K in UHV showed no big differences in spectral evolution in this coverage regime.

In principle, the observed coverage-dependent shift of the peak maximum to the lower temperatures may be indicative of a second-order (recombinative) desorption kinetics that would imply CO$_2$ dissociation. However, this seems to be hardly possible on well-ordered stoichiometric CaO(001) surfaces at pressures applied. In addition, blank experiments with pure CO did not reveal its adsorption on our films. Therefore, CO$_2$ desorption should be simply due to the formation, and the surface recombination reaction would be impossible, if CO could dissociate toward CO during exposure at 300 K. In another case of CO$_2$ decomposition up to elementary carbon, carbon deposits would result in continuous modification of the surface and hence nonreproducibility of the TPD results, which is not the case. Therefore, CO$_2$ dissociation can safely be excluded. Instead, the TPD spectra suggest nondissociative adsorption, with a peak shifting due to coverage effects.

Using a well-known Redhead analysis$^{20}$ of the first-order desorption kinetics using a typical prefactor 10$^{13}$ s$^{-1}$, one can calculate the desorption energies. For the $\gamma$ peak, it decreases from 129 to 122 kJ/mol with increasing CO$_2$ coverage. Accordingly, the estimates of the desorption energies for the $\alpha$, $\beta$, and $\delta$ states yield 185 kJ/mol (700 K), 150 kJ/mol (570 K), and 100 kJ/mol (360 K), respectively. Utilizing the prefactor 10$^{15}$ s$^{-1}$ in calculations results in all desorption energies increased by ∼15 kJ/mol.
On the basis of our previous DFT and IRAS study, focused on initial stages of CO$_2$ adsorption on CaO(001), we can tentatively assign the $\alpha$ state to adsorption on structural defects, whereas the principal $\gamma$ peak reflects CO$_2$ adsorption on terrace sites as monodentate carbonates. Following the formation of surface aggregates (like pairs and chains predicted by DFT), CO$_2$ binding becomes weaker at increasing coverage.

### 3.3. TPD Results: High Exposures

TPD experiments of CO$_2$ adsorption on CaO(001) at high exposures were performed in another chamber using a directional doser. In order to minimize desorption signals from a weakly bonded CO$_2$ “multilayer” (see Figure S3) that may obscure the chemisorbed states, CO$_2$ was dosed at 220 K.

Figure 3a,b shows two sets of TPD spectra obtained for a dense CaO(001) film prepared at 1100 K (a) and for a thinner film prepared at 1200 K, which additionally showed diffraction spots of Pt(001)-(1 x 1) (b). The experiments were carried out as follows. First, the films were several times exposed to 0.5 L CO$_2$ in order to examine data reproducibility. Then, the dosage was stepwise increased to 1, 5, and 10 L. Finally, the sample was again exposed to 0.5 L to compare with the first TPD runs at the same exposure. The spectra on these two films look very similar, although the $\alpha$ peak seems to be more pronounced on a thinner film. Apparently, another low-temperature desorption state exists at around 280 K (labeled $\epsilon$). Nonetheless, the spectra at low exposures (<0.5 L) are still dominated by the $\gamma$ peak.

Drastic spectral changes occur at higher exposures: the $\beta$ peak strongly gains in intensity, whereas the low-temperature signals attenuate, although the total CO$_2$ uptake does not change considerably, thus leading to the isosbestic point at around 500 K. The same behavior was observed for the CaO(001) films grown on Mo(001), as depicted in Figure 3c. (A full dataset is presented in Figure S4 in the Supporting Information.) Obviously, any structural changes caused by consecutive TPD runs can be ruled out because the last spectrum for 0.5 L well reproduces those measured before high dosage experiments.

As discussed above, the initial shift of the $\gamma$ peak to the lower temperatures (Figure 1) reflects gradual weakening of the CO$_2$ binding at increasing coverage. Such a behavior is well documented in the literature for many adsorbates on metal and oxide surfaces. CO$_2$ that becomes more strongly bound at high dosages is unusual. Adsorption of CO$_2$ molecules on top of the carbonate ad-layer or clustering of CO$_2$ molecules into three-dimensional particles, as a possible scenario, would only result in further lowered desorption temperature as intermolecular bonds in solid CO$_2$ (“dry ice”) are weak, thus resulting in the desorption signal at ~150 K (see Figure S3). Therefore, the spectral evolution observed at high exposures cannot be assigned to pure coverage effects. Moreover, the CO$_2$ uptake seems to reach saturation, and the $\beta$ state gains in intensity, in essence, at the expense of low-temperature states.

Bearing in mind that CaO readily reacts with water, we examined possible effects of traces of water in the UHV background on CO$_2$ adsorption. In principle, surface “contamination” with water was thought to be unavoidable even under UHV conditions in experiments with polycrystalline CaO films on Si.

Figure 4a shows desorption traces of CO$_2$ (44 amu) and H$_2$O (18 amu) in four TPD runs, all recorded upon exposure of 0.5 L CO$_2$ at 220 K. Note that cooling the sample down to 220 K after thermal flash usually takes a few minutes.

Comparison of the first two spectra indicates good reproducibility of desorption traces. At this exposure, the $\beta$ state only manifests itself as a shoulder to the $\gamma$ state. The amount of water desorbing between 400 and 500 K is negligible, thus suggesting that CO$_2$ is adsorbed on the clean CaO(001) surface. For the third run, the sample was kept at 220 K for 20 min prior to CO$_2$ exposure. It was thought that during this time, the sample will react with residual gases in the UHV background. This did cause dramatic effect on CO$_2$ desorption spectra which only showed the $\beta$ state. Concomitantly, the amount of water desorbing from the CaO surface is substantially increased. (No other species were observed in the multimass spectra.) The fourth spectrum recorded immediately after the third one well reproduces the first two spectra, with the $\gamma$ state dominating. These experiments show that the $\beta$ state can be formed upon interaction of CO$_2$ with the CaO(001) surface “contaminated” with water-ad-species. On the other hand, in TPD experiments shown in Figure 3, water molecules most likely interact with the surface already covered, at least partially, by monodentate carbonates first formed on the clean CaO(001) surface (see Figure 2). Nonetheless, the final effect is essentially the same: CO$_2$ desorption maximum shifts to much higher temperatures.

In another set of experiments, presented in Figure 4b, TPD spectra were measured for the same (0.5 L) exposures obtained either in 30 or 300 s exposure time by adjusting CO$_2$ pressure. Again, the spectra revealed considerable gain of the $\beta$ peak intensity at the lower CO$_2$ flux. In principle, a longer exposure increases probability for water in the background to react with the surface. However, the observed difference in water desorption signals is not as obvious as in Figure 4a. Since both CO$_2$ and H$_2$O strongly adsorb onto the clean CaO(001) surface, there seems to a competition for adsorption sites which result in complex kinetics.

### 3.4. IRAS Results

To gain information on the nature of ad-species formed during CO$_2$ adsorption at high exposures, we employed IRAS. Scheme 1 illustrates vibrational modes of CO$_2$ identified from our previous IRAS and DFT study of the initial stages of adsorption on the clean surfaces.

The most pronounced band at 1300–1320 cm$^{-1}$ is associated with symmetric stretching ($\nu_4$) of monodentate carbonate (CO$\textsubscript{3}^2$$^-$$^2$). The $\nu_4$ mode shows up as a relatively weak
associated with vibrations of adsorbed CO$_2$. At low CO$_2$ exposures, only the $\nu_4$ and $\nu_5$ modes are detected. Their shift to higher frequencies (from 1298 to 1310 cm$^{-1}$, see the first three spectra) nicely correlates with the shift of the $\gamma$ peak in TPD spectra (Figure 2) and can, therefore, be explained by increased density of monodentate carbonate species accompanied by their aggregation into pairs and linear chains as predicted by DFT.$^{18}$

However, at further increasing CO$_2$ dosage, new bands at 1366 and 1382 cm$^{-1}$ grow up, whereas the bands at 1310 and 1320 cm$^{-1}$ attenuate, thus resulting in the isosbestic point (inset in Figure 5), that is, in a similar way observed during the $\gamma \rightarrow \beta$ transformation in TPD spectra (Figure 3). The $\nu_4$ band at around 990 cm$^{-1}$ becomes broader and shifts to 1020 cm$^{-1}$ basically following the blue shift of the $\nu_5$ band. Last, a new band grows at 1510 cm$^{-1}$, and a small feature appears at 1440 cm$^{-1}$. Importantly, these changes occur simultaneously with the appearance and growth of sharp $\nu$(OH) bands at 3697 and 3506 cm$^{-1}$. Adsorption of isotopically labeled C$^{18}$O$_2$ onto the CaO(001) film (Figure S5 in the Supporting Information) led to all CO$_2$ related bands shifting to lower frequencies as anticipated on the basis of the reduced mass analysis, but not the $\nu$(OH) bands, again appearing at high C$^{18}$O$_2$ dosages. This finding indicates that hydroxyl species originate from adsorption of residual water in the UHV background rather than as impurity in CO$_2$ (expected to be isotopically labeled with the same O as CO$_2$). In any case, combined TPD and IRAS results provide compelling evidence that adsorption of CO$_2$ on the CaO(001) surface is affected even by traces of water in the ambient, the fact that has to be taken into account while discussing CO$_2$ adsorption on the “clean” surfaces and making comparison with theoretical calculations used for proper description of experimentally observed infrared bands.

Water adsorption on CaO was subject of numerous, both theoretical and experimental, studies using in particular IR spectroscopy. For example, the 3695 cm$^{-1}$ band that closely resembles the 3697 cm$^{-1}$ peak in our spectra was observed on CaO powders formed by high-temperature degassing of Ca(OH)$_2$.$^{22}$ This band has been assigned to a “free” hydroxyl to differentiate it from a “bound” hydroxyl that appears a considerably broader band centered at 3550 cm$^{-1}$. Recently, water adsorption has also been studied on the CaO(001)/Mo(001) films using IRAS, XPS, and STM, in combination with DFT.$^{3,24}$ Two $\nu$(OD) bands were detected upon D$_2$O adsorption at 300 K: a sharp peak at 2725 cm$^{-1}$ (3703 cm$^{-1}$ for OH counterpart$^{22}$) and a broad one centered at ~2600 (3516) cm$^{-1}$. On the basis of the DFT calculations, the high-frequency band was assigned to O$_2$D hydroxyls, where the subscript (w) denotes oxygen in dissociated water. Accordingly, the low-frequency band was assigned to O,D hydroxyls, where the subscript (s) designates surface oxygen in the oxide.

Note, however, that calculated values for the second band could only agree with the experimental ones in model structures containing at least 4 water molecules per CaO(001)-(3 $\times$ 4) unit cell.$^{24}$ Therefore, we assign the $\nu$(OH) bands observed in CO$_2$ adsorption experiments at high dosages to hydroxyl species formed directly on the CaO(001) surface. Some deviation from the spectra obtained for pure water adsorption likely originates from that water adsorbs onto the carbonate precoversed surface, not clean. In particular, the low-frequency band (3506 cm$^{-1}$) is much narrower, thus pointing to isolated OH species lacking H-bonds.

To gain further information about the role of water in CO$_2$ adsorption, the sample saturated with CO$_2$ at room temperature was dosed with D$_2$O at 300 K (Figure 6). (Note that the 3000–3450 cm$^{-1}$ region is usually obscured by continuous H$_2$O adsorption on the cold IR detector, which is one of the reasons of using D$_2$O instead of H$_2$O in water adsorption IR measurements.) The results show that OH species undergo H–D exchange with D$_2$O: the $\nu$(OH) bands at 3697 and 3506 cm$^{-1}$ disappear, whereas the $\nu$(OD) bands appear at 2733 and 2573 cm$^{-1}$. Concomitantly, the 1382 cm$^{-1}$ peak is strongly reduced, whereas the 1510 and 1440 cm$^{-1}$ bands gain in intensity. Again, spectral changes in this region show the isosbestic point (see inset). In addition, a broad band at ~1020...
formation of isosbestic point. This O\textsubscript{16}D v s O\textsubscript{18}D).\textsuperscript{25,26} Therefore, the H shift corresponds to oxygen exchange in surface hydroxyls (e.g., sharp band appears at 865 cm\textsuperscript{−1} alleviates, following the behavior of the 1382 cm\textsuperscript{−1} peak, and a new sharp peak appears at 868 cm\textsuperscript{−1}.

The same kind of experiment performed on the CaO(001) surface first exposed to isotopically labeled CO\textsubscript{2}\textsuperscript{18} revealed a similar behavior (Figure 7). Upon the H−D exchange, the band at 1365 cm\textsuperscript{−1} is strongly reduced, whereas the 1495 and 1427 cm\textsuperscript{−1} bands gain in intensity, although the relative ratio is different from their CO\textsubscript{2}\textsuperscript{16} counterparts (Figure 6). Again, a sharp band appears at 865 cm\textsuperscript{−1}. However, a closer look revealed both \(\nu(OD)\) bands red shifted by about 17 cm\textsuperscript{−1} as compared to OD bands observed in experiments with CO\textsubscript{2}\textsuperscript{16} (2716 and 2557 cm\textsuperscript{−1} vs 2733 and 2573 cm\textsuperscript{−1}, see direct comparison in Figure S6 in the Supporting Information). This shift corresponds to oxygen exchange in surface hydroxyls (e.g., O\textsuperscript{16}D vs O\textsuperscript{18}D).\textsuperscript{25,26} Therefore, the H−D exchange is accompanied by O\textsuperscript{16}−O\textsuperscript{18} exchange. Since the latter is only present in CO\textsubscript{2}\textsuperscript{18} ad-species, the following surface reaction seems to occur: 2O\textsuperscript{16}H + CO\textsubscript{2}\textsuperscript{18} + D\textsubscript{2}O\textsuperscript{16} (gas) \(\rightarrow\) 2O\textsuperscript{18}D + CO\textsubscript{2}\textsuperscript{16} + H\textsubscript{2}O\textsuperscript{16} (gas). Indeed, the principal band at 1365 cm\textsuperscript{−1} associated with CO\textsubscript{2}\textsuperscript{16} almost vanishes upon D\textsubscript{2}O adsorption (Figure 7), whereas the CO\textsubscript{2}\textsuperscript{18} related band at 1382 cm\textsuperscript{−1} attenuates to a lesser extent (Figure 6). [Moreover, this reaction may explain the relative increase and the blue shift of the 1495 cm\textsuperscript{−1} band (see Figure 7) as it has now strong contribution from newly formed CO\textsubscript{2}\textsuperscript{16} species.] Certainly, the process resulting in oxygen scrambling and H−D exchange involves a complex reaction of water with ad-layer containing both hydroxyl and carbonates species and yet remains poorly understood.

All in all, the IRAS results show no evidence for the bicarbonate (CO\textsubscript{3}H\textsuperscript{−}) formation in our conditions. Indeed, all bands observed in the 1600−800 cm\textsuperscript{−1} region still belong to vibrations in CO\textsubscript{2}\textsuperscript{16}. No new bands appear in this region upon H−D exchange observed in Figures 6 and 7, which could otherwise be associated with vibrations in COH entities in bicarbonates which strongly shift upon replacing H with D.\textsuperscript{27} On the other hand, a very sharp band at 868 (865) cm\textsuperscript{−1} in the above-presented experiments is very close in frequency to the \(\pi\)-band of monodentate carbonates (875 cm\textsuperscript{−1}, see Scheme 1), which was observed on CaO particles\textsuperscript{18} and not on the CaO(001) films solely because of the surface selection rules in IRAS.\textsuperscript{21} To gain more information on this issue, we performed additional IRAS experiments on a granular CaO film consisting of small particles (about 5 nm in size, see the inset in Figure 8) prepared on Ru(0001).

As oxide surfaces are randomly oriented with respect to the metal plane, all \(\nu_3−\nu_4\) bands associated with monodentate carbonate formed at low CO\textsubscript{2} coverage can be detected (Figure 8). With increasing CO\textsubscript{2} dosage, first the band at 1340 cm\textsuperscript{−1} and then the bands at 1434 and 1524 cm\textsuperscript{−1} gain in intensity, whereas all bands characteristic for low coverage regime are strongly reduced (see also Figure S7a in the Supporting Information). In addition, a small signal shows up at 1054 cm\textsuperscript{−1} together with a sharp band at 865 cm\textsuperscript{−1} which is very close to the original 875 cm\textsuperscript{−1} band, and it is difficult to conclude whether the latter remains or not. The \(\nu(OH)\) region did not show any evidence of the isolated hydroxyls (Figure S7b in the Supporting Information) as observed on a
CaO(001) film (Figure 5). It seems likely that such bands strongly overlap because of surface heterogeneity of CaO nanoparticles, resulting in a broad band which is, in addition, obscured by water adsorption signal in IR detector and hence difficult to resolve.

In principle, all bands in the last spectra presented in Figure 8 bear close similarity to those obtained in Figures 5–7, except their relative intensity. A considerably lower intensities of 1510, 1440, and 868 cm$^{-1}$ bands on a CaO(001) film as compared to that of CaO particles (1524, 1444, and 865 cm$^{-1}$, respectively) can readily be explained by the surface selection rules. Since these bands develop upon adsorption of H$_2$O respectively) can readily be explained by the surface selection rules. Since these bands develop upon adsorption of H$_2$O, most clearly seen on CaO(001) fi

The respective CO$_2$-related species do not involve H (D) ν$^+$ because of interaction with hydroxyl species in proximity. Moreover, this interaction would also affect the ν$_5$ and ν$^+$ vibrational modes, both the frequency and the intensity.

Finally, to link the IRAS and TPD results, we measured IRA spectra after dosing 10 L CO$_2$ at 110 K and thermal flash to specified (stepwise increasing) temperature. For clarity, Figure 9 only shows the spectra obtained after heating above 200 K to disappear due to CO$_2$ desorption. The observed spectral evolution is fully consistent with the above-mentioned hypothesis that surface hydroxyls do not change the carbonate nature of CO$_2$ ad-species, but cause their reorientation (e.g., tilting). In fact, attenuation and final disappearance of the bands in the 1400–1500 cm$^{-1}$ region on heating reflects the desorption of water (see Figures 4 and S4 in the Supporting Information) before CO$_2$ starts to desorb. Overall, water ad-species shows stabilizing effect on CO$_2$ binding to the CaO(001) surface.

4. CONCLUSIONS

Well-ordered CaO(001) films grown on Pt(001) and Mo(001) single crystals were used to study CO$_2$ adsorption by TPD and IRAS. The results show that CO$_2$ first adsorbs as monodentate carbonate (CO$_3^{2-}$), with the adsorption energy of 125 kJ/mol which decreases to 100 kJ/mol at increasing coverage due to agglomeration of carbonates in pairs and chains as predicted by DFT. However, at high exposures, CO$_2$ was found to desorb at considerably higher temperatures corresponding to adsorption energy of 147 kJ/mol. Analysis of TPD and IRA spectra revealed a critical role of residual water in the UHV background on CO$_2$ interaction with CaO. It is found that water molecules readily dissociate within the carbonate ad-layer formed at room temperature. Comparative IRAS study of CaO nanoparticles favors the model, where surface hydroxyls coexist with carbonate species and affect their adsorption geometry rather than form bicarbonate species. Therefore, surface hydroxyls show the stabilizing effect on the carbonate layer. Yet, theoretical calculations remain to be done to examine such scenario.

The results highlight the fact that even traces of water in the ambient atmosphere has to be taken into account while discussing CO$_2$ adsorption onto the “clean” surfaces and making comparison with theoretical calculations used for proper description of experimentally observed IR bands.

ASSOCIATED CONTENT

5 Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acs.jpcc.8b11415.

LEED and Auger spectra of CaO(001) films prepared on Pt(001) as a function of annealing temperature; TPD spectra of CO$_2$ adsorbed at 100 K on CaO(001)/Mo(001) films; and full dataset of TPD and IRA spectra of CO$_2$ on CaO(001) films and CaO particles at high exposures (PDF)

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**REFERENCES**


